

ABSTRACT

Title of Document: MODELING THE IMPACT OF SEDIMENT RESUSPENSION AND FLOCCULATION ON THE FATE OF POLYCHLORINATED BIPHENYLS.

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Hydrophobic organic contaminants (HOCs) are important pollutants in urban estuaries. HOCs include polycyclic aromatic hydrocarbon (PAH), polychlorinated biphenyls (PCBs), and polybrominated diphenyl ethers (PBDEs). Sorption to resuspended particles and sediments plays an important role controlling the water column residence times and spatial distributions of HOC in aquatic environments. Pollutant residence times and the time required to reach sorptive equilibrium are highly dependent on the chemical character, the surrounding environment, and particle types and compositions. If rates of sorption are slow relative to particle residence times, HOC behavior may be described using kinetically-limited partitioning behavior.

In this study, a flocculation model that simulates flocculation of activated carbon, organic carbon, and inorganic solids ranging in diameter from 2 to 1000 μm has been developed. A multi-class flocculation-based contaminant fate model is adapted to

describe desorption kinetics for contaminants associated with flocculated particles during a resuspension event. The model is effective in predicting transport of hydrophobic organic contaminants among different size flocs, water, and two sediment layers. The model also demonstrates the impact of fractal geometry, bottom shear stress, particle composition, floc size, fraction of organic carbon (f_{OC}), fraction of activated carbon (f_{AC}), organic carbon partition coefficient (K_{OC}), and total suspended solids (TSS) on contaminant desorption rate and residence time. Under different scenarios, this model's results support the importance of multi-floc cluster, sediment-water interaction, and of flocculation for the contaminant desorption rate in the water column.

In a floc-rich environment flocculation is an important mechanism redistributing contaminants among flocs. When flocculation is considered in a dynamic particle environment that includes sediment resuspension, settling, and kinetic-limited HOC partitioning, the steady state total PCB concentration in the water column is decreased by 20 % and water column HOC residence time decreased by 36%. When activated carbon is added to contaminated sediments, the total PCB concentration in the water column decreases by 90% (123.4 to 11.4 ng/L). If the activated carbon coagulates with the resuspended sediment, this decrease is partially offset by some activated carbon being entrained in slowly-settling flocs, and the steady-state PCB concentration is 61 ng/L.

MODELING THE IMPACT OF SEDIMENT RESUSPENSION AND
FLOCCULATION ON THE FATE OF POLYCHLORINATED BIPHENYLS

By

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Dedication

To Shih-Yun, Yo-Yi, Yo-Hua, and my grandparents, for love, patience, understanding, and support.

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Earning Ph.D. degree is a collaborative process and there are many people I'd like to thank for their assistance and fruitful discussions about this work. First and foremost I wish to express my sincere gratitude to my advisor Professor Joel E. Baker, for his great insight and dedication in helping with all aspects of my research. He is not only my advisor in academic, but also in my life. I also appreciate the thoughtful comments of my other committees, Professors Walter Boynton, Shenn-Yu Chao, Lawrence P Sanford, and Alba Torrents, whose help and knowledge played a major role for completion of my dissertation. They all taught me to have the right research attitude and encourage me throughout this research. Finally, I would like to thank the present and former members in Baker's lab including Jeff Ashley, Holly Bamford, Eileen Beard, Travis Burrell, Bernie Crimmins, Josh Dixon, Susan Klosterhaus, John Kucklick, Fung-chi Ko, Laura Lockard, Randy Larsen, Dan Liebert, Amy Merten, Taeko Minegishi, John Offenber, Megan Toasperm, Abby Schneider, and Heather Stapleton. Especially Fung-chi helped me to use to and overcome the most difficult times in the USA.

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Chapter 1: Introduction

1.1 Background

Hydrophobic organic contaminants (HOCs) are important pollutants in urban estuaries. HOCs include polycyclic aromatic hydrocarbon (PAH), polychlorinated biphenyls (PCBs), and polybrominated diphenyl ethers (PBDEs). Sorption to resuspended particles and sediments plays an important role controlling the water column residence times and spatial distributions of HOC in aquatic environments (Alkhatib and Weigand, 2001; Wu and Gschwend, 1986; Borglin *et al.*, 1996; Jepsen *et al.*, 1995; Ko *et al.*, 2003; Lick and Rapaka, 1996; Rounds and Pankow, 1990). Pollutant residence times and the time required to reach sorptive equilibrium are highly dependent on the chemical character, the surrounding environment, and particle types and compositions. If rates of sorption are fast *relative* to particle residence times, HOC behavior may be described using equilibrium partitioning. Several water quality models have been developed based on equilibrium partitioning, including TOXIWASP (Ambrose *et al.*, 1987) and DELPCB (DRBC, 2003). However, in highly dynamic particle environments such as algal blooms, shallow water sediment-water interfaces, and dredging operations, traditional water quality model may not accurately simulate the fate of HOCs. Several studies have explored the possible reasons, including that these models do not account for (1) different sorption behavior (Gong and DePinto, 1998), (2) complex particle resuspension and settling (Schneider *et al.*, 2007), (2) different particle composition (Lick and Rapaka, 1996), and (4) competition among different solid types for HOC partitioning (Sun and

Ghosh, 2007). In an urban estuary, both resuspended and sediment particles exist as non-uniform, non-normal size distributions. Contaminant water column residence times and times required to reach a steady state are controlled by the different floc sizes within the rapidly varying boundary condition. Therefore, models must simulate how the entire spectrum of particles varies in space (sediment versus water column) and time (*i.e.*, before, during, and after resuspension events). A new model that can simultaneously simulate the varying particle and contaminant should be developed to better describe resuspended contaminated sediment in urban estuaries. The model should also be able to simulate the dynamics of particle exchange across the sediment/water interface, settling of particles and contaminants through the water column, production of flocs, and HOC sorption and desorption rates. Finally, the model may be used as a tool to explore the behavior of activated carbon when added to sediments as an *in situ* remediation technology.

1.2 *Flocculation and Particle Transport Models*

The process of aggregation resulting from the attachment of particles colliding with each other is called flocculation (O'Melia, 1972). It is an important internal transport processes affecting the particle residence time in the water column. Coagulated particles are formed either in the water column or on the sediment surface, often facilitated by microbial communities (Kiorboe and Hansen, 1993). Recent research has noted the importance of flocculation in the fate of HOCs, including that flocs will affect the fate of HOC by altering particle organic carbon content and therefore HOC partition coefficients and sorption rates (Alkhatib and Weigand, 2001, Wu and Gschwend, 1986; Borglin *et al.*, 1996, Jepsen *et al.*, 1995;

Ko *et al.*, 2003; Lick and Rapaka, 1996; Rounds and Pankow, 1990). The flocculation rate is a function of the stickiness coefficient and collision probability (O'Melia, 1972). The three major mechanisms that control collision probability are Brownian motion, differential settling, and fluid shear stress (Yao *et al.*, 1971). The stickiness coefficient α is defined as the ratio of the particle attachment rate and the particle collision rate. Edzwald *et al.* (1974) indicated that this step is concerned with eliminating or nullifying the repulsive energy barrier that exists between the two particles. In general, exopolymeric material, suspended solids concentration, ionic strength, pH, temperature, algae type, and algae concentration are the major factors that determine the stickiness coefficient (Edzwald *et al.*, 1974, Gibbs, 1983, Kiorboe and Hansen, 1993, Liss *et al.*, 1996; Elmaleh *et al.*, 1996; Rengasamy *et al.*, 1996; Crump and Baross, 2000; Jun *et al.*, 2001; Han and Kim, 2001; Hamm, 2002). The time to form flocs may vary from a few seconds to a few days and depends on the particle character and the surrounding environment, including the water column depth, salinity, shear stress, and temperature (Jackson, 1988).

Smoluchowski (1917) was the first to develop a mathematical model to describe the aggregation of particles. Later, Edzwald *et al.* (1974) solved Smoluchowski's equation as a function of porosity ϕ , velocity gradient G , and the stickiness coefficient, α . O'Melia (1972) described flocculation as a two-step process. Recent studies have expended these concepts into three types of flocculation models. The multi-cluster flocculation model developed by Lick and Lick (1988) simulates the flocs being transported among differently-sized floc clusters. Each floc cluster represents a certain volume-based size of flocs. When two flocs start to form a new

floc, the volume of new floc equals to the sum of previous floc volumes, and the new floc is assigned to the corresponding volume-based size floc cluster. Variables like settling velocity and number concentration vary temporally in each cluster of flocs. A second type of flocculation model uses the steady state floc diameter (D_{SS}) to represent the floc characteristic. Lick and Lick (1988) conducted a series of experiments to examine the relationships among velocity gradient G , steady state flocs' diameter and TSS concentration in Lake Erie samples. They concluded that at steady state a simple approximate relationship among concentration, median diameter and shear stress is described by the product of TSS, shear stress, and square of D_{SS} . This model provides a quick tool to predict the diameter of flocs at steady state based on given TSS and shear stress. Winterwerp (1998) developed a one-dimensional vertical numerical model (1 DV Point model) that uses the median particle size, D_{50} , to represent the entire floc spectrum. This is the first flocculation model to use fractal geometry to describe flocculation. This model first describes the transport of the cohesive sediment's flocculation and includes the sediment's settling velocity. It simulates aggregation and disaggregation processes using sediment concentration and turbulence, reaching a reasonable agreement with observations from settling column tests (Winterwerp, 1998).

1.3 HOC Sorption Behaviors and Related Models

Adsorption and desorption ('sorption') on suspended particles play an important role in the fate of hydrophobic organic compounds (HOCs), like PCBs, PAHs, and PBDEs, in urban estuaries. Four types of partition behavior models have been used to simulate this process, including the equilibrium partition model (Di Toro, 1985), the

single box kinetic model (Oddson *et al*, 1970), the two-stage sorption/desorption model (Pignatello *et al*, 1993; Gong and DePinto, 1998), and the radial diffusion model (Wu and Gschwend, 1986).

Equilibrium partition behavior, the simplest strategy, assumes that HOCs immediately reach equilibrium between the dissolved and particulate phases. This behavior has been widely applied to simulate the fate of HOCs in models like the USEPA's WASP model (Ambrose *et al.*, 1987). With this behavior, the fraction of dissolved and particulate contaminants is a function of total suspended particle concentration (TSS), particle solid density, fraction of organic carbon (f_{OC}), and octanol-water partition coefficient (K_{OW}).

In contrast to equilibrium partitioning models, many studies try to simulate the distribution of HOCs between the dissolved and particulate phases kinetically. The basic form is the so-called single box model, which assumes that the sorption process is a first order function and controlled by a single rate constant, which may be estimated from laboratory experiments (Oddson *et al*, 1970). Therefore, the dissolved and particulate HOCs concentrations exponentially approach the equilibrium condition. While the single-box model only requires a single rate constant, it often does not fit experimental data well before reaching equilibrium (Wu and Gschwend, 1986).

To overcome the limitations of the single-box model, a two-compartment diffusion model has been developed (Pignatello and Xing, 1996). This model separates particles into two compartments: an exterior compartment that has a faster exchange rate, and an inner compartment that has a slower exchange rate. This model

is more realistic than the single box model because particle surfaces exchange more easily with the surrounding water phase while the HOC transport to and from the interior of particles is slow due to distance and retarding factors. Therefore, this model agrees better with experimental data than either equilibrium partitioning or single-box kinetic models (Gong and DePinto, 1998). However, the two-compartment diffusion model requires three independent sorption rate parameters, including the ratio of exterior particle volume to the total particle volume, the sorption rate between surrounding water and the exterior compartment, and the sorption rate between exterior and interior compartments.

The radial diffusion model considers the retardation to sorption resulting from transport within porous particles, including flocculated particles (flocs; Wu and Gschwend, 1986). This model simulates HOC desorption resulting from exchange within and between particles and the surface sediment layer with the surrounding water. Because many studies have showed intraparticle diffusion is the major retarding factor for HOCs in aggregates, the radial diffusion approach has been used in many models (Werner *et al.*, 2006).

1.4 Calibration Data: the STORM experiments

The model developed here was calibrated using Shear Turbulence Resuspension Mesocosms (STORM) tank experiments that mimic resuspension and settling of contaminated Hudson River sediment with realistic bottom shear stress (Schneider *et al.*, 2007). In those experiments, the maximum instantaneous bottom shear stress was about 1 dyne cm^{-2} and the volume-weighted average water column turbulence intensity and energy dissipation rate were 0.55 cm s^{-1} and $0.0032 \text{ cm}^2 \text{ s}^{-3}$,

respectively. Hudson River sediment was added to a depth of 5 cm and allowed to consolidate for 10 days. These experiment included both erosional (net sediment resuspension into clean water) and free settling (no erosion or mixing) periods. During erosion, the mixing paddle continuously generated bottom shear stress for 53 hours to ensure that the water column suspended solids reached steady state. During the subsequent one hour free settling period, mixing was turned off to allow suspended particles to settle. Particle size distribution, dissolved and particulate PCBs, TSS, dissolved organic concentration (DOC), particulate carbon, nitrogen, and chlorophyll *a* were measured throughout the resuspension and settling portions of each experiment. The detailed PCBs measurement methods were reported by Schneider *et al.* (2007).

1.5 Model application: Amendment of PCB-contaminated sediments with activated carbon

As an alternative to traditional methods to clean-up sites with contaminated sediments, several investigators have proposed adding granular activated carbon *in situ* (Zimmerman *et al.*, 2005). This added carbon effectively sequesters PCBs, reducing bioavailability and therefore the risk of the contaminated sediments (Zimmerman *et al.*, 2004). However these methods have been demonstrated mainly in laboratory experiments, and significant operational questions remain prior to large-scale field applications. Of primary importance is the fate of the activated carbon particles in the sediments. Are they stable? Are they susceptible to resuspension? Do they aggregate with sediment particles? To address these questions, the model developed here is modified to include activated carbon as a state variable. Scenarios

are examined in which the activated carbon particles are resuspended and aggregated with eroded sediments to explore the long-term performance of the *in situ* remediation.

1.6 Objectives

The primary objective of the model development and application in this study is to explore the influence of flocculation on the fate of PCBs in a dynamic particle environment that includes sediment resuspension, settling, and kinetic-limited HOC partitioning. Specific objectives addressed in each chapter include:

1. To develop a flocculation model that includes the interactions between organic and inorganic particles in a shallow water estuary.
2. To determine how flocculation affects the water column residence times and sedimentation rates of organic carbon under varying conditions.
3. To quantify the influence of flocculation on the fate of organic contaminants
4. To evaluate the ability of kinetic partitioning models relative to equilibrium models to simulate PCB release during resuspension events.
5. To develop a new linear sorption model based on the concept of a two stage model allowing both HOC concentrations and the amount and character of the sorbent phase to vary in time.
6. To explore the impact of kinetic processes such as flocculation and erosion/deposition on the fate of PCBs; for example, the time to reach a steady state, the desorption rate, the dissolved and particulate PCBs concentrations at steady state, and the temporal variation in the deviation from sorption equilibrium during the resuspension event.

7. To discuss the differences between D_{50} and multi-floc-size models. This objective responds to the conclusion from DePinto (1998) that the radial diffusion model is too simple and does not reflect the composition and properties of real sediment.
8. To use the model to explore the behavior of activated carbon when added to sediments as an *in situ* remediation technology

1.7 Strategies

In an urban estuary, both resuspended and sediment particles show non-uniform, non-normal size distributions. Contaminant residence times, or times required to reach a steady state, are influenced by the different floc size within the rapidly varying sediment-water boundary. Therefore, models need the ability to simulate how the entire spectrum of particles varies in space (sediment versus water column) and time (*i.e.*, before, during, and after resuspension events). A new model that can simultaneously simulate the varying particle and contaminant should be developed to better describe resuspended contaminated sediment in an urban estuary. The model should also be able to simulate the dynamics of particle exchange across sediment/water interface, settling of particles and contaminants through the water column, production of flocs, and HOC sorption and desorption rates.

In Chapter 2, development a new flocculation model that simulates the interactions between organic and inorganic particles ranging in diameter from 2 to 1000 μm in 1 μm intervals during resuspension events is described. Each floc cluster is a state variable and represents a specific size of floc. This model simultaneously calculates temporally variable, floc size-dependent properties, including number and

organic carbon concentrations, stickiness coefficient, floc density, settling velocity, fraction of organic carbon (f_{OC}), and settling velocity under different conditions. The major model equations include the concepts from Smoluchowski (1917), Lick and Lick (1988), O'Melia (1972), Kiorboe (1993), Winterwerp (1998), and Grant et al. (2001). The flocculation process is based on the same concept and equations used in Lick and Lick (1988). Flocculation involves two processes; gaining from the smaller particles and losing by forming larger particles by inter-particle collisions. There are several approaches to estimate the floc settling velocity. In this study, the settling velocity equation was adapted from Winterwerp's model (1998) because fractal geometry was used to describe the flocs mass, volume, and porosity. Floc porosity is an important parameter in this dissertation and is calculated for each floc cluster depending on its composition. Porosity is calculated as a function of fractal factor, the smallest particle size, and floc diameter (Kranenburg, 1994). Furthermore, Khelifa and Hill (2006) suggested that the stickiness coefficient and the dry floc density should have a linear relationship with f_{OC} and this concept has been applied to this model. The fraction organic carbon (f_{OC}), which is calculated as the ratio of organic carbon over total "dry" mass for each particle size at each time step, ranges between 0 and 0.5 and equals half of the fraction organic matter (f_{OM}).

Chapter 3 of this dissertation discusses the impact of flocculation on the fate of PCBs during resuspension events in an urban estuary. There are two pools in this study, those being the water column and the sediment. In Chapter 2, the model focused on the exchange of particles between the water and sediment surface layers and the related flocculation mechanisms. However, many HOC fate models have

emphasized the importance of a deeper sediment layer on the long term HOC fate (DRBC, 2003). Therefore, it is necessary to include the deeper sediment layer and related surface-deeper sediment exchange mechanisms, and sediment porewater in this study (Figure 1.1). Suspended flocs are formed in the water column, and HOCs in the intra-particle fluid is assumed to immediately reach partitioning equilibrium within the particles. In addition, we assume that the HOC concentration is uniform inside the floc. To achieve the study objectives, a new kinetically-limited diffusion model has been developed. Dissolved HOC concentration, particulate HOC concentration and mass transfer velocity are all calculated for each cluster simultaneously at each time step. This model tracks the fluxes of particle-related HOCs among all particle size clusters, including both flocculation and resuspension-settling. This design also includes the effect of intra-particle transport. To include the retarding factor from transport within the flocs, the model adapts the concept of an effective diffusion coefficient from the radial diffusion model. Furthermore, to include bioturbation and mixing from the deeper sediment layer, the model simultaneously simulates HOCs being transported among the water column, top sediment layer (0.1 cm), and deeper sediment layer (4.9 cm).

In Chapter 4, the previously developed dynamic sediment-water exchange/flocculation/HOC partitioning model was modified to include activated carbon as a state variable (Figure 1.2 and 1.3). To simplify model parameters and environmental settings, the STORM experiment conditions were also selected as the scenario background in this study. Two model stages are designed and used this modified model as a tool to simulate the competition between activated carbon and

organic carbon in resuspended activated carbon-amended sediment. Further, this model was also used to explore the fate of HOC in the water column when activated carbon flocculates with sediments.

1.8 *Executive Summary*

1.8.1 Chapter 2

The model predicted that the water column residence time of TSS and organic carbon decrease with increased flocculation and fractal factor. Further, with increased flocculation, the efficiency to form larger flocs which have faster settling velocities than small flocs is increased. Therefore, the mean size of particles (D_{50}) in the water column decreases and the particle size distribution is shifted by losing larger flocs.

Several settling velocity strategies were tested based on the same initial conditions for a heterogeneous source of particles. A fractal geometry-based settling velocity equation (Winterwerp, 1998) agreed best with the STORM observations. The results also suggested that using a modified but temporally-invariant particle density over-estimates the organic carbon settling velocity in organic carbon rich environments due to the influence of flocculation on settling velocities. Including flocculation in the multi-size scenario increases TSS settling velocity ($W_{s\ TSS}$) 55% compared to the single size scenario. This result suggests that modeling settling with a single particle of size D_{50} may underestimate the $W_{s\ TSS}$ for the broad size distribution of cohesive flocs.

The organic carbon size distribution is another important factor influencing the water column residence time, gross settling velocity, and steady state concentration of solids and organic matter. In this study, we found that the relationship among individual floc cluster density, gross organic carbon, and TSS settling velocity were impacted by the f_{OC} distribution. They have the same settling velocity and individual floc cluster density only when f_{OC} does not vary with particle size and is evenly distributed over all sizes of particles. The settling velocity of organic carbon plays an important role in predicting the fate of particle-reactive chemicals in the water column. When a wide spectrum of heterogeneous flocs is present, it is necessary to apply a multi-size-floc strategy to simulate gross particle behaviors.

1.8.2 Chapter 3

A multi-class flocculation-based contaminant fate model was adapted to describe desorption kinetics for contaminants associated with flocculated particles during a resuspension event. The model was effective in predicting transport of hydrophobic organic contaminants among different size flocs, water, and two sediment layers. The model also demonstrated the impact of fractal geometry, bottom shear stress, particle composition, floc size, f_{OC} , organic carbon partition coefficient (K_{OC}), and TSS on contaminant desorption rate and residence time. Under different scenarios, this model results support the importance of multi-floc cluster, dirty sediment erosion-deposition, and of the flocculation process on the contaminant desorption rate in a water column. Comparing the equilibrium and kinetically-limited models, both behaviors predict the same dissolved and particle contaminant concentrations at steady state. However, during the first three hours of a simulated resuspension event,

the equilibrium and radial diffusion models overestimated the PCB desorption rate by 50% and 20% respectively. This result suggests equilibrium behavior may not be the best choice for prediction of desorption kinetics during fast events, like dredging, tidal events, or storm water runoff. The radial diffusion model, a common tool to describe desorption kinetics for a single floc, is limited by several factors during a resuspension event, as it fails to include the contaminant exchange with surrounding flocs, it has numerical difficulties in calculating the impact of various boundary conditions, and it ignores indirect impacts from sediment-water exchange.

Further, in a floc-rich environment flocculation is an important mechanism redistributing contaminants among flocs. When flocculation is considered in a dynamic particle environment that includes sediment resuspension, settling, and kinetic-limited HOC partitioning, the steady state total PCB concentration in the water column is decreased by 20 % and water column HOC residence time decreased by 36%.

1.8.3 Chapter 4

Activated carbon (AC) has a stronger sorption capacity than flocs of the same size and organic carbon content. However, AC has a slower sorption rate (smaller mass transfer velocity) due to its relatively lower porosity and higher K_{AC} value. Over long time scales, a majority of HOCs will be sequestered in the activated carbon. However, in the absence of flocculation in a dynamic erosional and depositional environment, the initial HOC distribution in the sediment and the relatively slower sorption rate for activated carbon are the dominant factors controlling the distribution of HOCs in the water column. The interaction time becomes a very important issue under conditions

when activated carbon does not have enough time to adsorb or compete with other organic particles. The above analysis considers the case where amended activated carbon does not interact with natural sediments or other aquatic particles to form flocs. We found that the lower porosity and higher settling velocities of activated carbon relative to natural organic carbon result in lower steady-state activated carbon concentrations in the water column and longer times to sorptive equilibrium.

With enriched AC in sediment, the total water column HOC concentration significantly decreases. Further, when AC aggregates with sediment particles, this flocculated AC has the same physical properties as the floc, resulting in a slower settling velocity, longer residence time, and higher AC-associated HOC concentration in the water column. When activated carbon is added to contaminated sediments, the total PCB concentration in the water column decreases by 90% (123.4 to 11.4 ng/L). If the activated carbon coagulates with the resuspended sediment, this decrease is partially offset by some activated carbon being entrained in slowly-settling flocs, and the steady-state PCB concentration is 61 ng/L.

1.9 Figure Captions

Figure 1.1: The conceptual diagram of the HOC fate model

Figure 1.2: The strategy diagram of the flocculation model

Figure 1.3: The strategy diagram of the PCBs model

Figure 1.1:

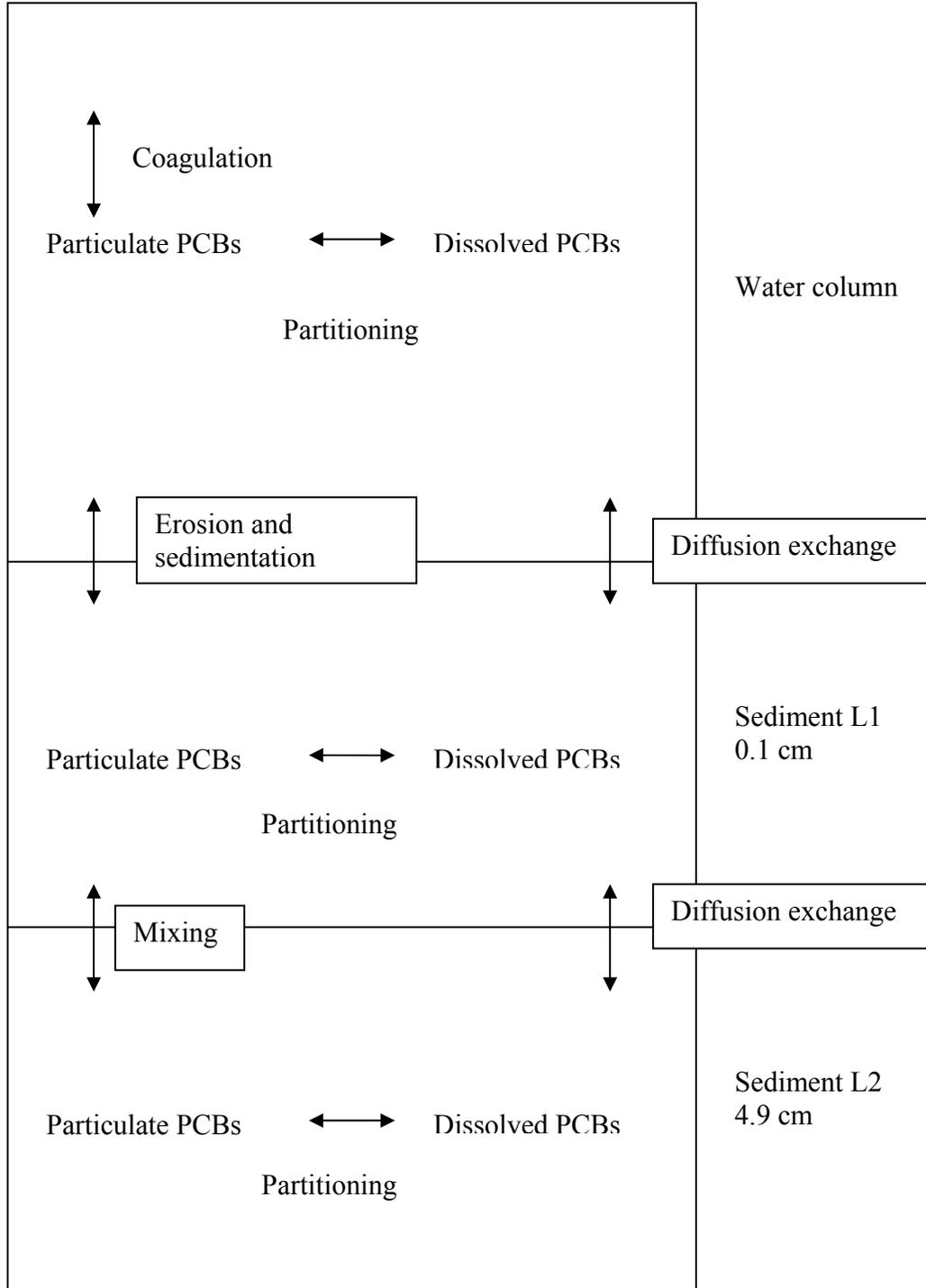


Figure 1.2:

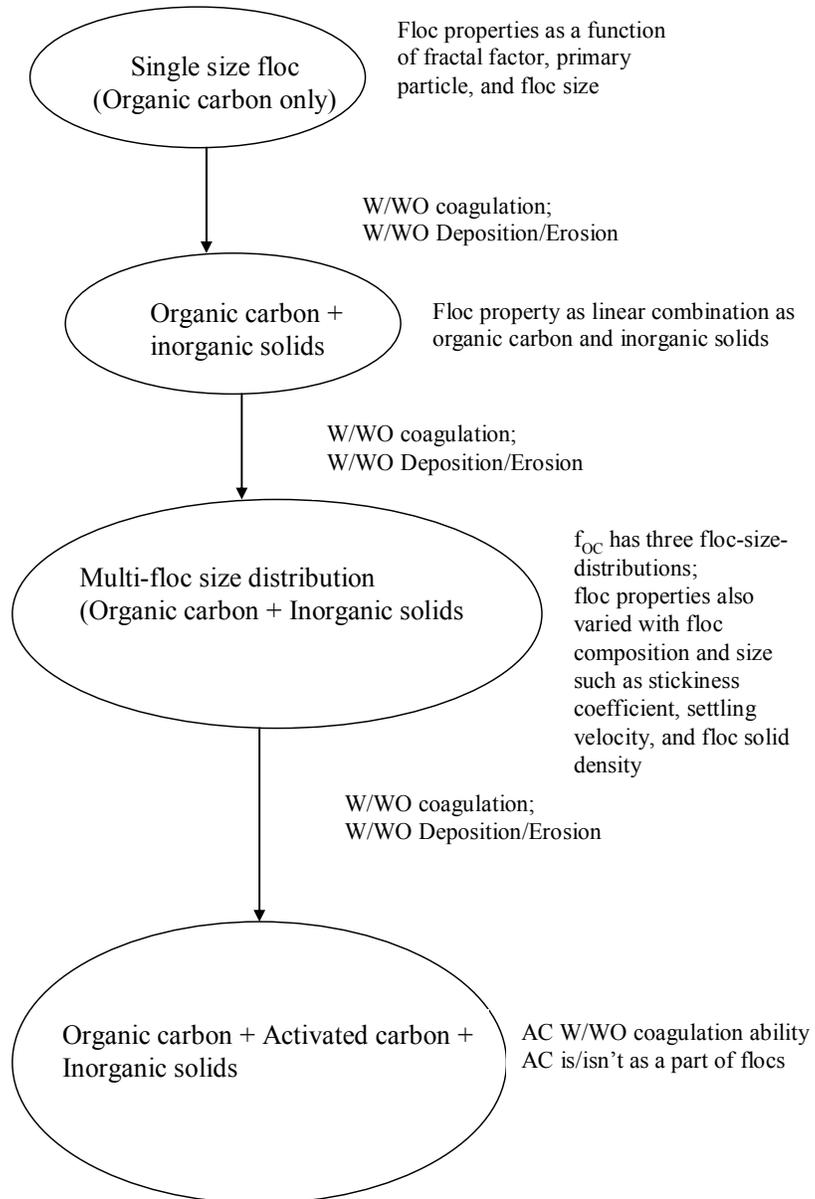
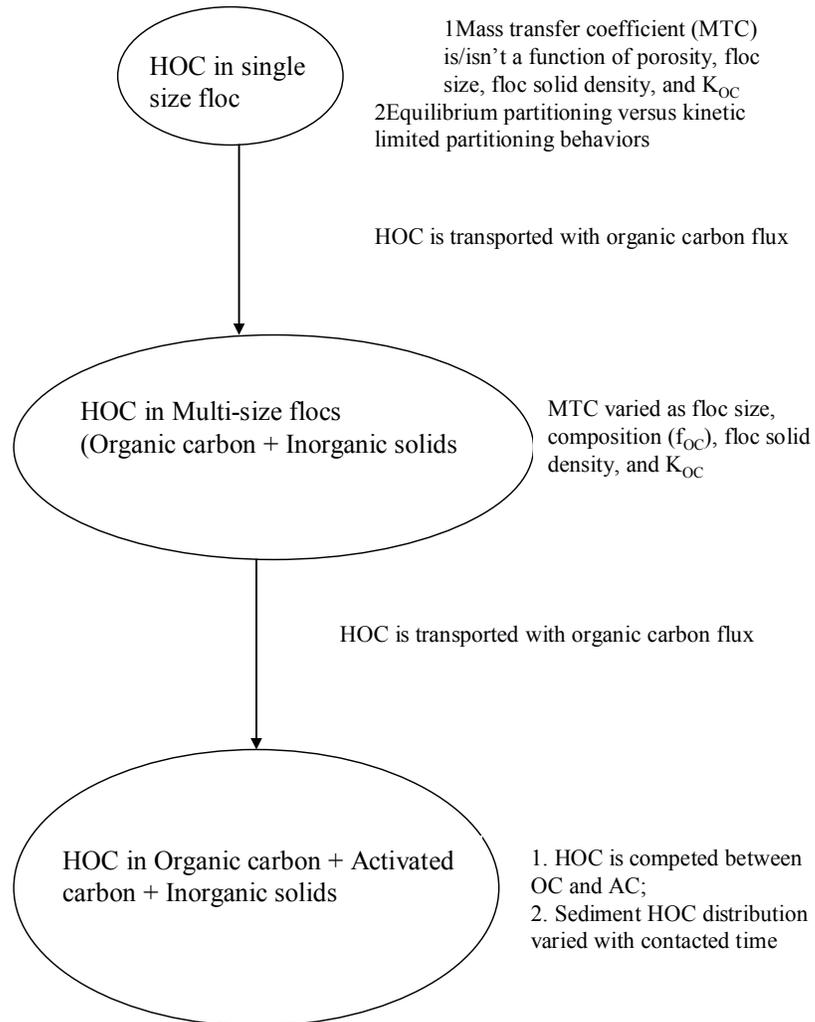


Figure 1.3:



Chapter 2: Modeling the Impact of Flocculation on the Fate of Organic and Inorganic Particles during Resuspension Events in an Urban Estuary

2.1 Abstract

Organic particles play an important role in the fate of organic contaminants in natural waters. These organic particles move by sedimentation and erosion and their size distribution may be influenced by flocculation and disaggregation. Previous organic carbon transport models assume organic particles have the same settling velocity either as slow as phytoplankton or as fast as total suspended solids (DELPCB, 2003 and Chang 2002). Di Toro (2001) used sediment oxygen demand to back-calculate the field organic carbon settling flux. He suggested the effective net organic particle settling velocity should be between these two types of particles, which hints that flocculation plays an important role in organic carbon settling in a heterogeneous particle environment. Later studies also suggested flocculation is a major factor controlling the suspended particle residence time in the estuarine turbidity maxima (ETMs) zone, where significant amounts of inorganic solids, plankton, and organic detritus are trapped (Rengasamy *et al.*, 1996; Sanford *et al.*, 2004). Two types of flocculation models that have been developed are the multi-cluster flocculation model (Lick and Lick, 1988; O'Melia and Tiller, 1993; Jackson, 1995) and the D_{50} flocculation model (Winterwerp, 1996). These models assume that flocs are composed of a single primary particle type and that floc properties are the same for all sizes of flocs, in contrast to many field observations (Lee, 2004). In this study, a new flocculation model that simulates the flocculation for both organic carbon and inorganic solids ranging in diameter from 2 to 1000 μm has been developed.

This model simultaneously calculates temporally variable, floc size-dependent properties, including the organic carbon content, density, stickiness coefficient, and the settling velocity under different conditions. The model was calibrated using Shear Turbulence Resuspension Mesocosms (STORM) tank experiments that mimic resuspension and settling of contaminated Hudson River sediment with realistic bottom shear stress (Schneider et al., 2007). The objective of this study is to develop a flocculation model that includes the interactions between organic and inorganic particles in estuaries. In addition, including flocculation of heterogeneous particle populations will improve models of organic contaminants in estuaries.

2.2 *Introduction*

Flocculated particles (flocs) are cohesive particles formed in the water column or on the sediment surface by aggregation of the complex matrix of microbial communities, organic detritus, and inorganic particles (Dyer and Maning, 1999). Flocs are fragile (Krone, 1962) and have higher fractional organic carbon (f_{OC}), porosity, water content, contact area and intraparticle viscosity compared to same-sized solids. Aqueous flocs have been studied since the 1970s because of their importance in water treatment. Recent research notes the importance of flocs in the fate of hydrophobic organic contaminants (HOCs) including their impact on HOC partitioning and sorption rates, and on particle organic carbon contents (Alkhatib and Weigand, 2002; Wu and Gschwend, 1986; Borglin *et al.*, 1996; Jepsen *et al.*, 1995; Lick and Rapaka, 1996; Rounds and Pankow, 1990).

The importance of flocs on the fate of natural organic matter has also been reported by Kiorboe and Hansen (1993), Hill (1998), Kiorboe *et al.* (1998), Serra and Logan (1999), and Droppo (2001). Richardson and Jackson (2007) demonstrated that floc

formation scavenges picoplankton from surface waters, resulting in an additional important organic carbon source for zooplankton in the deep ocean. Tiselius *et al.* (1998) and Peperzak *et al.* (2003) indicated that although single diatoms settle slowly their aggregates rapidly settle and are enriched in carbon content. Variable flocculation may explain the wide range in organic carbon settling rates reported in field studies. For example, Graf and Rosenberg (1997) report carbon settling fluxes ranging from 0.025 to 70 gCm⁻² d⁻¹, and those of large marine snow and fecal pellets settling velocity can exceed 100 m d⁻¹ (Stemmann *et al.*, 2004).

When inorganic solids aggregate with organic matter, these flocs exhibit different sedimentation rates and residence times. Rengasamy *et al.* (1996) reported the effectiveness of adding clay to control algal blooms in open water. Sanford *et al.* (2004) also suggested flocculation is the major factor controlling the suspended particle residence times in the estuarine turbidity maxima (ETM) zone, where significant amounts of inorganic solids, plankton, and organic detritus are trapped.

The net result of particle aggregation is a wide, mixed heterogeneous particle distribution in the water column, with size-variable organic carbon settling velocity, f_{oc} and effective density. Flocculation may play an important role in transporting organic carbon and associated pollutants through the water column. When organic substrates surround inorganic solids, the settling velocity of organic carbon is influenced by the properties of both the organic and inorganic substrates. Many researchers have tried to find a universal strategy to simulate organic carbon sedimentation rate. Some particle transport models use rules of thumb to estimate an invariant sedimentation rate, such as 0.1 m d⁻¹ for biotic particles and 1.0 m d⁻¹ for particulate organic particles in the

DELPCB model (DRBC, 2003). Other models apply field-measured suspended solid dry density in the Stokes law equation to estimate the particle settling velocity (Chang *et al.*, submitted). Alternatively, Di Toro used Sediment Oxygen Demand (SOD) to constrain organic carbon settling velocity delivering oxygen-depleting materials to the sediment surface. That analysis suggests that the organic carbon settling velocity lies between those of discrete clay particles and algal cells (Di Toro, 2002). Therefore, a model that simulates the flocs composed by organic detritus and inorganic solids can be a better tool to simulate the fate of organic carbon in natural waters.

The process of the aggregation of particles resulting from the attachment of those particles colliding with each other is called flocculation. Salinity, pH, shear stress, total suspended solids concentration (TSS), and the floc character, density and porosity affect the flocculation rate (O'Melia, 1972; Farley and Morel 1986; Lick and Lick, 1988; and Lick *et al.*, 1993). O'Melia (1972) described flocculation as a two-step process: particle transport resulting in collisions (as parameterized by the collision probability β) and particle destabilization (*e.g.*, the probability of each collision resulting in an aggregate, as parameterized by the 'stickiness coefficient' α).

There are three major mechanisms that control the collision probability: Brownian motion, differential settling, and fluid shear, which depend on temperature, particle size, dynamic viscosity of the fluid, shear stress, effective particle density, and particle settling velocities (Ali *et al.*, 1985, Burd and Jackson, 2002; Jackson, 1995; Lick and Lick, 1988; Lick *et al.*, 1993; Dyer and Manning, 1999). In estuaries, Brownian motion plays a minor role in particle transport relative to those processes driven by turbulent mixing.

The stickiness coefficient α is the ratio of the particle attachment rate to the particle collision rate (O'Melia, 1972). Edzwald *et al.* (1974) indicated that this step is concerned with eliminating or nullifying the repulsive energy barrier that exists between two approaching particles. Environmental factors like ionic strength, pH, salinity, temperature, and flocs composition factors (*i.e.* exopolymeric material, algae type, and algae concentration) are the major controls determining the stickiness coefficient (Edzwald *et al.*, 1974, Gibba, 1983, Kiorboe and Hansen, 1993; Winterwerp, 2002). Ali *et al.* (1984) reviewed and summarized reported stickiness coefficients, which range from 0.01 to 1 depending on the transport method, coagulant, and colloid composition. The effect of organic substances on the stickiness coefficient is complex. Gibbs (1983) compared sediment flocculation rates with and without organic substances such as humic acid and dissolved organic carbon (DOC). The sample with the organic substance removed had a stickiness coefficient α four times larger than the natural sample for salinity ranges between 0.6 to 20‰. Weilenmann *et al.* (1989) observed that humic materials act as a stabilizing agent in Swiss lakes. But, algae like diatoms and green algae often act as strong destabilizing agents in the summer. Furthermore, the way that the stickiness coefficient relates to diatom concentration is controversial. Some studies determined that the diatom stickiness coefficient increases during blooms, while others have observed that the coefficient actually decreases during this period because diatoms release mucus that prevents the diatom from sticking together.

Smoluchowski (1917) was the first to develop a mathematical model to describe the aggregation of particles. Later, Edzwald *et al.* (1974) solved Smoluchowski's equation as a function of porosity ϕ , velocity gradient G and the stickiness coefficient α . O'Melia

(1972) described flocculation as a two-step process. Recent studies have expanded these concepts into three types of flocculation models. The multi-cluster flocculation model developed by Lick and Lick (1988) simulates the flocs being transported among differently sized floc clusters. Each floc cluster represents a certain volume-based size of flocs. When two flocs form a new floc, the volume of new floc equals the sum of the previous floc volumes, and the new floc is assigned to the corresponding volume-based size floc cluster. Variables like settling velocity and number concentration vary temporally in each cluster of flocs. Later Jackson (1995) applied this concept to simulate organic carbon flocculation within algae populations. He extended this concept and assumed that flocculation involves the balance between gain from smaller particles and loss to form larger particles by inter-particle collisions.

The traditional multi-cluster flocculation models use either volume (O'Melia and Tiller, 1993) or solid mass (Jackson, 1995) as the floc cluster unit interval. In other words, floc cluster N 's volume or mass equals N times the volume or mass of floc cluster $N=1$. Therefore, the particle diameter does not increase linearly with floc cluster number N . This method simplifies tracking the conservation of volume or mass, but requires very large arrays in the computer program. Therefore, the model can simulate only a very limited floc size spectrum. Furthermore, these models only simulate the flocculation by either pure organic or inorganic solids and lack the ability to describe a wide, mixed heterogeneous particle distribution.

A second type of flocculation model uses the steady state mean floc diameter (D_{SS}) to represent the floc characteristic. Lick and Lick (1988) conducted a series of experiments to examine the relationships among velocity gradient G , steady state floc

diameter and TSS concentration in Lake Erie samples. They concluded that at steady state a simple approximate relationship among concentration, median diameter and shear stress is described by the product of TSS, shear stress, and square of D_{SS} . This model provides a quick tool to predict the steady state flocs' diameter based on given TSS and shear stress.

Winterwerp (1998) developed a 1 DV Point flocculation model that uses the median particle size, D_{50} , to represent the entire floc spectrum. This is the first flocculation model to use fractal geometry to describe flocculation. This model first describes the transport of the sediment's flocculation and includes the sediment's settling velocity. The model was tested against field measurements in the estuarine turbidity maxima (ETM) along the border between Netherlands and Germany. It simulated the aggregation and disaggregation processes using sediment concentration and turbulence, reaching a reasonable agreement with observations from settling column tests. This model can properly describe the impact of flocculation and disaggregation on D_{50} , and has a good agreement with calibrated data. This model also provides a equation to simulate the relationship between fractal geometry and flocs settling velocity.

However, Winterwerp's model has a few disadvantages. First, there are more than ten empirical parameters and formulations in this model. To adjust these parameters to fit our study would require several experiments to determine parameter values. Second, this model assumes that f_{OC} and the stickiness coefficient are constant throughout the entire period. This assumption goes against field and mesocosm observations, where measured particle characteristics vary both temporally and spatially (Ko *et al.*, 2003; Richardson and Jackson, 2007). Therefore, the model could misestimate the flocs density and the

flocculation rate. Third, unlike the O'Melia and Lick models, this model simulates the mean particle diameter of the flocs instead of that of the individually sized class particles, and assumes that all sizes of flocs have the same porosity, composition, and density. However, the observations by Ko *et al.* (2003) demonstrate that neither D_{50} nor D_{SS} can accurately represent the behavior and characteristics of all sizes of organic carbon in a heterogeneous mixture of particles. Flocculation between organic matter and inorganic particles results in flocs with time-variable characteristics, such as f_{OC} (Ko *et al.*, 2003), stickiness coefficient (O'Melia and Tiller, 1993), and effective floc density (Khelifa and Hill, 2006), which in turn affects the erosion and settling flux of organic carbon (Burns and Rosa, 1980). Again, these reports indicate the necessity of simulating these kinetically varied processes to better predict heterogeneous particle behavior.

Several strategies have been suggested to model particle collision rates and flocculation rates, including fractal geometry (Winterwerp, 1996), disaggregation (Alldredge *et al.*, 1990), a curvilinear collision kernel (Han and Lawler, 1991), the adjusted settling velocity equation (Allen, 1985), and the inclusion of heterogeneous particles (Jackson, 1995; Winterwerp, 1996; Hill, 1998). Fractal geometry has recently been used to describe the structure, porosity, and settling velocity of flocs (Winterwerp, 1996). Jackson (1998) and Flesch *et al.* (1999) have combined the fractal factor concept with a multi-cluster flocculation model to derive relationships among particle length, mass, and fractal scaling. Aggregates are modeled as fractals comprised on primary particles, and the fractal factor value varies between one and three (Mandelbrot, 1983). Flocs are assumed to be composed of pure inorganic solids (Yao *et al.*, 1971) or algae cells (Jackson, 1995). Aquatic particles are composed of many different solids including

clay and plankton. Many fractal factor indexes have been reported to address this fact, such as the Sierpinski Carpet fractal dimension D_{SC} , and the fractal dimension of the pore boundaries D_B (Lee, 2004). Several authors have employed unique methods to estimate fractal factors for field-collected aggregates, including a multi-stage-fractal-factor (Li and Logan, 1997), a gross field fractal factor estimated using a computer image technique (Lee, 2004), and a linear combination between oil and clay (Sterling *et al.*, 2005). Many authors suggest the fractal factor value in estuarine water ranges from 1.7 and 2.3 (Winterwerp, 1998). However, Sterling *et al.* (2005) reported the fractal factor for clay-oil system was higher than previous values and between 2.6 and 3.0.

Disaggregation is another important issue in many flocculation simulation models. However, Alldredge *et al.* (1990) concluded that physical disaggregation is important only in the upper ocean layer during storm events when the energy dissipation rate ε is larger than $10^{-3} \text{ cm}^2 \text{ s}^{-3}$. Stemmann *et al.* (2004) also suggested that disaggregation is more likely controlled by biological mechanisms such as bacteria-mediated dissolution and biologically-derived shear stress. Several disaggregation equations have been reported. Most of them use a similar form as the flocculation equation but with a disaggregate coefficient replacing the stickiness coefficient, and the floc is broken into two, three, or four pieces of smaller floc. In the absence of erosion or other external particle sources, the particle size distribution shifts to smaller particles (the D_{50} value decreases) when disaggregation is stronger than aggregation.

Several floc settling velocity equations have been developed by including a shape factor, effective viscosity, hindered effect, Reynolds number, or fractal dimension factors to modify the Stokes law equation (Allen, 1985; Sanford and Halka, 1993; Schnoor,

1996; Winterwerp, 1998). Among these equations, the method developed by Winterwerp agrees well with observed data (Winterwerp, 2002).

A heterogeneous particle size distribution may also impact the flocculation mechanism and alter the organic carbon settling velocity. Crump and Baross (2000) used an Owens tube to compare the relationship between settling velocity and the percentage of organic carbon in particles from flood tide estuarine turbidity maxima (ETM) samples. McCave (1984) also used multi-effective density equations for differently sized particles in his sedimentation model to match the observed data. These reports all suggest a wide, mixed heterogeneous particle distribution in the water column and that the organic carbon settling velocity, f_{OC} and effective density vary in each size of heterogeneous particles.

2.3 *Model Developments*

2.3.1 Objectives

The first objective of this study is to develop a flocculation model that includes the interactions between organic and inorganic particles in a shallow water estuary. The second objective is to determine how flocculation affects the water column residence times and sedimentation rates of organic carbon under varying conditions.

2.3.2 Model Processes

2.3.2.1 General Model Structure

In this study, the flocculation model simulates the movement and concentration of flocculated particles in the water column and in the top sediment layer. The major model equations include the concepts from Smoluchowski (1917), Lick and Lick (1988), O'Melia and Tiller (1992), Kiorboe and Hansen (1993), Winterwerp (1998), and Grant *et*

al. (2001). Equation 2.1 describes how the number of particles per volume per time varies with settling velocity, flocculation, and bed-water exchange.

$$\frac{\partial N_i}{\partial t} = F_i + \left(\frac{W_{s,i}}{x} \times N_i \right) + E_{b,i} \quad \text{Equation 2.1}$$

where, N is number of particle of class i per volume, $W_{s,i}$ is the floc setting velocity, F_i is the flocculation effect on particle number balance, x is the water column depth, and $E_{b,i}$ is the pure sediment resuspension flux that is calculated with a sediment erosion model.

The flocculation process is based on the same concept and equations used in Lick and Lick (1988). Flocculation involves two processes; gaining from the smaller particles and losing to form the larger particle by inter-particle collisions. The first term on the right-hand side Equation 2.2 represents the rate of formation of flocs of size k by cohesive collision between particle sizes i and j . The second term is the loss of size k flocs or solids by cohesive collision with other size particles.

$$F_k = \frac{1}{2} \sum_{i+j=k} \alpha_{i,j} \beta_{i,j} N_i N_j - N_k \sum_{i=1}^{\infty} \alpha_{i,k} \beta_{i,k} N_i \quad \text{Equation 2.2}$$

where α is the stickiness coefficient, $\beta_{i,j}$ represents the cohesive collision frequency between particle i and j , and N is the floc number concentration.

There are three major mechanisms involved in estimating collision frequency: fluid shear, differential settling, and Brownian motion (Equations 2.4 to 2.6). Because only particles larger than $2 \mu\text{m}$ are included in this model, Brownian motion is ignored.

Although organic carbon and number concentration have different units, both variables use the same flocculation rate at each time step for a given floc cluster. Further, a proper selection of the collision kernel is necessary for a particular simulated environment.

Differential settling and fluid shear may be modeled by either a rectilinear or a

curvilinear form of the collision algorithm (Han and Lawler, 1991). The rectilinear kernel assumes that any flocs within a radius of two flocs' centerline would be intercepted by the settling particle. The curvilinear kernel assumes only smaller particle hit the centerline of larger one is swept. In general, the predicted collision rate by curvilinear kernel and rectilinear kernel ratio varies from 10^0 to 10^{-5} (Lawler, 1993). In this study, the curvilinear kernel is applied to the flocculation model because this type kernel is a better approach to the field environment in that particles would move following the streamline (Jackson, 1998). To convert the kernel from rectilinear to curvilinear, we used the formula from Han and Lawler (1991) and Li and Logan (1997).

$$\beta_{i,j} = \beta_{shear_i,j} + \beta_{settling_i,j} + \beta_{Brownian_i,j} \quad \text{Equation 2.3}$$

$$\beta_{rectilinear_shear} = \frac{G}{6}(D_i + D_j)^3 \quad \text{Equation 2.4}$$

$$\beta_{rectilinear_settling} = \frac{\pi}{4}(D_i + D_j)^2 |W_{s,i} - W_{s,j}| \quad \text{Equation 2.5}$$

$$\beta_{Brownian} = \pi(Diff_{floc,i} + Diff_{floc,j})(D_i + D_j) \quad \text{Equation 2.6}$$

$$\beta_{curvilinear} = \left(\frac{(D_i - 2D_j)^2}{(D_i + D_j)^2} \right) \times \beta_{rectilinear} \quad \text{Equation 2.7}$$

where β_{shear} , $\beta_{settling}$, $\beta_{Brownian}$ are the collision frequencies from shear stress, differential settling, and Brownian motion, G is shear stress gradient, D_i or D_j is the floc diameter, $Diff$ is the ideal floc diffusion coefficient, and $W_{s,i}$ is the floc settling velocity

Disaggregation is another important process. In this study, the LISST temporal volume concentration distribution profiles represent net aggregation in the STORM tank experiments, which are described in a later section (Schneider *et al.*, 2007). The shear

stresses in the STORM experiments were likely not sufficient to disaggregate particles. Therefore, this study does not explicitly simulate gross disaggregation but focus instead on net aggregation. We assume that the disaggregation process would only slow the flocculation process in the low shear stress environment. In practice, the stickiness coefficient value could be decreased to model disaggregation.

In this study, an essentially unlimited sediment supply has been assumed and the top layer sediment solid mass is assumed to be the same as the steady state total suspended solid mass in the water column in the STORM experiment. The mass erosion equation was based on the concept from Sanford and Maa (2001):

$$E_{b,i} = \sigma \left(\frac{\tau_b}{\tau_c} - 1 \right) \quad \text{Equation 2.8}$$

where $E_{b,i}$ is the pure sediment resuspension flux ($\text{g m}^{-2} \text{sec}^{-1}$), σ is a calibration coefficient, τ_b is bottom shear stress, and τ_c is critical shear stress.

There are many equations to estimate the floc settling velocity. In this study, the settling velocity equation was adapted from Winterwerp's model (1998), because fractal geometry was used to describe the flocs mass, volume, and porosity:

$$W_s = \frac{a}{18b} \frac{(\rho_{floc,dry} - \rho_w)}{\mu} D_p^{3-nf} \frac{D_f^{nf-1}}{1 + 0.15 \text{Re}^{0.687}} \frac{1}{1 + 2.5\phi} \quad \text{Equation 29}$$

where W_s is the floc settling velocity, a and b are shape parameters, $\rho_{floc,dry}$ and ρ_w are the solid floc and water densities, nf is the fractal dimension, Re is the Reynolds number, D_f is the floc diameter, D_p is the diameter of primary particle, and ϕ is the floc porosity.

2.3.2.2 Special Methods in This Study

Simulating the formation and movement of a wide range of mixed heterogeneous particles is the major task in this model. The model requires that both dry mass and dry

density are conserved, and is based on the multi cluster flocculation model from Yao *et al.* (1971) and Lick and Lick (1988). In this model, the term volume means the conserved spherical equivalent volume. This model calculates spherical equivalent volume and dry mass concentrations with the flocculation equations for each size floc cluster simultaneously:

$$V_{dry,i} = N_i \times V_{spherical_equivalent_cluster,i} \quad \text{Equation 2.10}$$

$$C_{dry,i} = V_{dry,i} \times \rho_{dry,i} \times (1 - \phi_i) \quad \text{Equation 2.11}$$

where V_{dry} is the solid volume concentration ($m^3 m^{-3}$), N_i is the floc number concentration, $C_{dry,i}$ is the solid floc mass concentration ($g m^{-3}$), ρ is the solid floc density ($g m^{-3}$), and the ϕ is the floc porosity.

The fraction organic carbon (f_{OC}), which is calculated as the ratio of organic carbon over total “dry” mass for each particle size at each time step, ranges between 0 and 0.5 and equals half of the fraction organic matter (f_{OM}). Two equations estimate the mass in this model. The first estimates the dry mass concentration which is converted from the number concentration and dry floc density of each floc cluster. The second estimates the organic carbon mass resulting from mass transport mechanisms and uses mass concentration as a unit. This study assumes two basic types of particles; clay ($f_{OC}=0$, density = $2.65 g/cm^3$) and biotic-substrate ($f_{OC}=0.5$, density = $1.05 g/cm^3$).

Floc porosity is assumed to be a constant in each floc cluster, and calculated from the fractal factor, the smallest particle size, and floc diameter with the compaction concept (Kranenburg, 1994).

$$\phi_f = [1 - (\frac{D_f}{D_p})^{n_f} (\frac{D_p}{D_f})^3] \quad \text{Equation 2.12}$$

where the ϕ is the floc porosity, D_f is the floc diameter, D_p is the diameter of primary particle, and nf is the fractal factor.

Khelifa and Hill (2006) suggested that the stickiness coefficient and the dry floc density should have a linear relationship with f_{OC} and this concept has been applied to this model. We ignored the possibility that the stickiness coefficient for the biotic substrate could lessen during phytoplankton blooms. Similarly, the bulk floc density has a linear relationship with porosity and dry floc density and varies temporally. All variables are calculated by:

$$\rho_{dry,i} = 2 \times f_{OC,i} \times \rho_{biotic} + (1 - 2 \times f_{OC,i}) \times \rho_{clay} \quad \text{Equation 2.13}$$

$$\rho_{bulk,i} = \rho_{dry,i} \times (1 - \phi_i) + \rho_{water} \times \phi_i \quad \text{Equation 2.14}$$

$$\alpha_i = 2 \times f_{OC,i} \times \alpha_{biotic} + (1 - 2 \times f_{OC,i}) \times \alpha_{clay} \quad \text{Equation 2.15}$$

where $\rho_{dry,i}$ is dry floc density (g/m^3), $\rho_{bulk,i}$ is bulk floc density (g/m^3), α_i is floc stickiness coefficient for floc at size i , f_{OC} is fraction of organic carbon for floc at size i , ρ_{biotic} , ρ_{clay} are organic carbon and inorganic carbon dry density (g/m^3), ρ_{water} is water density (g/m^3), ϕ_i is floc porosity for floc at size i , and α_{biotic} , α_{clay} are organic carbon and inorganic carbon stickiness coefficient respectively.

It is difficult to use traditional multi-floc-cluster flocculation models to simulate a wide range of floc sizes due to limitations on the floc cluster intervals and computer array sizes. In the STORM tank experiment, the flocs sizes measured varied between 2 to 1000 μm . However, with the traditional approach using fixed volume intervals, a model would require a matrix with more than 150 thousand elements to represent the particle properties when the primary particle size range from 2 μm ($8\mu\text{m}^3$ in volume) to 1000 μm ($10^9\mu\text{m}^3$ in volume). In addition, variables like number and organic carbon

concentrations, stickiness coefficient, collision efficiency, floc density, settling velocity, and f_{OC} are all calculated for each floc size at each time step, which is not practicable. To solve this problem, a fixed floc cluster diameter interval is used, meaning that floc cluster N 's diameter is N times longer than floc cluster one. The floc spherical equivalent volume is calculated based on floc diameter in each floc cluster. The volume and mass concentrations are varied temporally in each floc cluster, but the gross dry mass and volume are conserved. In this model, the particle sizes vary between 2 to 1000 μm with a 1 μm interval. Therefore, the model has one thousand particle floc clusters or variables to represent each individual floc property.

2.3.3 Calibration Data

Schneider *et al.* (2007) conducted Shear Turbulence Resuspension Mesocosms (STORM) tank experiments that mimic resuspension and settling of contaminated upper Hudson River sediment with realistic bottom shear stress and water column turbulence. The average instantaneous bottom shear stress was about 1 dyne cm^{-2} and the volume-weighted average water column turbulence intensity and energy dissipation rate were 0.55 cm s^{-1} and 0.0032 $\text{cm}^2 \text{s}^{-3}$, respectively. Hudson River sediment was added to a depth of 5 cm and allowed to consolidate for 10 days. This experiment included both erosion and free settling periods. During erosion, the mixing paddle continuously generated bottom shear stress for 53 hours to ensure suspended solids reached steady state. During the one hour free settling period, the paddle was turned off to allow suspended particles to settle. Particle size distribution, dissolved and particulate PCBs, TSS, DOC, particulate carbon, nitrogen, and chlorophyll *a* were measured throughout the resuspension and settling portions of each experiment (Schneider *et al.*, 2007).

Because observations were made at the mid-depth of the tank, the free-settling scenario assumes the water column depth to be 0.5 m. The erosion scenario assumes the water column depth is 1 m because of the tank is well-mixed. In the STORM experiment, the suspended particle size analysis was conducted using LISST-100C (Sequoia Scientific, Inc.). The LISST measures the particle size distribution between 2 and 500 μm , in 32 size bins in a log scale, providing the steady-state particle size distribution for the erosion period and the initial particle size distribution for the free settling period.

The sediment grain size analysis was conducted using the Beckman Coulter LS100 Laser Diffractometer. This instrument works on the same principles as the LISST and measures the size distribution of suspensions of nonsieved sediments 0.4 to 1000 μm in diameter using the laser diffraction technique. The size distribution of the homogenized Hudson River sediment was tri-modal with a large peak in the volume size distribution at 145 μm and two lesser peaks at 61 and 473 μm respectively. The volume median diameter of the sediment grains was 63 μm . This analysis result was applied to this model as an initial sediment particle size distribution (Figure 2.1).

2.3.4 Numerical Method

The system was solved on 1000 floc clusters for each variable. The above equations form the basis of our model, which was written in double precision Fortran 90 and used a fourth-order Runge-Kutta numerical algorithm to solve the set of nonlinear ordinary differential equations. The model was stable with a time step of one minute and a 53 simulation hours run was completed in 75 CPU minutes on a Sun Ultra-60 workstation. Both dry mass and dry volume were conserved during the two period simulations. At the end of the free settling period simulation, 7.25 g organic carbon and 411 $\mu\text{m}^3\text{m}^{-3}$ dry floc

volumes remained distributed among two reservoirs, water and sediment, demonstrating that there is no systematic numerical drift within the model.

2.3.5 Parameterization and Initial Conditions

Several parameter values and initial conditions were estimated from a combination of observations during the STORM experiment and literature reports, including the fractal factor, the primary particle character, the porosity, and initial particle volume and the f_{OC} size distribution and corresponding values for each individual floc cluster.

In this model, the fractal factor is the dominant parameter controlling particle porosity, and therefore the settling velocity. The sampled particles were predominately flocs (Schneider, 2007), so the fractal factor is less than 3.0. We used the highest estuarine fractal factor value suggested by Winterwerp (1998) of 2.3 because the eroded sediment was consolidated for 10 days prior to analysis. The influence of the fractal factor was evaluated in the sensitivity analysis described below. The primary particle is assumed to be a solid of 2 μm diameter which does not disaggregate and is the lowest LISST reading in the STORM experiment. The porosity of each floc cluster is estimated by Equation 2.7.

The model is initialized with a particle size distribution derived from the STORM experiments. To interpolate the measured size distribution, we used the strategy of Mikkelsen *et al.* (2005), in which both the total volume concentration and the shape of the particle size distribution are conserved. The same method was used to interpolate the sediment size distribution measurements to initialize the sediment particle size distribution in the model.

The LISST integrates all flocs $> 500 \mu\text{m}$ into one group, and in the STORM experiments Schneider *et al.* (2007) ignored all flocs larger than $250 \mu\text{m}$. However, Mikkelsen *et al.* (2005) used a digital camera to measure the larger floc size distribution from Hudson River sediment, and showed that the particle size distribution for the larger floc decreased with particle size. Present model results show the similar behavior for the larger particles. Therefore, the initial particle size distribution was extrapolated to $1000 \mu\text{m}$ iteratively based on steady state model results for the sediment layer during the erosion period simulation and suspended particles during the free settling period simulation.

The next step was to set the eroded particle size distribution. Published models estimate the erodable particle size as a function of particle density, size, critical shear stress, and bottom shear stress. However, the eroded particle size distribution was also impacted by aggregation-disaggregation along with the given bottom shear stress. To simplify the model simulation, the first measured suspended particle LISST profile was selected to represent the net eroded particle size distribution (Figure 2.1) and assumed no suspended particles in the water column at time zero. This distribution does not vary temporally in the model.

The f_{OC} distribution plays an important role in this model. Because there is no method to measure the f_{OC} in the individual sized particles, we have to make a reasonable assumption for this variable. Three types of distributions were tried (Figure 2.2): a uniform f_{OC} distribution across all sizes, f_{OC} concentrated in the small flocs, and a size-variable f_{OC} distribution based on the predicted steady state f_{OC} distribution resulting from a two type particle erosion scenario which described in detail later. Because the model is very sensitive to the f_{OC} distribution, each assumption causes a different trend in

the steady state. Compared the model simulation results with STORM experiment measurement, the first two f_{OC} distributions could only match the particle size distribution during the early stage of resuspension and each missed the later trend. The third scenario gave the best overall agreement with measurements at steady state. However, the third f_{OC} distribution could not be directly applied to the model because of the conservation of TSS, TVC, and f_{OC} . To solve this problem, the individual floc cluster f_{OC} value was back calculated from equation 2.5 to 2.9, because TSS, total organic carbon concentration, and total volume concentration were measured values, and fractal factor value and individual volume concentration were assigned.

To ensure that the model eroded a realistic amount of mass under the given shear stress, we compare model results with Upper Chesapeake Bay data (Sanford, unpublished). In Figure 2.3, we compared the relationship between eroded mass and eroded flux from STORM experiment and field data. This figure shows that the relationship between eroded flux and eroded mass under a given shear stress in STORM tank experiment was similar to field observations.

2.3.6 Sensitivity Analysis

The flocculation model was first evaluated by exploring the sensitivity of key outputs (the gross floc settling velocity for TSS ($W_{s_{TSS}}$), residence time, steady state TSS concentration, and particle size distribution trends) to several parameters. The model was initialized with 1 to 3 discrete sizes of flocs under the same conditions as the STORM experiments, including the experiment duration, eroded shear stress, and tank size during the both the free settling and erosion periods. In this chapter, the residence time was calculated as the concentration weight average for size of flocs.

$$Residence_Time = \frac{\sum [C_{dry,i} \times \left(\frac{V_{tank}}{Area \times W_s} \right)]}{\sum C_{dry,i}} \quad \text{Equation 2.16}$$

where V_{tank} is the experimental tank volume (m^3), area is the sediment surface area (m^2), $C_{dry,i}$ is the solid floc mass concentration ($g\ m^{-3}$), and $W_{s,i}$ is the floc setting velocity.

The first sensitivity test examines the impact of the fractal factor on the gross settling velocity for TSS ($W_{s\ TSS}$) during a one hour free settling period. The gross settling velocity for TSS during the free settling period was estimated as a first order equation which was only used in this section:

$$\frac{\partial TSS}{\partial t} = - \frac{W_{s,TSS}}{x} TSS \quad \text{Equation 2.17}$$

In this test, the model first simulated 0.1 g of clay packaged into 20 μm flocs settling in a one meter tank for 1 hour without flocculation. Three fractal factor values (1.0, 2.0, and 3.0) were evaluated. When the fractal factor equals 3.0 the 20 μm particles are solids, and particle settling is predicted by Stokes' law. When the fractal factor is less than three the porosity increases and the floc bulk density decreases, resulting in a smaller settling velocity. This analysis demonstrates that the fractal factor strongly controls particle settling, as $W_{s\ TSS}$ decreased from 19 m/day to 0.1 m/day as fractal factor decreased from three to one.

The second sensitivity test examined the impact of flocculation on $W_{s\ TSS}$. In this test, the model simulated 0.1 g/m^3 of clay particles in flocs with an initial 20 μm diameter and a fractal factor of 2.0 which settled through one meter for 1 hour. The stickiness coefficient (α) varied between zero (no flocculation) and 1 (every collision results in

flocculation). Under these conditions, $W_{s\ TSS}$ is 30% higher when $\alpha=1$ than without flocculation due to the formation of larger flocs.

Previous particle transport models use D_{50} to represent the entire particle spectrum and use this value to estimate the $W_{s\ TSS}$ (Winterwerp, 1996). However, as discussed before, a size distributed suspension may increase the flocculation rate due to differential settling-induced collisions. Thereafter, the gross $W_{s\ TSS}$ for cohesive particles may have different value between multi-size-particle and D_{50} runs. The third test compares the impact of flocculation on $W_{s\ TSS}$ between D_{50} and the multi-size-particles. The model again simulated 0.1 g/m^3 clays with a D_{50} of $20\ \mu\text{m}$ diameter settling through one meter for 1 hour with a fractal factor of two. Three sizes of floc clusters, 10, 20 and $30\ \mu\text{m}$, having the same average volume concentration size value as D_{50} run, were selected to compare with the D_{50} simulation results. The stickiness coefficient varied between one and zero to control the extent of flocculation. Without flocculation, there is less than 2% in difference for the $W_{s\ TSS}$ between multi- and single size flocs scenarios. This result implicates that D_{50} is a workable tool to estimate $W_{s\ TSS}$ for a broad size distribution of non-cohesive particles. After including flocculation, the multi-size scenario increases $W_{s\ TSS}$ around 55% compared to the single size scenario. Flocculation occurs more rapidly in the multi-size scenario since there is a higher collision probability. This result implies that assuming a single particle of size D_{50} maybe underestimate the $W_{s\ TSS}$ for the broad size distribution of cohesive flocs.

The fourth test examines the impact of particle characteristics on $W_{s\ TSS}$. Separate simulations with $20\ \mu\text{m}$ diameter clay (density = 2.65 g/cm^3) and biotic (density = 1.05 g/cm^3) particles, each with and without flocculation, are compared. For both particle

types $W_{s\ TSS}$ increases with flocculation due to the formation of larger, faster settling flocs. Without flocculation, $W_{s\ TSS}$ calculated for clays was 2.5 times faster than for biotic substrates because of differences in the floc density. With flocculation, the $W_{s\ TSS}$ for clay was only 1.2 times larger than that of the biotic substrates during the one hour simulation. Flocculation impacts the settling velocity of the biotic particles more than clay. Unlike clay, most low density biotic flocs remained in the water column in both runs, having a better opportunity to coagulate with other flocs during the one hour simulation period and one meter water depth.

Clay has been widely used to remove algae from wastewater treatment plants and from open water (Rengasamy *et al.*, 1996). To ensure that the flocculation model could work with erosion and sedimentation processes, the fifth test simulates flocculation's impact on a simple clay and diatom mixture. The model used two sizes of particles that were constantly eroded by bottom shear stress in a one meter deep tank: a 20 μm ($f_{OC} = 0$) clay floc and a 22 μm ($f_{OC} = 0.5$) biotic floc. The critical shear stress and bottom shear stress were the same as in the STORM experiments. Three tests were done in which the model constantly eroded clay only, diatoms only, and both together from the sediments at rates of 2.4 ($f_{OC} = 0$), 1.05 ($f_{OC} = 0.5$) and 3.45 ($f_{OC} = 0.14$) $\mu\text{g m}^{-3}\text{s}^{-1}$, respectively. The model was run with a one minute time interval for a 1000 hour to ensure that the model reached steady state. Further, three runs were done in each test, and all runs had the same parameters and settings, except the stickiness coefficients, which varied between 0, 0.5 and 1.0 to represent the strength of flocculation

Figure 2.4 shows the steady state total suspended solids concentration, calculated residence time, D_{50} , and simulated particle size volume concentration distribution both at

steady state and in transitory periods for all runs. In general, the steady state TSS and water column residence times decrease with stronger flocculation due to forming larger flocs. Before reaching steady state, we observed that the model run with more flocculation had a higher volume concentration of the larger flocs in all tests because those runs had a faster flocculation rate to form larger flocs. Comparing the three tests, the residence time value corresponded to floc densities. In this study, the residence time for the diatom-only test decreased almost 53% and 41% with half and full flocculation effects, respectively. When two types of particles started to coagulate, the residence time decreased slightly faster than in the diatom test. The difference in residence time could be more significant if we eroded a more complex floc group to alter the flocculation efficiency with increasing the differential settling and collision number concentration.

After including the flocculation effect, the particle size volume concentration distribution showed a similar shape but different quantities for both runs at steady state in each eroded test. In addition, we observed the periodic volume concentration peak alone with increasing floc size, which has the similar trend as observed in the field (Sanford, personal communication). This is because the model started with one size of particles, and the assumption of volume conservation-floc cluster size relationship. Compared to the steady state clay-only results, the diatom-only test had larger flocs, which because of their lighter density settled more slowly than clay.

Without flocculation, the steady-state TSS concentration from the clay plus diatom mixed-erosion scenario was equal to the sum of TSS concentrations from the clay only and diatom only runs. This result suggested that the steady state TSS concentration from a group of non-cohesive particles could be estimated from the summation of the

individual particle size run result. However, after including flocculation both TSS and volume concentrations from the mixed-eroded-scenario were less than the sum of TSS and volume concentrations from the clay and diatom runs at steady state. These results showed the impact of flocculation on the kinetic variation in number concentrations, floc density and settling velocity.

2.3.7 Calibration with STORM Experimental Data

To understand the effect of shear stress on TSS and D_{50} with flocculation, we applied different shear stresses in this model. Model initial conditions and environment parameters were the same as STORM tank erosion calibration scenario. Shear stress was varied from 1 to 4 dyne cm^{-2} in 1 dyne cm^{-2} increments. Each shear stress was applied constantly for 53 simulation hours to ensure that the model reached steady state. As before, disaggregation was not included in this simulation. Figure 2.5 shows how the model simulated D_{50} , and the TSS and organic carbon concentrations varied with shear stress. In general, at steady state the D_{50} , TSS, and organic carbon concentrations increased with higher shear stress because more mass was eroded. Further, higher shear stress and higher concentrations also increased the flocculation efficiency to form larger flocs. Thereafter, the time to reach steady state was shorter with higher shear stress. D_{50} had several peaks when shear stress was changed in the model. The sharp D_{50} increase was a combination of steady state particle size distribution under previous shear stress and additional eroded particle size distribution under new shear stress.

The flocculation model was calibrated with STORM experiment results by adjusting the stickiness coefficient and erosion coefficient. The final calibrated parameter values are listed in Table 2.1. Figure 2.6a-f showed measured and model calculated organic

carbon, total suspended solids, and total volume concentrations during the erosion and free settling periods.

2.4 *Example Applications*

2.4.1 Comparison of Organic Carbon Settling Rates Predicted by the Flocculation Model with those from a Traditional Particle Transport Model

A traditional particle transport model does not include the formation and fractal nature of flocs and often fixes the organic carbon sedimentation rate at a value between those of inorganic solids (e.g., clays) and plankton (Di Toro, 2002). To compare the TSS and organic carbon settling velocities with and without flocculation, four scenarios were designed. Scenario one and two ignored flocculation, assumed flocs had constant densities as 1.05 and 1.35 g/cm³, respectively, and used settling velocities calculated using the Stokes' law equation. The third scenario included flocculation, allowing floc density to vary temporally with floc composition (but not with floc size). The last scenario is the same as scenario three except that the setting velocity is calculated using the equation of Winterwerp (1998), which is the strategy in the model described in this paper.

Figure 2.7 compares the measured and modeled TSS and organic carbon concentrations among the four scenarios. The diatom-based scenario one results underestimated the TSS and organic carbon settling velocities due to low particle densities. This under-prediction is similar to those seen in other organic carbon particle transport models like the DELPCB model (DRBC, 2003). These results suggest that the impact of

inorganic solids on organic matter settling cannot be ignored in estuaries and other algal-rich waters.

The other interesting finding is the result from scenario two, in which the particle density was estimated from field observations of settling, which is the approach commonly used in many particle transport models. The result suggested that in organic rich environments, modifying the particle density to improve the organic carbon settling velocity does not work well and under-estimates the organic carbon residence time because flocculation continues to alter the particle size and gross bulk density.

Scenario three and four started with the same initial condition, giving us an initial look at the impact of equations 2.4 and 2.11 on the fate of flocs. Scenario three predicted that both TSS and organic carbon concentrations decreased much faster than observations. But scenario four had a good agreement with observations. These results agreed with the suggestion from Winterwerp to use a fractal geometry-based equation to estimate the settling velocity for flocs.

2.4.2 The Impact of Eroded f_{OC} Distributions on the Behavior of Organic Carbon

As discussed before, the model performance is sensitive to the f_{OC} distribution. Because there is no method to measure the f_{OC} of the individual sized particles, we have to make a reasonable assumption for this variable. Three types of distributions were considered in this study: a uniform f_{OC} distribution across all sizes, which has the same effect as assuming all particles have the same composition, f_{OC} concentrated in the small flocs, which has been adapted by the other models (Gong and DePinto, 1998), and a size-variable f_{OC} which we believed to be closest to the actual distribution (Figure 2.2). Each

distribution has a unique impact on the gross particle residence time, settling velocity, and steady state concentration, after including the flocculation, erosion, and settling. In this section, we use this model to explore the impact of the f_{OC} distribution and flocculation on the fate of eroded sediment carbon.

We designed a series of model runs to estimate the steady state concentrations and calculated the residence time for bulk organic carbon and TSS with varying stickiness coefficients and three types of f_{OC} distributions. All runs had the same gross organic carbon and inorganic solids erosion fluxes, and erosion was based on the same particle size distribution and fractal factor. To simplify the model settings, here clay and organic matter were assumed to have the same stickiness coefficients. The stickiness coefficient varied from zero, representing no flocculation, to 1, representing full flocculation, in 0.2 increments. Each set was tested with the three f_{OC} distributions. There were total 18 runs in this case, and each run was executed for 200 simulated hours to reach steady state.

The steady state D_{50} increases and both the organic carbon concentrations and the time to reach steady state decrease when the model includes flocculation (Table 2.2). However with flocculation the variations in steady state TSS and organic carbon concentrations were very limited among different stickiness coefficient runs for each f_{OC} distribution because at steady state all runs predicted similar organic carbon and TSS volume size distributions. The only difference was the time to reach steady state. Increasing flocculation (higher stickiness coefficient) reduced the time to reach steady state. Compared with the other f_{OC} distributions, the varied distribution runs reached steady state more quickly and had the least difference between the residence time and steady state organic carbon concentrations with and without the flocculation because

these runs were initialized with a steady state f_{OC} . These runs did not reach steady state immediately because the initial particle size distribution was not the same as the steady state particle size distribution.

These trends agree with the other reports that when particles start to form flocs, the gross settling velocity increases as larger particles form. Schneider (2005) used a first order equation to estimate the TSS based floc settling velocity with a mean of 0.12 cm s^{-1} for the upper Hudson River. Fugate and Friedrichs (2002) reported a TSS-based settling velocity at the Cherrystone site in the Chesapeake Bay by an acoustic doppler ranging from 0.07 to 0.13 cm s^{-1} . Sanford *et al.* (2005) used a video camera technique to measure upper Chesapeake Bay floc settling velocities ranging from 0.02 to 0.50 cm s^{-1} , with a mean of 0.15 cm s^{-1} . The model calculated TSS-based settling velocities ranged between 0.09 to 0.16 cm s^{-1} , and organic carbon based settling velocities ranged between 0.03 to 0.18 cm s^{-1} , which was within the report range of values.

When the eroded organic carbon is initially concentrated in small flocs, flocculation more significantly impacts the organic carbon steady state concentration. In addition, the organic carbon residence time and steady state concentration with full flocculation ($\alpha = 1$) were one third of the results without flocculation because without flocculation all organic carbon was in small, low density flocs with slower settling velocities. As they coagulated, the organic matter aggregated into larger, denser particles, resulting in faster settling and lower steady state organic carbon concentrations. On the contrary, the runs under size-variable f_{OC} distribution are least impacted by flocculation. The organic carbon residence time and steady state concentration with full flocculation are only 5% different than without flocculation.

Furthermore, because f_{OC} varies temporally among different floc sizes in this model, bulk TSS and organic carbon may settle with different velocities depending on the initial f_{OC} distributions. When f_{OC} is evenly distributed among all sizes of flocs, all flocs have the same density and there is no difference between organic carbon and inorganic solids steady state concentration distributions. However, when f_{OC} varied with particle size distribution, density varies with floc size, altering the size-dependent settling velocity distribution. In the small flocs dominated distribution runs, the TSS settling velocity was almost three times faster than organic carbon settling velocity. For the varied distribution scenario, the organic carbon settled 13% faster than TSS.

The floc dry density, which controls the settling velocity, depends on the f_{OC} of each floc size cluster. In this model, within the same floc size cluster organic and inorganic substrates have the same settling velocity. As a result, under the uniform f_{OC} distribution both gross TSS and organic carbon are predicted to have the same settling velocity. On the contrary, TSS and organic carbon should have the different settling velocity when the model is initialized with an uneven f_{OC} distribution. When organic carbon is enriched in the larger flocs, faster settling velocities lower steady state organic carbon concentrations compared to the same amount of organic substrates beginning in smaller flocs. Therefore, the organic carbon in the small f_{OC} dominated distribution had the longest time to reach steady state compared to the other two f_{OC} distributions.

These results suggest that the organic carbon size distribution in the sediment is an important factor controlling the residence time, gross settling velocity, and steady state concentration of solids and organic matter. Most particle transport models assume that f_{OC} is constant and evenly distributed across the entire range of particle sizes, and that

organic carbon has the same settling behavior as TSS. In this study, we found these models might either under or over estimate the fate of organic carbon when f_{OC} is not evenly distributed, with results depending on the nature of the f_{OC} distribution in the sediment.

2.4 *Summary*

The behavior of organic particles, which strongly influences the fate of organic contaminants in natural waters, is controlled by settling and erosion fluxes. However, there is no universal method to estimate the behavior of organic matter when it forms flocs with itself or with inorganic solids. In this study, a flocculation model that simulates the flocculation of both organic carbon and inorganic solids ranging in diameter from 2 to 1000 μm has been developed. This model simultaneously calculates the temporally varying organic carbon content, density, flocculation coefficient, and settling velocity for each size of particles under different scenarios. The model was calibrated using the Shear Turbulence Resuspension Mesocosms (STORM) tank experiments that mimic resuspension and settling of contaminated Hudson River sediment with realistic bottom shear stress. This model was effective in predicting the temporal variability in the behavior of wide spectrum heterogeneous flocs in the mesocosm.

The model predicted that the water column residence time of TSS and organic carbon decrease and the median size of particles (D_{50}) is less with increased flocculation and fractal factor. Several settling velocity strategies were tested based on the same initial condition for a heterogeneous source particle population. A fractal geometry-based settling velocity equation agreed best with the STORM observations. The results also suggested that using a modified but temporally-invariant particle density over-estimates

the organic carbon settling velocity in organic carbon rich environments due to the influence of flocculation on settling velocities. Including flocculation in the multi-size scenario increases $W_{s\ TSS}$ 55% compared to the single size scenario. This result suggests that modeling settling with a single particle of size D_{50} may underestimate the $W_{s\ TSS}$ for the broad size distribution of cohesive flocs.

Organic carbon size distribution is another important factor influencing the water column residence time, gross settling velocity, and steady state concentration of solids and organic matter. In this study, we found that the relationship among individual floc cluster density, gross organic carbon, and TSS settling velocity were impacted by the f_{OC} distribution. They have the same settling velocity and individual floc cluster density only when f_{OC} does not vary with particle size and is evenly distributed over all sizes of particles. The settling velocity of organic carbon plays an important role in predicting the fate of particle-reactive chemicals in the water column. When a wide spectrum of heterogeneous flocs is present, it is necessary to apply a multi-size-floc strategy to simulate gross particle behaviors. The next step of this study will apply this flocculation model to simulate the fate of organic contaminants in an urban estuary.

2.5 *Figure Captions*

Figure 2.1: The initial volume size distribution for sediment and eroded flux during the erosion period and of suspended particles at the beginning of free settling period.

Figure 2.2: Three f_{OC} distribution trends: small, uniform, and a size-variable. All trends are estimated by the same gross f_{OC} (0.115), gross TSS (43.5g/m³), and gross TVC (190uL/L) as in the STORM experiment

Figure 2.3. Model predicted TSS, residence time, D_{50} , and simulated particle size volume concentration distribution both at steady state for clay eroded only ($f_{OC}=0$), biotic-substrates eroded only ($f_{OC}=0.5$), and clay-biotic-substrate-co-eroded ($f_{OC}=0.14$) scenarios. In each scenario, the model was tested with three different stickiness coefficients (0, 0.5, and 1).

Figure 2.4: Model predicted TSS, organic carbon concentrations, and D_{50} variation with different shear stress values. Shear stress varied from 1 to 4 dyne cm⁻². Each shear stress was applied for 53 simulation hours. Fractal factor is 2.3 with TSS equal to 43.5 g m⁻³, TVC equal to 191 μ l/L, and gross f_{OC} equal to 0.115.

Figure 2.5: Comparison of the experimental eroded flux-eroded mass relationship with Upper Chesapeake Bay field measurements (Sanford, unpublished)

Figure 2.6 a and b: Comparison of model-predicted and measured TSS and organic carbon concentrations during the STORM free settling period. Fractal factor is 2.3 with TSS equals to 63 g/m³, TVC equal to 410 μ l/L, and gross f_{OC} equal to 0.115

Figure 2.6 c, d, e, and f: Comparison of model predicted and measured TSS, TVC, organic carbon concentrations, and D_{50} during the STORM resuspension period. Fractal factor is 2.3 with TSS equals to 43.5 g/m³, with TVC equal to 191 μ l/L, and gross f_{OC} equal to 0.115.

Figure 2.7: Comparison of measured and model-predicted TSS and organic carbon concentrations among four scenarios. All runs started with the same initial conditions for an one hour duration at one meter water depth: Fractal factor is 2.3 with TSS equal to 63 g/m³, TVC equal to 410 μ l/L, and gross f_{OC} equal to 0.115.

Figure 2.1

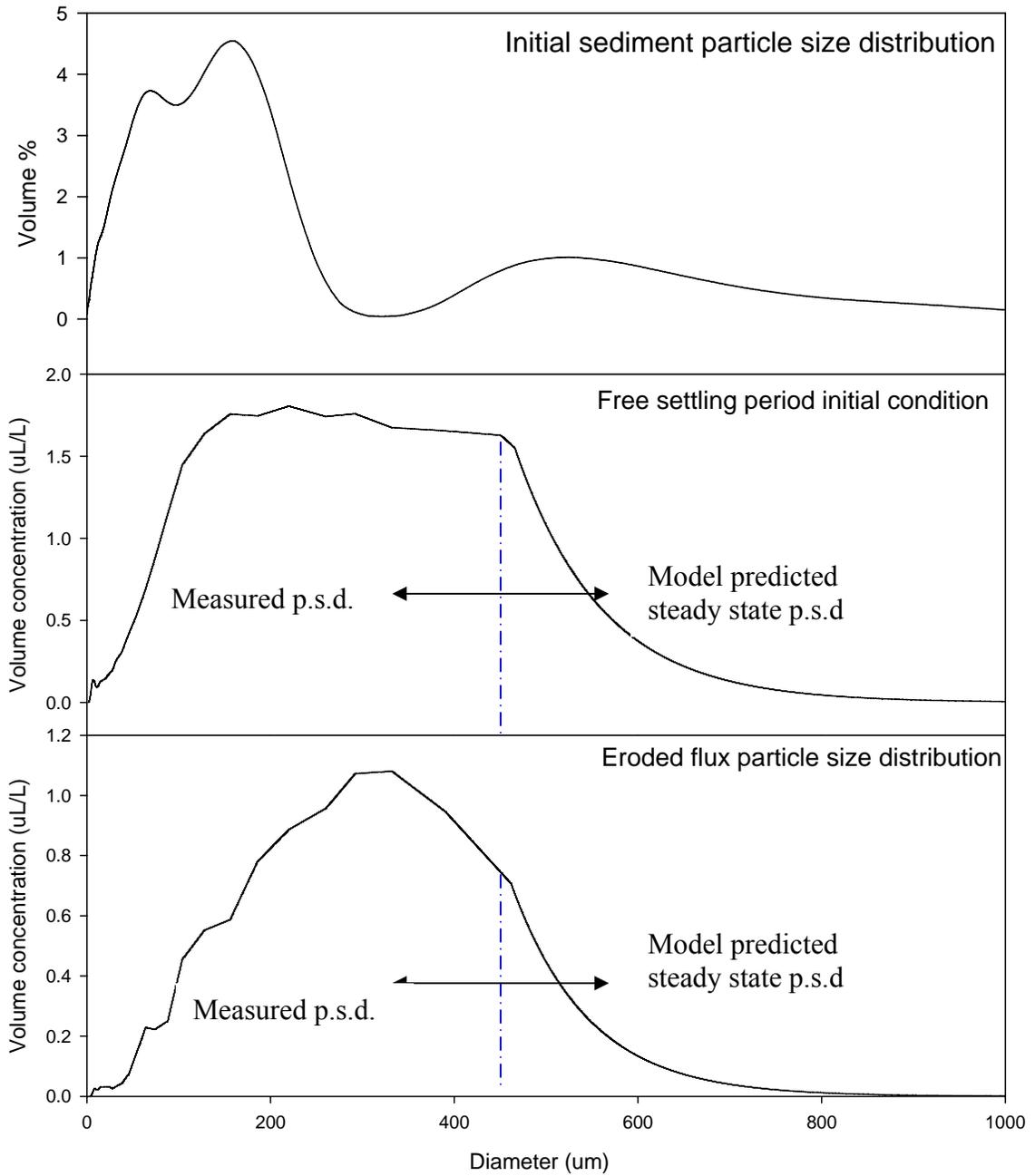


Figure 2.2

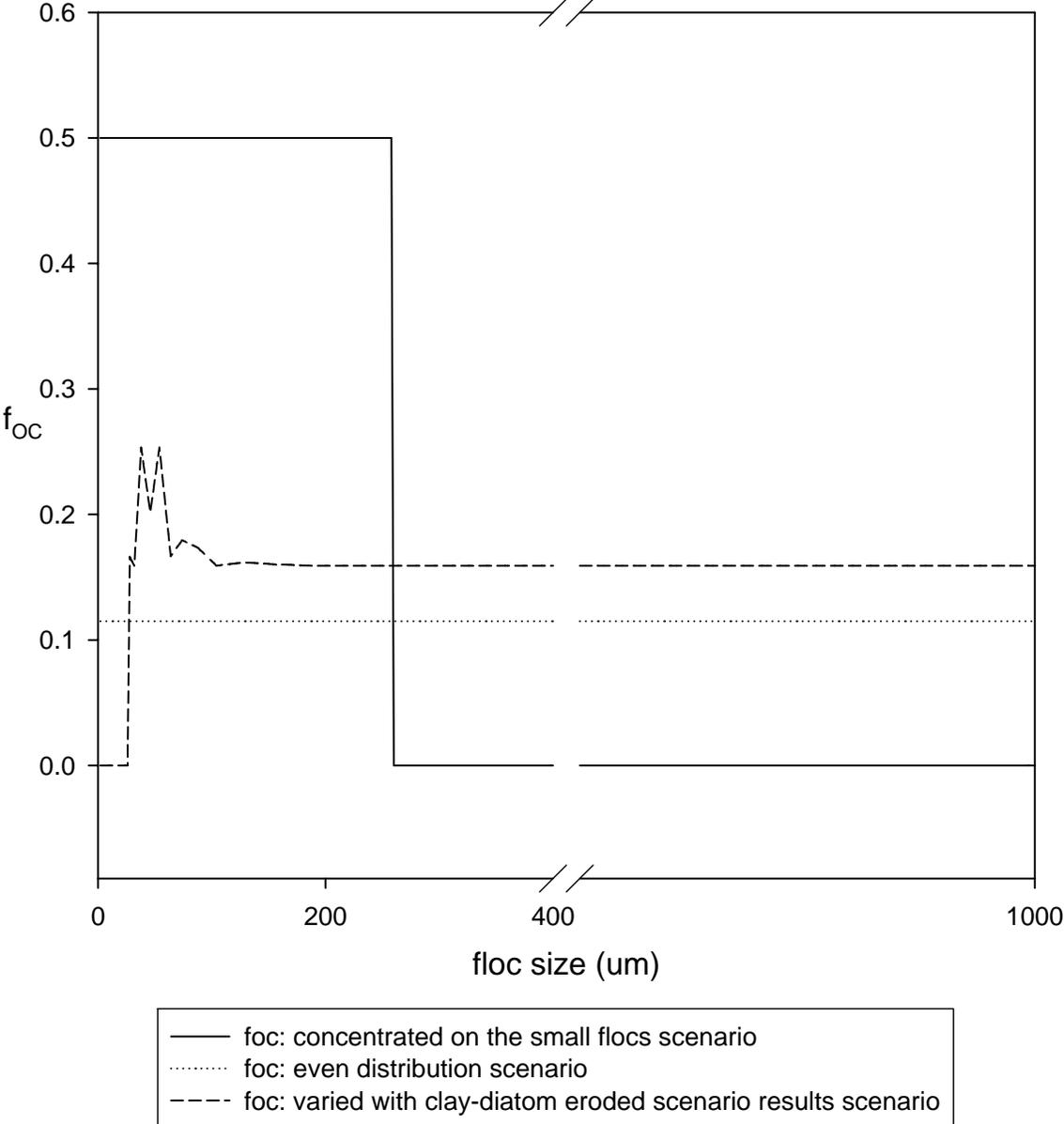


Figure 2.3a

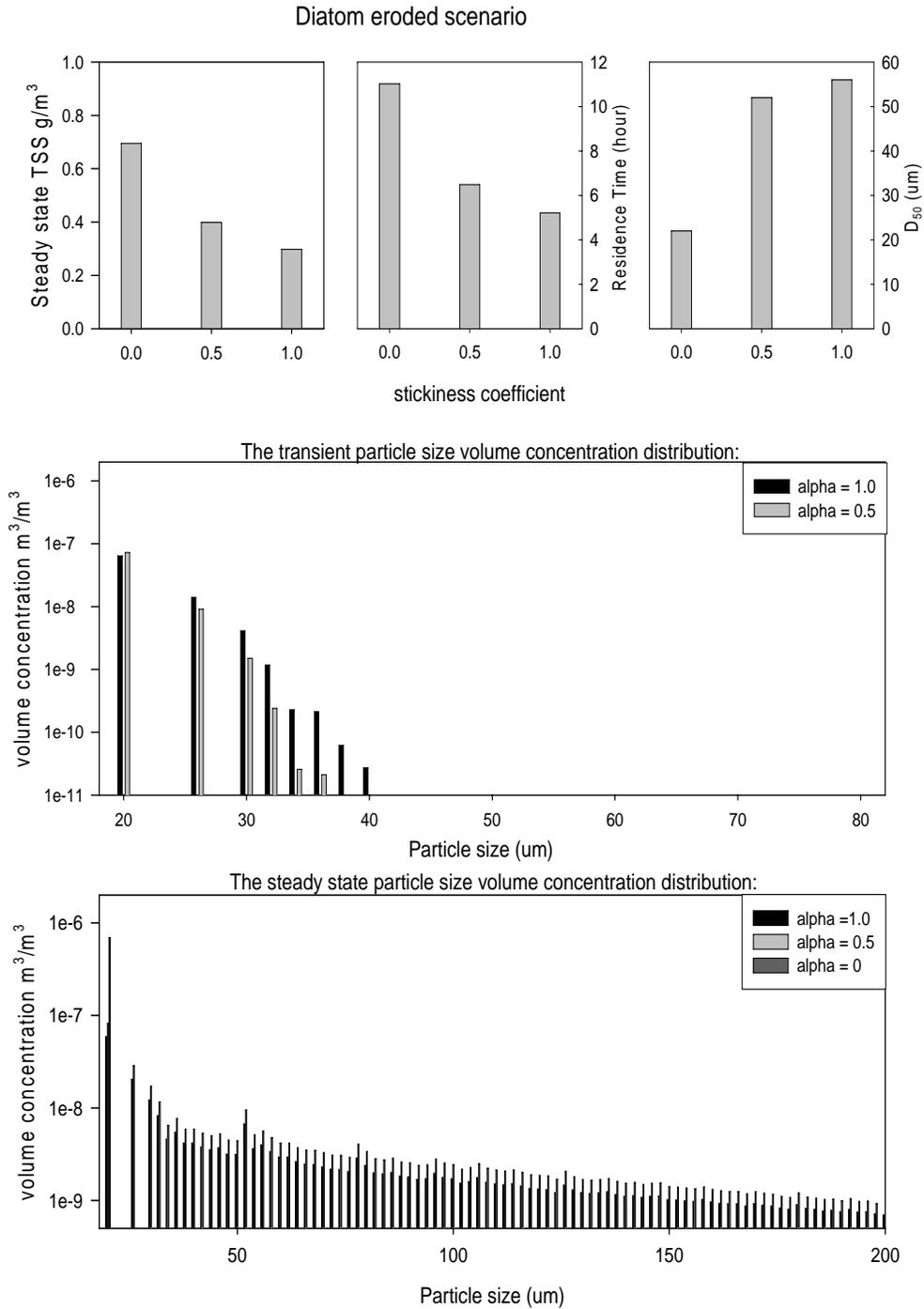
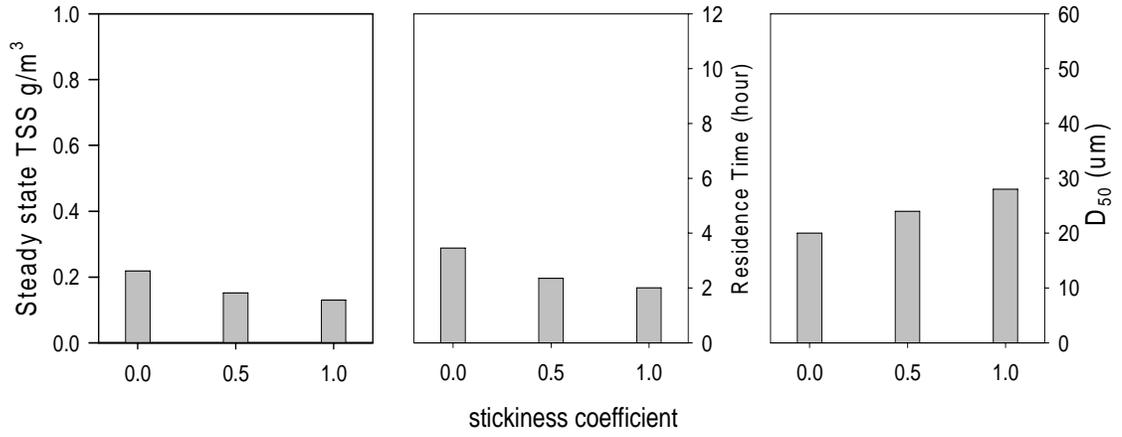
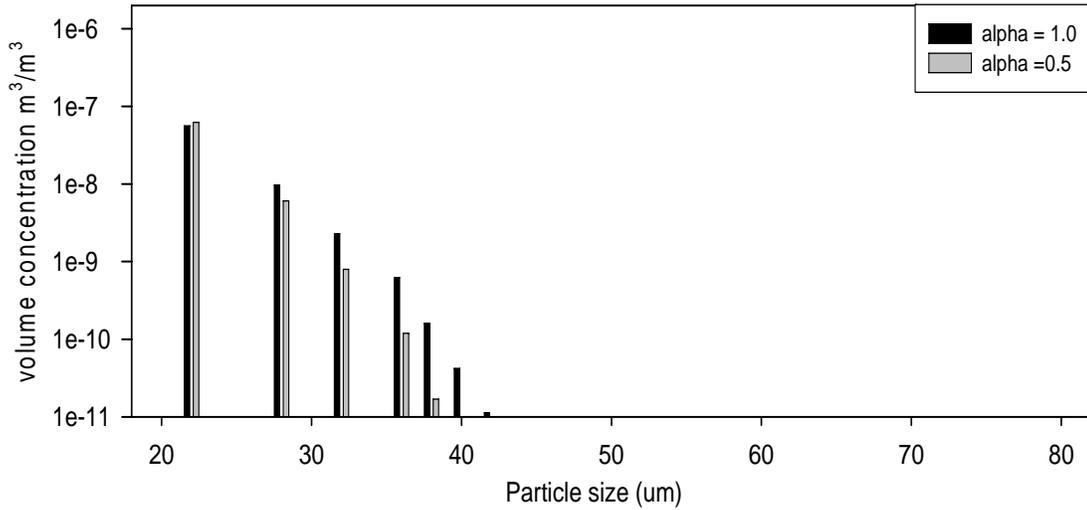


Figure 2.3 b

Clay eroded scenario



The transient particle size volume concentration distribution:



The steady state particle size volume concentration distribution:

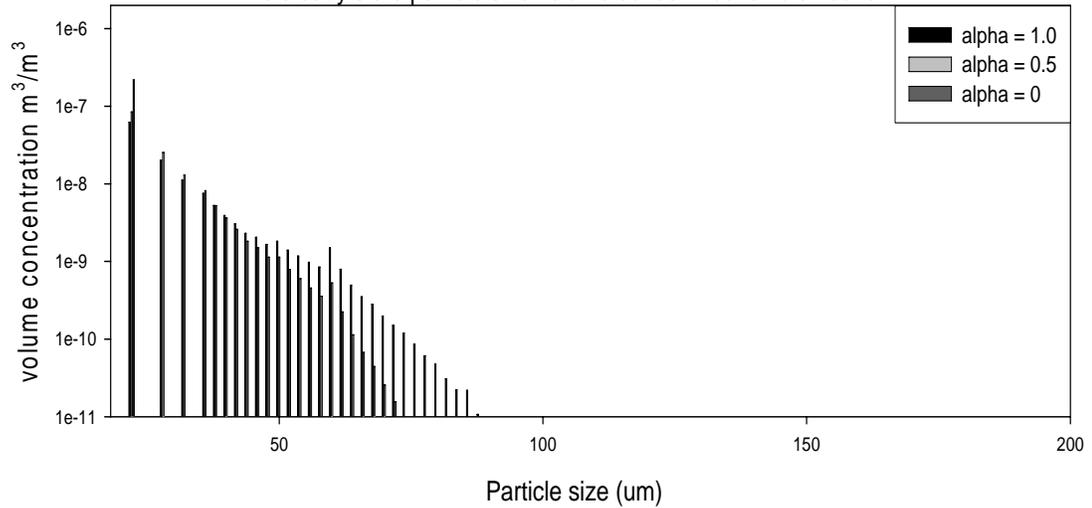
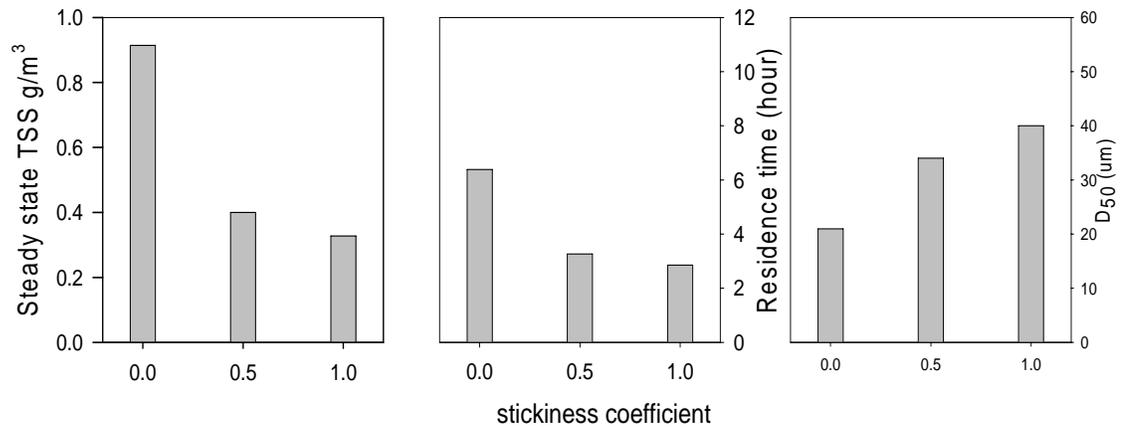
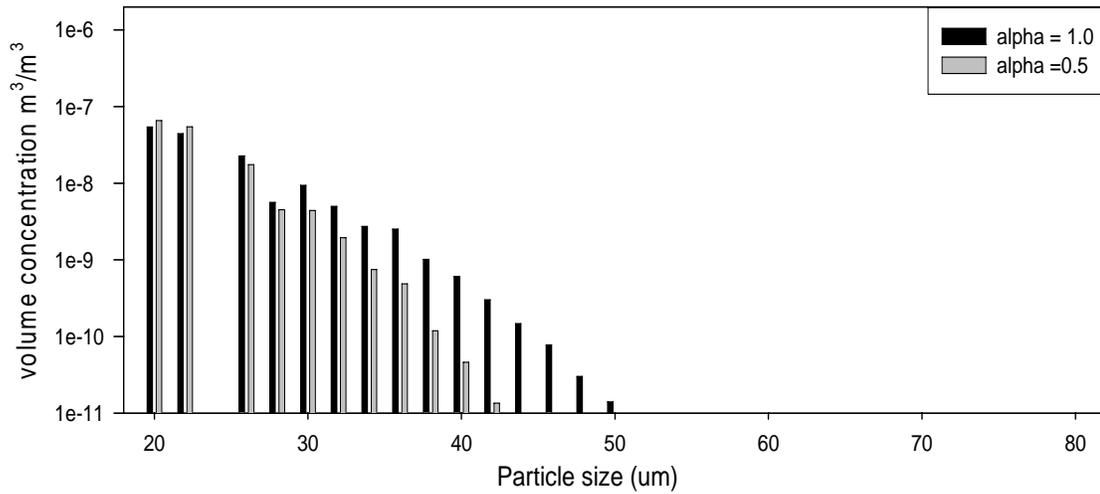


Figure 2.3c

Clay and diatom mixed eroded scenario



The transient particle size volume concentration distribution:



The steady state particle size volume concentration distribution:

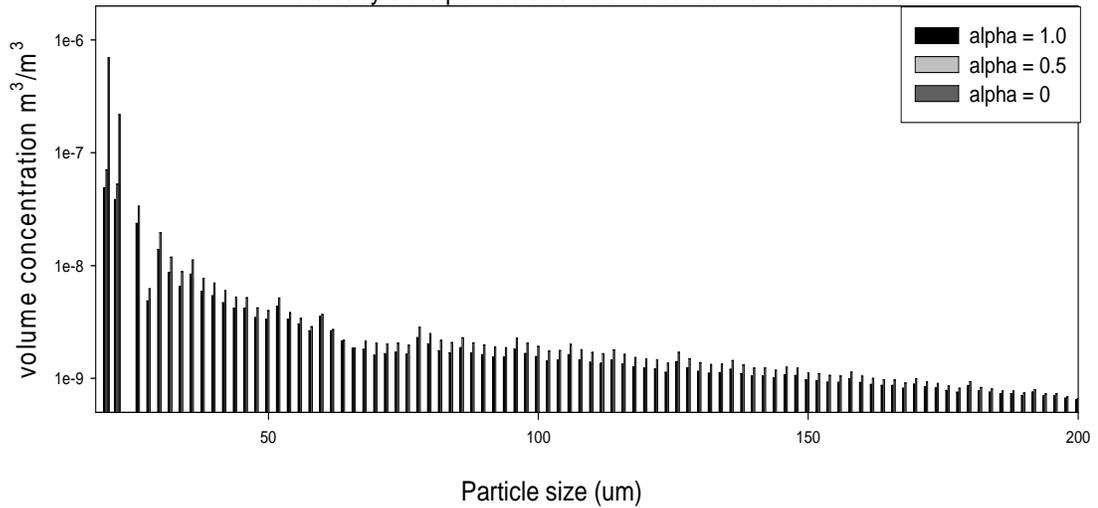


Figure 2.4:

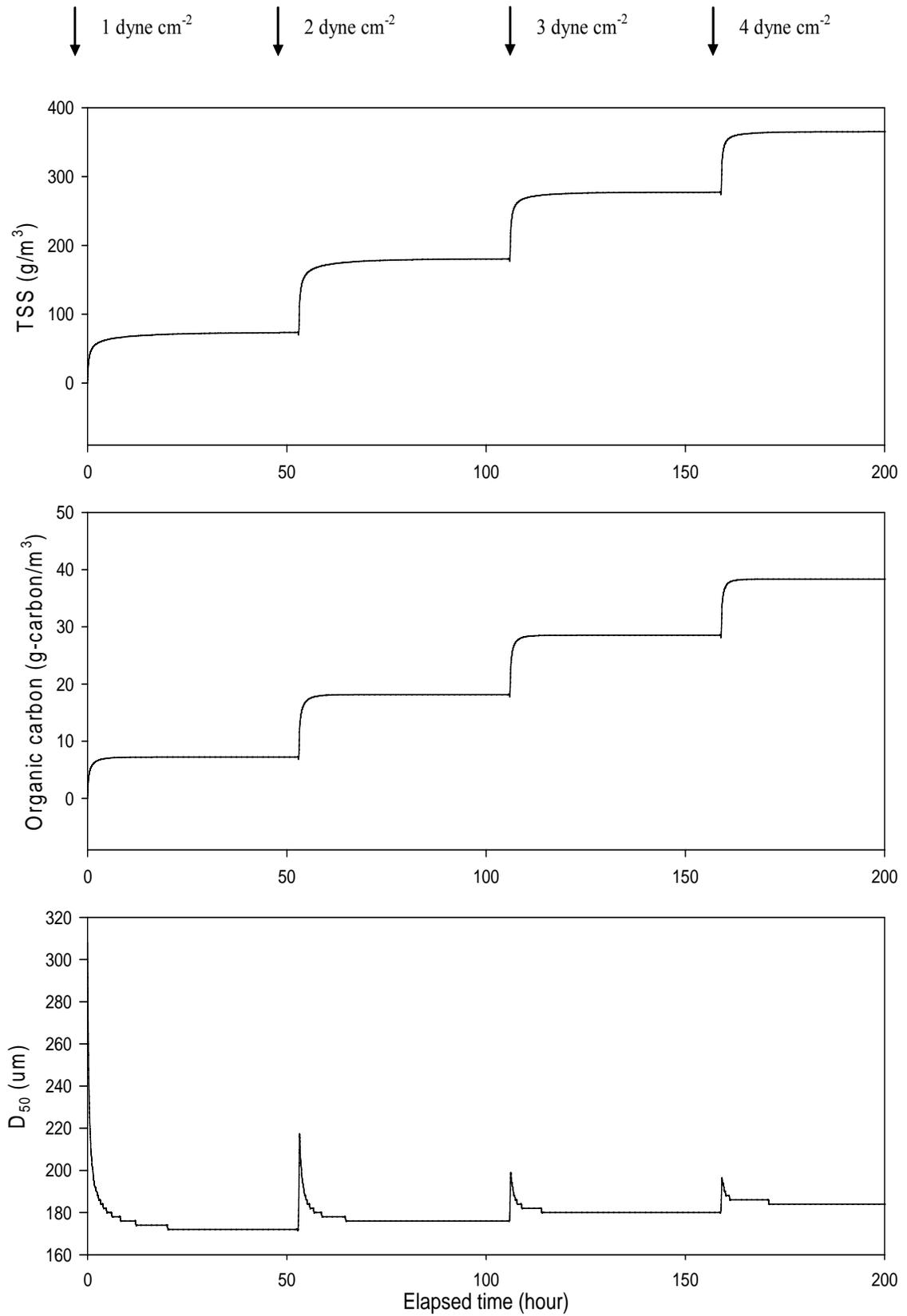


Figure 2.5:

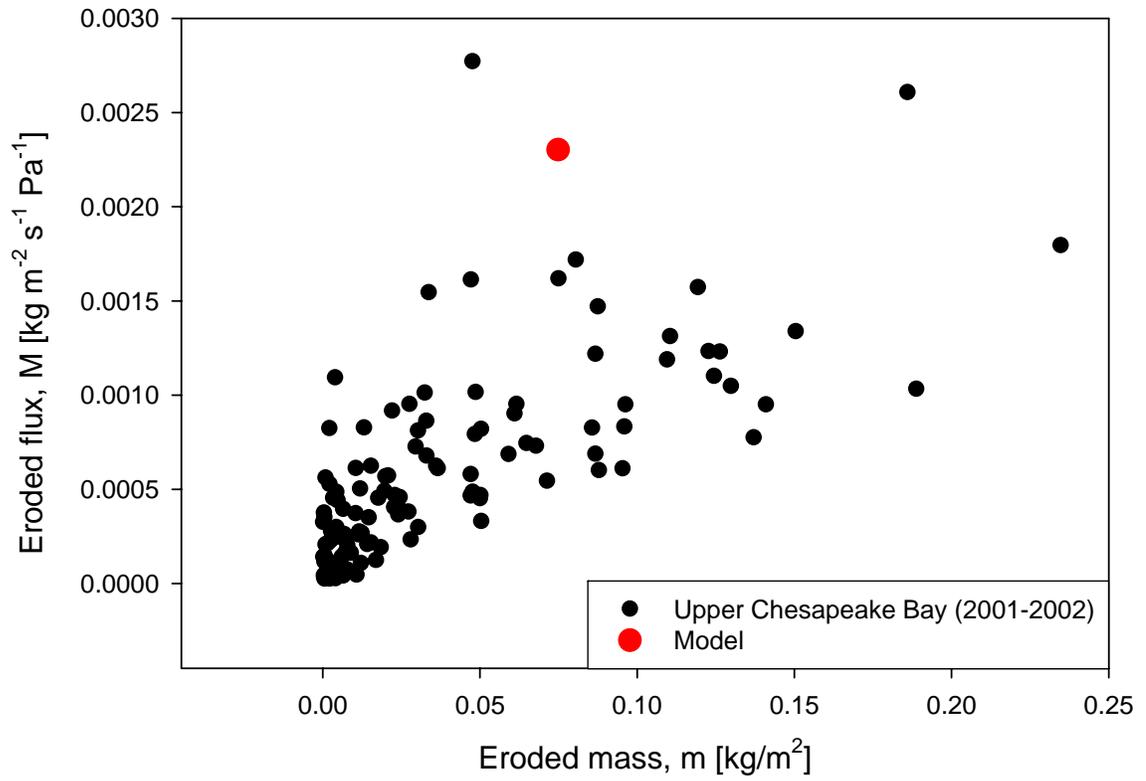


Figure 2.6 a and b

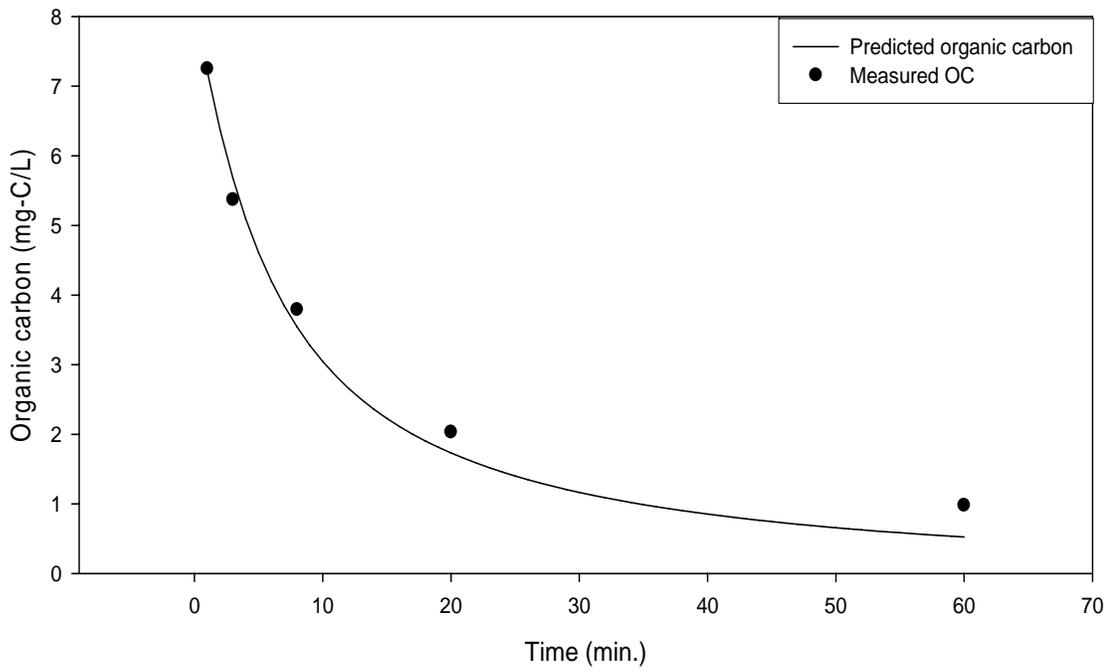
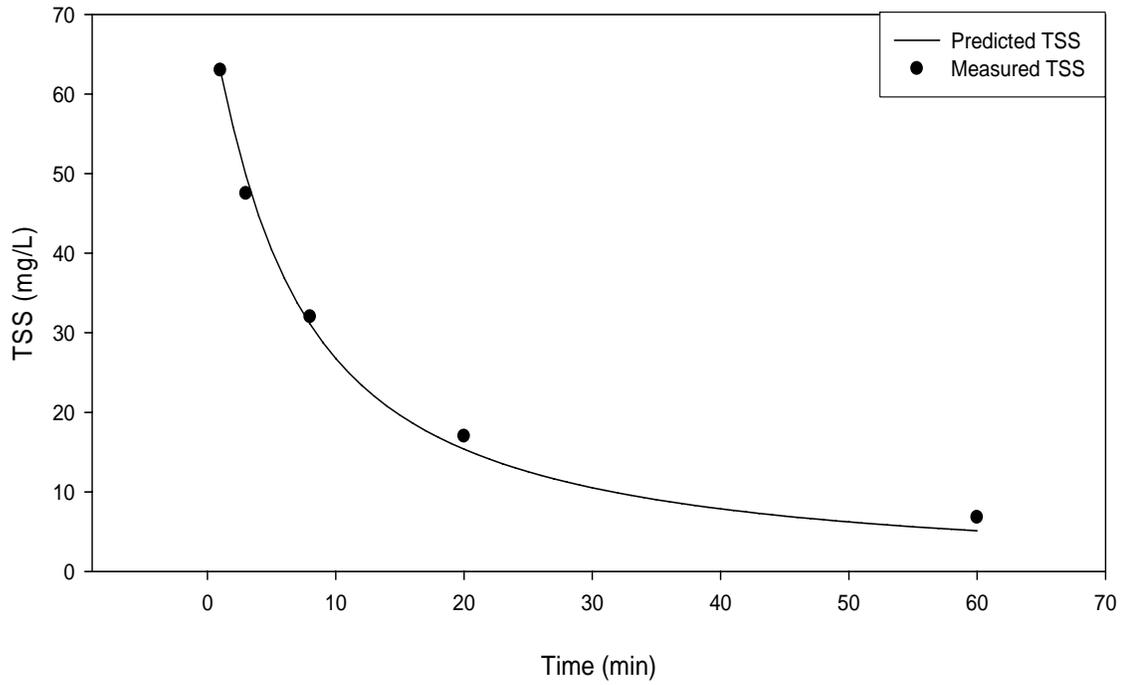


Figure 2.6 c and d

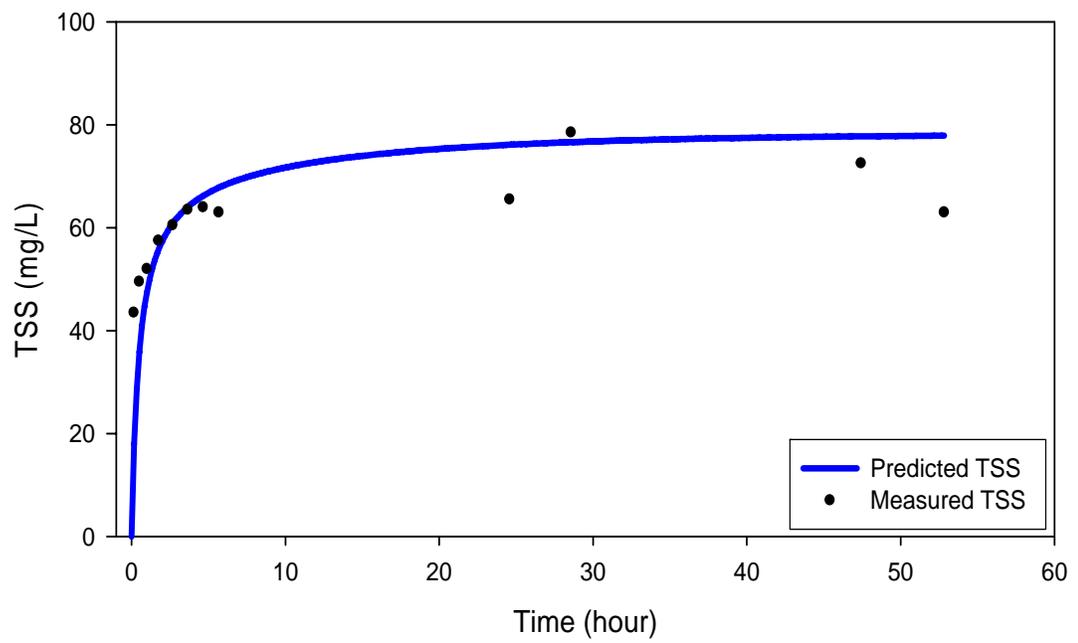
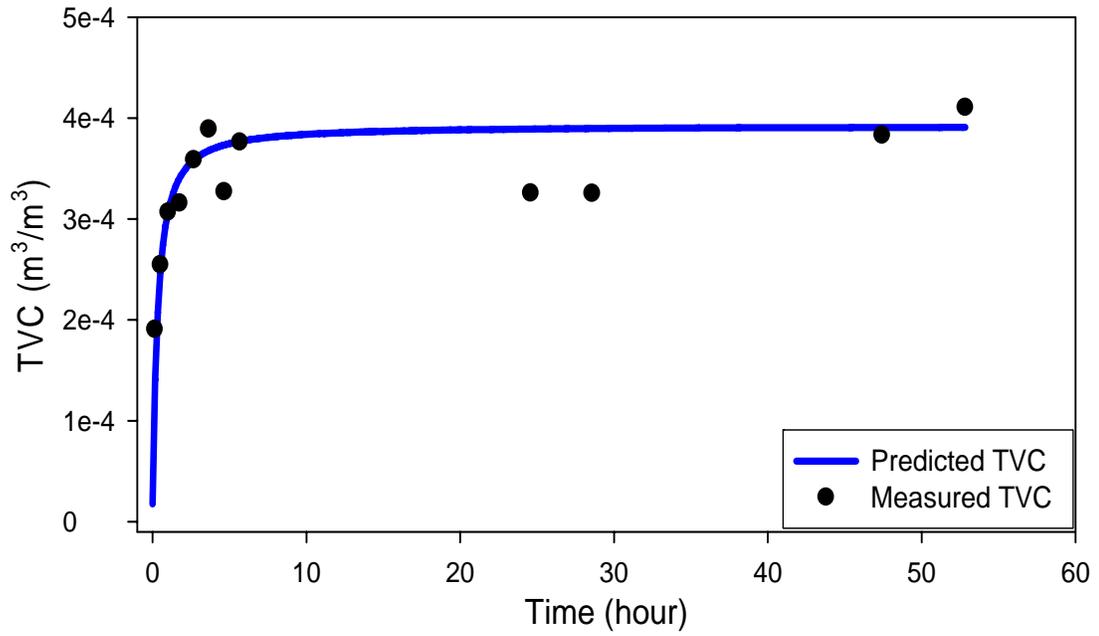


Figure 2.6 d and f

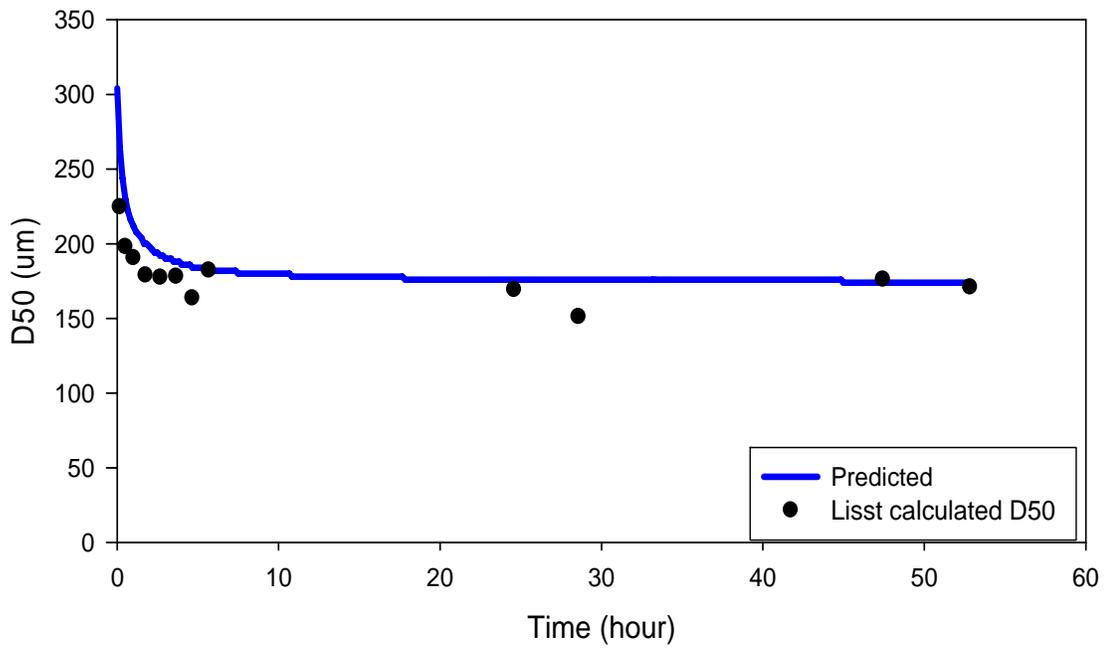
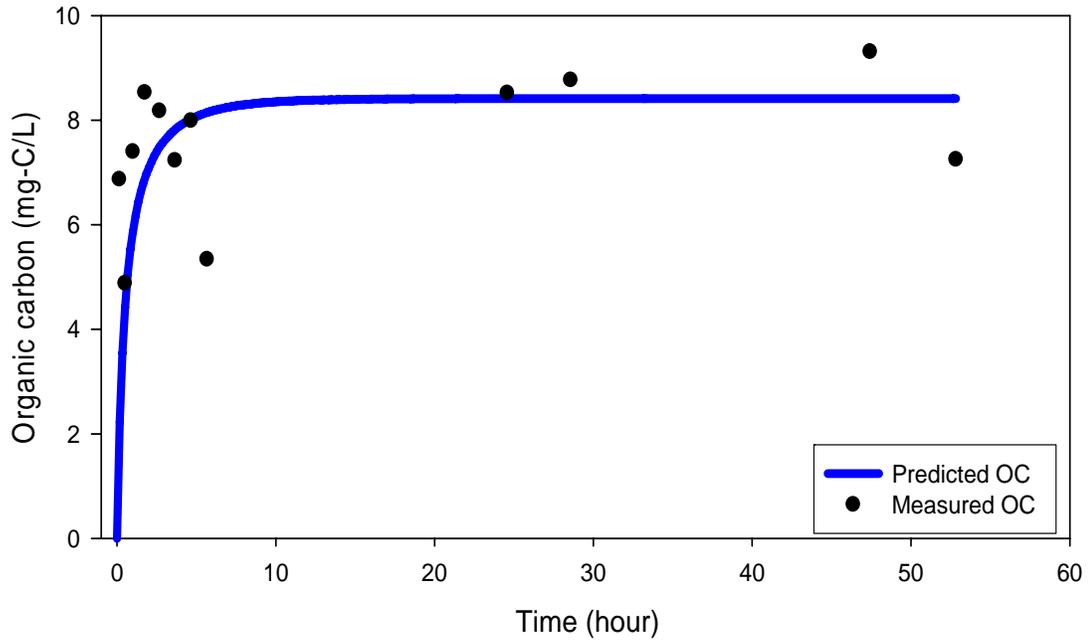


Figure 2.7:

	Scenario name	density (g/cm ³)	Settling velocity equation.
Three scenarios without the flocculation:	1. diatom	1.05	Stokes' Law
	2. field data estimated based	1.65	Stokes' Law
	3. clay	2.65	Stokes' Law
This model: One scenario with the flocculation:	4. flocculation	varied with floc composition	Winterwerp

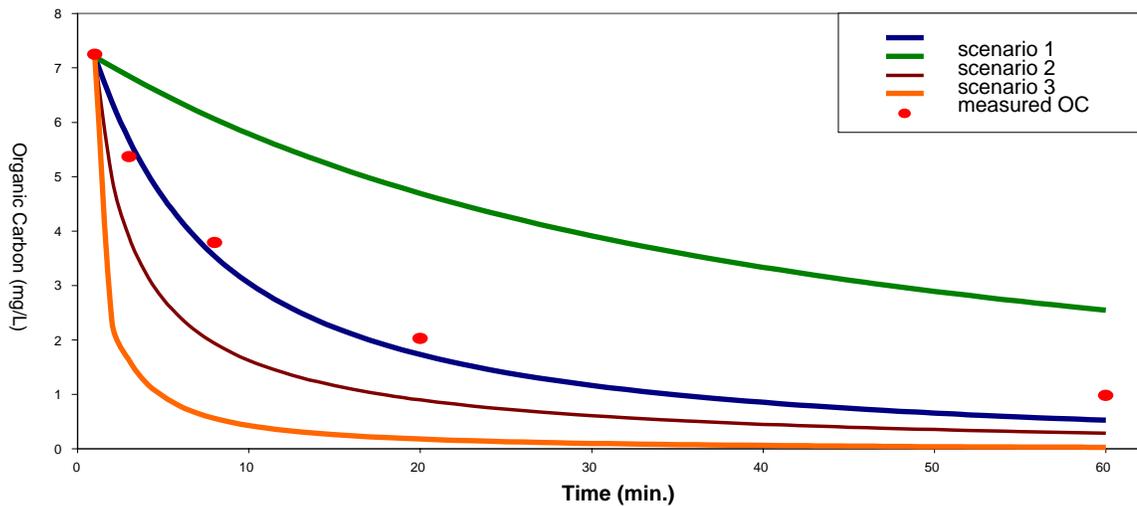
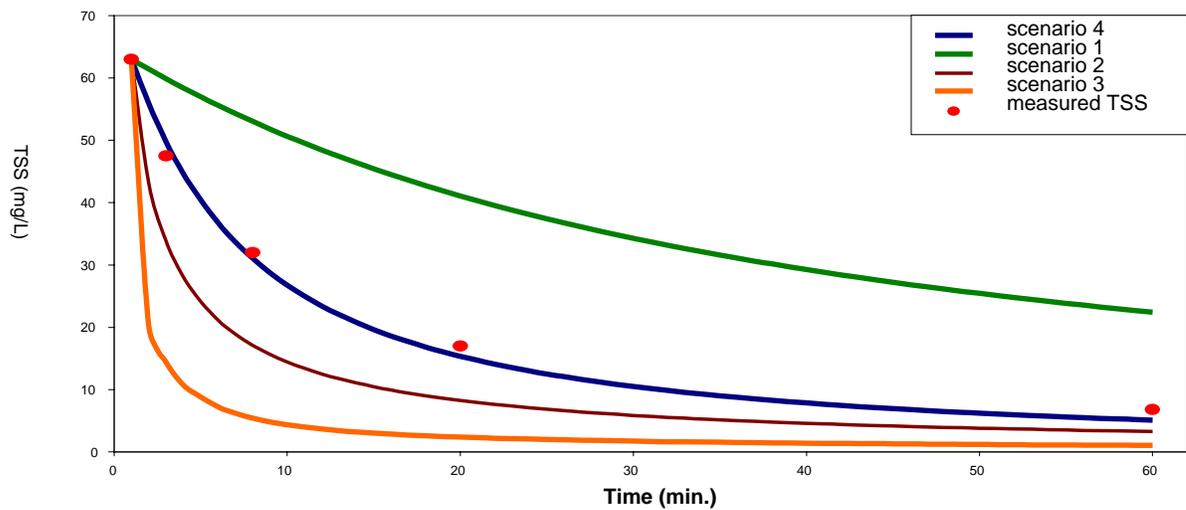


Table 2.1: The final parameter values at the end of calibration

nf	2.3	Fractal factor (dimensionless)	Model assumption
D_p	2	Primary particle size (μ m)	Model assumption
α_{Clay}	0.1; 0.4	Stickiness coefficient (w/wo stress)	Calibrated value
$\alpha_{biotic_substrate}$	0.5; 0.8	Stickiness coefficient (w/wo stress)	Calibrated value
M	2.3E-9	Eroded mass flux (kg/m ² sec Pa)	Calibrated value
ρ_{water}	1E6	density of water (g/m ³)	Model assumption
$\rho_{biotic-substrate}$	1.05E6	density of biotic substrates (g/m ³)	Model assumption
ρ_{clay}	2.65E6	density of inorganic solids (g/m ³)	Model assumption
μ	1	dynamic viscosity of fluid (g/m-sec)	STORM data
τ_c	0.0998	critical stress (dynes cm ⁻²)	Calculated value
τ_b	1	bottom shear stress (dynes cm ⁻²)	STORM data
g	9.81	Gravitational acceleration (m/sec ²)	Default value

Table 2.2: Impact of f_{OC} distribution on the behavior of organic carbon

f_{OC} = small flocs dominated initial distribution

stickiness coefficient	0	0.2	0.4	0.6	0.8	1
[TOC] at SS	36.3	30.7	28.6	27.3	26.0	25.8
TOC [time hr.] reaching SS	190.5	123	71.3	50	41.5	31.5
TSS Residence time [sec]	1107	1002	909	909	909	870
OC Residence time [sec]	2987	2529	2353	2249	2174	2120
Ws TSS weight avg [cm/sec]	0.09	0.10	0.11	0.11	0.11	0.11
Ws OC weight avg [cm/sec]	0.03	0.04	0.04	0.04	0.05	0.05

f_{OC} = even initial distribution

stickiness coefficient	0	0.2	0.4	0.6	0.8	1
[TOC] at SS	9.2	8.7	8.3	8.0	7.8	7.6
TOC [time hr.] reaching SS	50	47.3	43	37.5	32.2	27.7
TSS Residence time [sec]	769	667	667	625	625	625
OC Residence time [sec]	769	714	667	667	625	625
Ws TSS weight avg [cm/sec]	0.13	0.15	0.15	0.16	0.16	0.16
Ws OC weight avg [cm/sec]	0.13	0.14	0.15	0.15	0.16	0.16

f_{OC} = varied initial distribution

stickiness coefficient	0	0.2	0.4	0.6	0.8	1
[TOC] at SS	7.3	7.2	7.1	7.0	7.0	6.9
TOC [time hr.] reaching SS	10.8	10.2	9.5	9	8.5	8.2
TSS Residence time [sec]	711	667	667	625	625	616
OC Residence time [sec]	600	588	588	588	588	567
Ws TSS weight avg [cm/sec]	0.14	0.15	0.15	0.16	0.16	0.16
Ws OC weight avg [cm/sec]	0.17	0.17	0.17	0.17	0.17	0.18

Chapter 3: Modeling the Impact of Flocculation on the Fate of PCBs during the Resuspension Event in an Urban Estuary

3.1 *Abstract*

A multi-class flocculation-based contaminant fate model was adapted to describe desorption kinetics for contaminants associated with flocculated particles during a resuspension event. The model was effective in predicting transport of hydrophobic organic contaminants among different size flocs, water, and two sediment layers. The model also demonstrated the impact of fractal geometry, bottom shear stress, particle composition, floc size, f_{OC} , K_{OC} , and TSS on contaminant desorption rates and residence times. Under different scenarios, results from this model supported the importance of multi-floc-size, sediment-water interaction, and of flocculation on the contaminant desorption rate in the water column. Both equilibrium and kinetically-limited models predict the same dissolved and particle contaminant concentrations at steady state. However, during the first three hours of a simulated resuspension event, equilibrium and radial diffusion models overestimate the PCB desorption rate by 50% and 20% respectively. This result suggests equilibrium behavior may not be the best choice for prediction of desorption kinetics during fast events, like dredging, tidal events, or storm water runoff. The radial diffusion model, a common tool to describe desorption kinetics for a single floc, is limited by several factors during a resuspension event, as it fails to include the contaminant exchange with surrounding flocs, it has numerical difficulties in calculating the impact of various boundary

conditions, and it ignores indirect impacts on contaminant concentration from sediment-water exchange.

Further, in a floc-rich environment flocculation is an important mechanism redistributing contaminants among flocs. When flocculation is considered in a dynamic particle environment that includes sediment resuspension, settling, and kinetic-limited HOC partitioning, the steady state total PCB concentration in the water column decreases by 20 % and the water column HOC residence time decreases by 36%.

3.2 *Introduction*

Hydrophobic organic contaminants (HOCs) are important pollutants in urban estuaries. HOCs include polycyclic aromatic hydrocarbon (PAH), polychlorinated biphenyls (PCBs), and polybrominated diphenyl ethers (PBDEs). Sorption to resuspended particles and sediments plays an important role controlling the water column residence times and spatial distributions of HOC in aquatic environments. Pollutant residence times and the time required to reach sorptive equilibrium depend on properties of the chemical, of the particles, and of the surrounding environment. If rates are sorption are fast *relative* to particle residence times, HOC behavior may be described using equilibrium partitioning models. In contrast, in highly dynamic particle environments such as algal blooms, shallow water sediment-water interfaces, and during dredging operations, HOC behavior may be better described using kinetically-limited partitioning models. In laboratory studies, sorption of HOC takes from a few hours to a hundred days to reach equilibrium (Lick and Rapaka, 1996; Jespen *et al.*, 1995). The time for HOCs to reach sorption equilibrium with algae also

varies from a few minutes to a few days (Ko, 1994). Many studies suggest that PCB sorption can be described as sequential rapid (labile) and slow (resistant) steps. Carroll *et al.* (1994) suggest that the labile stage be defined as that which is adsorbed within 24 hours and a resistant stage as that which requires over one year to fully absorb. Other studies have made similar observations, finding that the first stage could be as fast as 15 minutes for PCBs in algae (Ko, 1994) and 4.5 hours for PCBs in Hudson River sediment (Schneider, 2006). The second stage may last more than 200 days for PCBs in intact sediment (Werner *et al.*, 2006). The time for each stage varies according to the type of particles examined and the molecular weight of the PCBs.

In the water column, particles are transported by many mechanisms, including water flow, sedimentation, resuspension, and grazing by biota. In the field, the particle settling velocity has a wide range (from 0.01 m day^{-1} to in excess of 100 m day^{-1}) depending on the particle size, shape, and bulk density (Graf and Rosenberg, 1997; Stemmann *et al.*, 2004). Therefore the particle residence time in a 1 meter water column varies from less than 15 minutes to longer than 24 hours. Sediment resuspension, another important particle transport mechanism, occurs when the bottom shear stress generated by wind, current, or human activity exceeds the critical shear stress required to lift a particle off of the bed. In the Upper Chesapeake Bay, the resuspension flux reached up to $82.5 \text{ g m}^{-2} \text{ day}^{-1}$ (Sanford, unpublished). Particle residence times in the deep ocean or in closed engineered systems may be months or years, but in a dynamic environment such as Baltimore Harbor the particle residence time in the water column maybe less than 4 hours due to strong water currents and

upstream flushing. Human activity including boating and dredging may also increase particle residence times.

The process of aggregation resulting from the attachment of particles colliding with each other is called flocculation. It is an important internal transport processes affecting the particle residence time in the water column. Coagulated particles are formed either in the water column or on the sediment surface, often facilitated by microbial communities. Recent research has noted the importance of flocculation in the fate of HOCs, including that flocs will affect the fate of HOC by altering the measured partition coefficient, organic carbon content, and sorption rates (Alkhatib and Weigand, 2002, Wu and Gschwed, 1986; Borglin *et al.*, 1996, Jepsen *et al.*, 1995; Ko *et al.*, 2003; Lick and Rapaka, 1996; Rounds and Pankow, 1990). The flocculation rate is a function of the stickiness coefficient and collision probability. The three major mechanisms that control collision probability are Brownian motion, differential settling, and fluid shear stress (O'Melia, 1972). The stickiness coefficient α is defined as the ratio of the particle attachment rate and the particle collision rate. Edzwald *et al.* (1974) indicated that this step is concerned with eliminating or nullifying the repulsive energy barrier that exists between the two particles. In general, exopolymeric material, suspended solids concentration, ionic strength, pH, temperature, algae type, and algae concentration are the major factors that determine the stickiness coefficient (Edzwald *et al.*, 1974, Gibbs, 1983, Kiorboe and Hansen, 1993, Liss *et al.*, 1996; Elmaleh *et al.*, 1996; Rengasamy *et al.*, 1996; Crump and Baross, 2000; Jun *et al.*, 2001; Han and Kim, 2001; Hamm, 2002). The time to form flocs may vary from a few seconds to a few days and depends on the particle

character and features of the surrounding environment, including water column depth, salinity, shear stress, and temperature (Jackson, 1989).

The question follows as to whether the population of particles changes fast enough that HOCs do not have sufficient time to reach sorptive equilibrium. If the answer is positive, how does rapid particle transport affect the HOC's sorption rate? In the water column, under conditions with rapidly changing particle populations like resuspension events and algae blooms, dissolved HOC concentrations in the water column are affected in many ways. For example, if the time for a resuspended particle-associated contaminant to resettle to the bottom is shorter than the time for the HOC to reach sorptive equilibrium, less contaminant is desorbed into the water column. Further, rapid increases in phytoplankton biomass results in an under-saturation of HOC concentrations in the algae relative to equilibrium (Swackhamer and Skoglund, 1991; Ko, 1994). Under these conditions, the chemical contaminant diffusion gradient from water to phytoplankton is strongest during peak primary production (Dachs and Eisenreich, 2000). Resuspension events also last on the order of a few hours to a few days and are significantly shorter than the time needed to reach sorptive equilibrium. Cheng *et al.* (1995) compared the particle residence time under the impact of turbulence with the PCB residence time for the Buffalo River. They concluded that the PCB residence time was much longer than the particle residence time during the resuspension period under all scenarios.

For an internal process like flocculation, the time to reach sorptive equilibrium for HOC and the time to form flocs both vary widely, and we can not easily determine which process is faster under a given condition. Lick and Rapaka (1996)

conducted a series of experiments to examine the relationship between the time required for floc size to reach steady state and the time required to reach sorptive equilibrium under a given suspended solids concentration and water column shear stress. Under certain environments, like those with higher TSS concentrations, the time to form flocs is faster than the time to sorptive equilibrium (Lick and Rapaka, 1996). However, their experiments also involved a dynamic disaggregation process. Therefore, the time difference between forming a single floc with flocculation process and a sorbing PCBs to the reach equilibrium is still not clear.

Recent research has noted the importance of flocculation on the fate of HOCs, including that flocs will affect the fate of HOCs by altering the partition coefficient, the organic carbon content of the flocs, and the sorption rates (Alkhatib and Weigand, 2002, Wu and Gschwed, 1986; Borglin *et al.*, 1996, Jepsen *et al.*, 1995; Ko *et al.*, 2003; Lick and Rapaka, 1996; Rounds and Pankow, 1990). Flocculation can impact the fate of HOC in two ways. The direct impact is changing the mass transfer velocity that controls the sorption process. Mass transfer velocity is a function of floc size, porosity, K_{OC} , f_{OC} , the molecular diffusion coefficient, and floc dry density. The porosity, floc size, contact area, f_{OC} , and dry density vary as flocs are formed. The second indirect impact of flocculation on HOC cycling results from changing the particle water column residence times. When flocs are formed, the particle settling velocity changes and the adsorbed contaminant have different residence times in the water column.

The particle size distribution varies with time in the water column as flocs are formed and change properties. Contaminant partitioning also varies in the water

column due to changes in f_{OC} , the particle porosity, the diffusion distance from particle distance from center to the shell, and the particle dry density. Several papers have discussed the relationship between grain size and field measured partition coefficient for methyl-mercury and HOCs in the sediment, water, and atmosphere (Krauss and Wilcke, 2002; Hayes *et al.*, 1998). These papers suggest the contaminant residence time changes along with the variation of particle contact area. Schneider *et al.* (2007) conducted mesocosm experiments that mimic resuspension and the settling of contaminated upper Hudson River sediment with realistic bottom shear stress and water column turbulence. In that study, they also observed that both the measured PCB partition coefficient (K_p) and the particle size distribution varied temporally throughout the experiments.

Therefore, the objectives of this paper are to explore the question of how HOC sorption is affected in a rapidly changing particle environment, especially changes resulting from flocculation. Four approaches to modeling HOC sorption have been previously developed to simulate HOC-particle interactions, including equilibrium partitioning (Di Toro, 1985), a single box partitioning model (Oddson *et al.*, 1970), a two compartment diffusion model (Pignatello and Xing, 1996; Gong and DePinto, 1998), and a radial diffusion model (Wu and Gschwend, 1986). Equilibrium partitioning is the simplest approach that assumes dissolved and particulate HOC concentrations are invariant and a function of particle porosity, total organic carbon concentration (TOC), fraction of organic carbon, and the octanol-water partition coefficient. A single box partitioning model assumes that sorption process can be calculated as a first order function and controlled by a single rate constant estimated

from laboratory experiments. The two compartment diffusion model separates particles into two compartments: an exterior compartment that has a faster exchange rate, and an interior compartment that has a slower exchange rate in order to have better agreement with the experimental data. The radial diffusion model considers the sorption-retarded diffusion within flocculated particles (flocs). Wu and Gschwend (1988) used the radial diffusion model to simulate desorption of HOCs within and between porous particles, the sediment surface layer, and the surrounding water.

In the field, the sorption process shows a two-step sorption behavior and often takes a long time to reach equilibrium (Gong and DePinto, 1998). Therefore, the equilibrium and single box model behaviors are often too simple to predict the fate of HOCs over short periods. Further, both models assume a static particle composition. Therefore, they may poorly estimate the HOC residence time and dissolved and particle-sorbed concentrations in dynamic environments. However, these relatively simple partitioning models may still be applicable for long-term predictions if the short term dynamics average out. Dynamic environments also affect the performance of the radial diffusion model, because this model assumes a fixed boundary on a single particle (Gong and DePinto, 1998). Natural sediment particles have complex composition and therefore it may not be proper to use a single-sized particle like median particle diameter (D_{50}) to represent the entire particle spectrum. Therefore, modelers often adjust the particle size along with other physical and chemical parameters to fit experimental data from resuspended sediment HOC desorption experiments (Werner *et al.*, 2006; Schneider, 2007). Furthermore, the complex boundary condition setting required for the radial diffusion model limits its

application when particle properties are varying. Application of the two compartment diffusion model in a dynamic environment has a different problem. Three independent fitting sorption rate constants, including the fraction of exterior particle volume to the total particle volume, the sorption rate between surrounding water and the exterior compartment, and the sorption rate between exterior and interior compartments, must be specified. These input parameters are often difficult to estimate and vary among different experiments or environments.

In an urban estuary, both resuspended and bottom sediment particles show non-uniform, non-normal size distributions (Schneider, 2007). The contaminant also has different residence times or times required to reach steady state, which correspond to the different floc size within the rapidly varying boundary condition. Therefore, models need the ability to simulate how the entire spectrum of particles varies in space (sediment versus water column) and time (*i.e.*, before, during, and after resuspension events). A new model that can simultaneously simulate the varying particle population and contaminant partitioning would better describe resuspended contaminated sediment in an urban estuary. The model should also be able to simulate the dynamics of particle exchange across sediment/water interface, settling of particles and contaminants through the water column, production of flocs, and HOC sorption and desorption rates.

3.3 *Model Development*

3.3.1 Objectives and Strategies

1. Develop a model to investigate how flocculation influences the fate of organic contaminants when both HOCs and sorbent concentrations vary rapidly.
2. Evaluate whether a kinetic sorption model more accurately describes the fate of PCBs during a resuspension event compared to equilibrium.
3. Discuss the impact of kinetic processes such as flocculation and erosion/deposition on the fate of PCBs; for example, the time to reach a steady state, the desorption rate, the dissolved and particulate PCBs concentrations at steady state, and the temporal variation in the deviation from sorption equilibrium during the resuspension event.
4. Discuss the differences between D_{50} and multi-floc-size models. This objective responds to the conclusion from DePinto (1998) that the radial diffusion model is too simple and does not reflect the composition and properties of real sediment.

3.3.2 Model Assumptions and Structure

In this work, the structure of the model is optimized for hydrophobic organic contaminants (HOCs), a broad class of chemicals that include PCBs, dioxins and furans, PAHs, many brominated flame retardants, and organochlorine pesticides. The model tracks particles, organic carbon, and HOCs in the water column and in the sediment. In principle, any chemical contaminant with a known affinity for organic carbon can be modeled within this framework. In the previous chapter, the model

described the exchange of particles between water and sediment surface layers and related flocculation mechanisms. However, many HOC fate models have emphasized the importance of a deeper sediment layer on the long term HOC fate. Therefore, a deeper sediment layer and related surface-deeper sediment exchange mechanisms are added to the model in this study (Figure 3.1). Porewater also plays an important role in this contaminant model by linking HOC among surficial sediment, deeper sediment, water, and resuspended particles. In this model, porewater HOC is defined as the dissolved contaminant concentration in the sediment layer; it connects the HOC source from the deeper sediment layer and allows diffusive exchange with the overlying water column.

In the STORM resuspension experiments (Schneider *et al.*, 2007), properties of the resuspended particles vary with time in many ways. The median diameter (D_{50}) of particles in the water column varied, flocs in the sediment and water column had different D_{50} and particle size distributions, and flocs rapidly exchanged between the sediment and the water column. These observations indicate that flocs properties are not constant during the experiment. In addition, the dynamic resuspension-settling and flocculation processes might cause particle-related HOCs to not have enough time to reach steady state in the water column. The traditional radial diffusion model may fail to simulate the fate of field HOCs because of the complexity of particle properties and the broad particle size distribution, which complicates parameterization and setting boundary conditions. The traditional one box or two-compartment diffusion models cannot directly model situations where particles are

changing rapidly, because both HOCs and related flocs varied rapidly, which violates assumptions of constant particle concentrations and characteristics.

The HOC sub-model simulates chemicals being transported among particles, and uses particle models as a platform to transport HOCs among flocculated particles, surrounding water, sediment porewater and sediment particles. This model also tracks the fluxes of particle-related HOCs among all particle size clusters. These fluxes include both flocculation and resuspension-settling. To resolve the above concerns and achieve the study objectives, a new kinetically-limited diffusion model was developed.

First, we assume that the HOC will immediately reach equilibrium or become well-mixed between pore water and solids within each flocculated particle to simplify the boundary condition and numerical techniques. For example, it is difficult to set the intra-floc HOC distribution and related boundary condition after two different size flocs coagulate to form a new larger floc. To compensate for the overestimated HOC concentration gradient, we modified the chemical diffusive exchange rate by including the floc size effect and adjusted diffusion coefficient. Further, to simplify the model, we assumed no stagnant film layer between floc exterior shell and surrounding water in this study. The detail of these assumptions will be explained in the next section.

Second, flocs sizes are varied between 2 and 1000 μm in 1 μm intervals. Each cluster is a state variable and represents a specific size of floc. Number and organic carbon concentrations, particulate HOC concentration, mass transfer velocity, stickiness coefficient, floc density, settling velocity, and f_{OC} are calculated for each

cluster simultaneously at each time step. The floc spherical equivalent volume and area are calculated based on floc diameter in each cluster. The volume and mass concentrations vary temporally in each cluster, and the gross dry mass and volume are conserved. In addition, only flocs that contain organic carbon transport HOCs and we assume no diffusion exchange of contaminant into pure inorganic flocs. To include bioturbation and mixing from the deeper sediment layer, the model simultaneously simulates HOCs being transported among the water column, top sediment layer (0.1 cm), and deeper sediment layer (4.9 cm).

3.3.3 Major Model Equations

The water column particle model in this chapter adapts the same equations and assumptions used in the previous chapter:

$$\frac{\partial N_i}{\partial t} = F_i + \left(\frac{W_{s,i}}{x} \times N_i \right) + E_{b,i} \quad \text{Equation 3.1}$$

where N represents the number of particles of class i per volume, $W_{s,i}$ is the floc settling velocity, F_i is the flocculation effect on particle number balance, x is the water column depth, and $E_{b,i}$ is pure sediment resuspension flux that is given by the sediment erosion model.

Regarding the sediment layer, a few additional terms are added in this chapter. To simulate the particle exchange between the surface and the deeper sediment, this model adapts the strategy from Di Toro (2001) that simulates this process as a simple diffusive exchange. In order to simplify this process, the model ignores consolidation in both layers and focuses on bioturbation exchange. The sediment consolidation rate is relatively slow compared with other mechanisms in a 53 hour simulation period.

$$\frac{\partial N_{i,surface}}{\partial t} = -\frac{m_{1,2}}{(sz_1 \times sz_2)} (N_{i,surface} - N_{i,deeper}) - E_{b,i} \quad \text{Equation 3.2}$$

$$\frac{\partial N_{i,deeper}}{\partial t} = \frac{m_{1,2}}{(sz_1 \times sz_2)} (N_{i,surface} - N_{i,deeper}) \quad \text{Equation 3.3}$$

where N represents the number of particles of class i per volume in the surface or deeper sediment layer, $m_{1,2}$ is the particle mixed rate between two sediment layers, which is controlled by the bioturbation effect; sz_1 is the top sediment thickness, and sz_2 is the deeper sediment thickness.

The contaminant model simulates HOCs being transported among resuspended flocs, dissolved water, porewater in the sediment, and sediment particles. There are 3000 state variables, including 1000 state variables in the water column and 2000 state variables in the two sediment layers, to represent the HOCs in the flocs from size 2 to 1000 μm . In addition to the state variables for solid HOCs, the model also includes three state variables representing the dissolved HOC concentrations in the water column and in porewater in two sediment layers. In general, each floc state variable gains or loses contaminants through diffusive exchange, and by contaminants 'piggybacking' on the organic carbon flows calculated by the flocculation submodel

$$\frac{\partial HOC_i}{\partial t} = HOC_flux_{F_n, floc} + J_{i,j} \quad \text{Equation 3.4}$$

where HOC_i represents the HOC concentration in the floc cluster i with the unit of $\mu\text{g}/\text{m}^3$, $HOC_flux_{F_n, floc}$ represents the HOC flux due to particle transport by flocculation, resuspension, settling, or bioturbation in the water column or sediment layer, and $J_{i,j}$ represents the HOC being transported between particles variable i and surrounding water variable j.

In this study, the diffusive contaminant uptake and release has a specific equation to emphasize the temporal varied organic carbon concentration. The rate of diffusive exchange between the dissolved contaminant pool and each carbon phase is calculated as the product of a diffusional gradient and a mass transfer velocity. The diffusional gradient is defined as the difference between the instantaneous contaminant concentration in the carbon phase and that at equilibrium with the surrounding dissolved contaminant concentration. The equilibrium condition, as parameterized by K_{oc} , is an input to the model. The mass transfer velocity equals a velocity term multiplied by the specific interfacial area of the carbon phase. Overall, the bi-directional diffusive flux equals:

$$J_i = k \times A \times N_i \times \left(HOC_{dissolved} - \frac{HOC_i}{OC_i \times K_{oc}} \right) \quad \text{Equation 3.5}$$

where represents the HOC being transported between particles variable i and surrounding water variable j , k is the mass transfer velocity (m/sec), A is the specific interfacial exchange area (m²/floc), OC_i is the organic carbon state variable concentration (g-carbon/m³), $HOC_{dissolved}$ is the dissolved HOC concentration (ng-HOC/m³), HOC_i is the particulate HOC state variable i concentration (ng-HOC/m³), and K_{oc} is the organic carbon normalized partition coefficient (m³/g-carbon).

The first important element in the diffusion flux equation is that the model includes the mass transfer velocity (k). In this study, we therefore calculate the effective diffusion coefficient based on the dry density, porosity, K_{oc} and f_{oc} of the floc (Schwarzenbach *et al.*, 1993),

$$k = \frac{D^*}{z} = \frac{\phi_f f D_m}{z \times (K_{OC} f_{OC} \times (1 - \phi_f) \times \rho_{f,solid} + \phi_f)} \quad \text{Equation 3.6}$$

where k is the mass transfer velocity (m/sec), z represents the volume-weighted effective interior diffusion distance, f is the tortuosity factor, ρ_f is floc dry density (g/m³), ϕ_f is porosity (m³/m³), D^* is the adjusted diffusion coefficient (m²/sec), and D_m is the molecular diffusion coefficient of the chemical in water (m²/sec).

Since effective interior diffusion thickness, porosity, f_{OC} , and floc dry density vary, either with floc size or temporally, the chemical diffusive exchange rate is not a constant and also varies temporally along with floc size in this study.

The second important element in the diffusion flux equation is that the model includes the total floc contact area. The traditional radial diffusion model describes diffusion within a single flocculated particle. However, the sorption/desorption rate has been observed to vary with particle mass or number concentrations (Schwarzenbach *et al.*, 1993). For example, the dissolved PCB concentration should be higher after resuspending many contaminated particles than for adding a single particle during a given time period. To simulate this effect, the total diffusion contact area has been added to the diffusion equation as the product of area per floc (A , m²/floc) and total floc number concentration (N , number of flocs/m³). This raises the question of how to estimate the area per floc. Our model is based on the floc volume being equal to that of spherical primary particles combined in flocs following fractal geometry (Winterwerp, 1998). Therefore, the method of estimating the interior floc contact area per aggregated particle becomes complex. The distance between the floc

center and the shell is floc radius r . Therefore, the contact area per floc particle is simply estimated as $4\pi r^2$.

Many papers suggest effective radial diffusional distance should be equaled to half the floc radius to estimate the effective of floc thickness (Gong and DePinto, 1998; Werner *et al*, 2006). In this study, the volume-weighted average distance (z) from the interior of a sphere to the outer shell of the sphere is $R/4$ instead of $R/2$, where R is the radius of the sphere (Sanford, personal communication) (Figure 3.2). The radial diffusion model calculates that during resuspension into clean water the HOC concentration will be highest at the floc center and will approach $\text{HOC}_{\text{dissolved}}$ at the outer shell, where $\text{HOC}_{\text{dissolved}}$ represents the dissolved concentration in the surrounding water or stagnant film layer. However, in this study we assume a uniform HOC concentration distribution within the modeled floc particle sphere shell as explained in the previous section. Therefore, with this assumption, the mass flux may overestimate the flux because the concentration difference is too great. As a result, we need to compensate for the concentration gradient (Sanford, personal communication). To calculate the volume-weighted mean distance from any point in the interior of a sphere to the outer shell of the sphere, the distance from the center of the sphere is defined as r and the radius of the sphere as R : The distance to the outer shell is $(R-r)$ and the appropriate weighting for each of these distances is the area of the internal spherical shell of radius r . Thus, the volume-weighted average radial diffusional distance is $R/4$ instead of $R/2$:

$$(R - r) = \frac{1}{\frac{4}{3}\pi R^3} \int_0^R 4\pi r^2 (R - r) dr = \frac{\frac{1}{3}\pi R^4}{\frac{4}{3}\pi R^3} = \frac{1}{4}R = z \quad \text{Equation 3.7}$$

3.3.4 Numerical Methods, Parameterization, Boundary Conditions, and Initial Conditions

The model requires an initial concentration for each floc size, and STORM experiment measurements (Schneider *et al.*, 2007) are applied as model initial conditions. Initially only sediment layers contain HOCs and we assumed HOCs reach equilibrium between flocs and porewater in the sediment layers for all size of flocs, calculated using the given K_{OC} , organic carbon size distribution, and experimental measurements. The f_{OC} distribution derived in Chapter Two is used to initialize the sediment carbon.

Several parameter values and initial conditions were estimated from the combination of observations during the STORM experiment and literature reports, including the fractal factor, the characteristics of the primary particle, porosity, and initial particle volume and f_{OC} size distribution. In this chapter, these parameters are derived using the strategies and references in Chapter 2. K_{OC} was estimated by measurements in the STORM tank experiments at steady state and varied among the different PCB congeners. The molecular diffusion coefficient was back calculated by algae PCB uptake rate from Ko (1994), the porewater-water diffusion coefficient was selected from the Lake Michigan PCB model (2007), and the sediment bioturbation rate was chosen from Di Toro (2001).

The system of equations was solved for 1000 floc sizes for each variable. The model was written in double precision Fortran 90 and uses a fourth-order Runge-Kutta numerical algorithm to solve the set of nonlinear ordinary differential equations. The model is numerically stable with a time step of one minute, and a 53

hour simulation is completed in 75 CPU minutes on a Sun Ultra-60 workstation. Total dry mass, total particle dry volume, and total contaminant mass in the system are monitored to insure conservation of solids and contaminant mass during the simulations.

3.3.5 The Impact of Floc Property on the Mass Transfer Velocity and Desorption Rate

Mass transfer velocity (k) is a function of floc size, porosity, K_{OC} , f_{OC} , D_m , and floc dry density. It is the most important parameter in the radial diffusion model controlling the sorption/desorption rate. The first objective of this section is to test the response of mass transfer velocity by using our new diffusion equation with the given conditions. The second objective in this section is to discuss the impact of these variables on the HOC desorption rate.

To have a better understanding of the mass transfer velocity in this new diffusion model, several simple model scenarios and assumptions are tested in this section. Initially, the impact of fractal geometry, total contact area, and the concentration gradient between dissolved and particulate phases are ignored to explore two different porosity-size distributions. The first and simplest relationship is that all flocs have the same porosity. The second relationship is that the porosity and floc size follow fractal geometry.

3.3.5.1 The Influence of Constant Porosity on the Desorption Rate

In the first case, the floc porosity is constant for all sizes, resulting in the mass transfer velocity decreasing with floc size, which implies that HOC require a longer time to diffuse to the floc outer edge in larger flocs (Figure 3.3a). Further, based on

the equation 3.5, a faster desorption rate should be observed in the higher mass transfer velocity run. Table 3.1 shows the settings and the corresponding assumptions of Runs 1 to 7. Based on Equation 3.6, the mass transfer velocity will increase with increasing D_m and porosity, and decrease with increasing floc size, K_{OC} , f_{OC} , and solid density. As shown in Figure 3.4 and Table 3.1, the desorption rate has a strong positive correlation with the mass transfer velocity. Results from Run 1 to Run 7 demonstrate that our new HOC fate model follow the proper trends, namely that the desorption rate increases with increasing D_m and porosity, and decreases with a increasing floc size, K_{OC} , f_{OC} , and dry floc density.

We next discuss the impact of more than one parameter or variable on the desorption rate. From Run 1 to Run 7, we observed that the mass transfer velocity increases with some parameters and decreases with others. However, when two or more parameters influence the mass transfer velocity calculation, their combined effect on the desorption rate is not clear. For example, when f_{OC} is increased the particle solid density also decreases since organic carbon is less dense than clay. An increase in K_{OC} should lead to a decrease in D_m and an increase in floc size should result in an increase in porosity. However based on the Run 1 to Run 7 results, these parameters have the opposite influence on the mass transfer velocity calculation and trend in the desorption rate.

Lick and Rapaka (1996) also observed that the run with larger floc, lower solid density, and higher porosity has a faster desorption rate than the run with smaller floc, higher solid density, and lower porosity. Lick's experimental data can be used to examine the impact of more than one parameter or variable on the desorption rate

(Runs 8 and 9, Table 3.1 and Figure 3.5). These model runs ignore the impact of TSS, fractal geometry, total contact area, and the gradient between dissolved and particle HOCs. The model simulation results agree with the observations of Lick and his colleagues that the larger and more porous flocs showed a faster desorption rate than small and lower porosity flocs.

3.3.5.2 The Influence of Size-Dependent Porosity on the Desorption Rate

In the model runs above, we assumed porosity is constant for all floc sizes to isolate the influence of porosity independent of size-specific behavior. Here the particle sub-model is based on fractal geometry and to emphasize the importance of flocs the fractal factor was set less than 3.0 and greater than 1.0 in all runs. Here we explore the influence of fractal geometry on the HOC mass transfer velocity (Figure 3.3b). The floc character described by Lick and Rapaka (1996) has a similar trend in the mass transfer velocity as predicted using fractal geometry. When the fractal factor is less than 3.0, the larger floc has a lower solid density and a higher porosity. One interesting finding is that for a given fractal factor the mass transfer velocity only varies significantly for smaller flocs. For mid- to large-sized flocs, the dependence of the mass transfer velocity on floc size is very limited. There are two points in this section: First, discuss how fractal geometry impacts the mass transfer velocity. Then discuss the impact of fractal geometry on the total contact area, concentration gradient, desorption rate, and diffusion equation.

To simplify the model setting, in this section the influence of the total contact area and the concentration gradient between the dissolved and particulate phases are ignored. Therefore, for a given floc size the runs with the higher fractal factor

contains less porous flocs and smaller mass transfer velocity (Table 3.2). For a given fractal factor, larger flocs have higher porosity and higher mass transfer velocity than smaller flocs. In general, based on these assumptions, model results showed a similar trend as the fixed porosity runs in that the desorption rate increased with increasing D_m , and porosity, and decreasing floc size, K_{OC} , f_{OC} , and dry floc density (Figure 3.6).

Next, the model runs allow the number concentration and contact area to vary with mass concentration, fractal factor, and floc size. We have discussed the relationship between fractal geometry and the mass transfer velocity in the previous section. The fractal factor not only impacts the mass transfer velocity because porosity is a function of the fractal factor, but also affects the number concentration for a specified total solid volume and mass concentrations. Also, the diffusion flux is controlled by the mass transfer velocity, total contact area, and concentration gradient between the dissolved and particulate phases. The total contact area is a product of contact area per floc and number concentration, which is determined by the mass concentration, the fractal factor, and the floc size. The detailed model settings are listed in the Table 3.3. The model was executed for 300 simulation days without flocculation, deposition, and erosion. All model runs start with a concentration of 1 ng/g PCB 52. Run 18 behaviors are a reference for runs 19 to 22. Run 19 applies the same parameters except the fractal factor equals 1.5 to increase floc porosity to 0.996; Run 20 increases the floc size from 50 to 120 μm ; run 21 doubles the initial f_{OC} of the floc, and run 22 increases the TSS concentration ten fold to examine the influence of total contact area and concentration gradient on the desorption rate.

Most model run results showed similar trends to those in Tables 3.1 and 3.2 that increasing the solid density or decreasing the porosity decreases the diffusion rate (Figure 3.7). However, run 20 showed an opposite result from run 13 regarding to the floc size influence that the net diffusion flux is faster for the smaller floc. A larger floc has a higher mass transfer velocity than a smaller one with the same fractal factor. However, under the same solid mass concentration, the smaller floc has a higher total contact area. Thus, the net product of mass transfer rate (k) and total contact area ($A \cdot N$) is higher for the smaller floc since the diffusion flux includes the total contact area. Run 22 results also showed the importance of the total contact area. Comparing runs 18 and 22, the mass transfer velocity are the same for both runs because they have the same floc characters. However, run 22 has a higher total contact area in the water column. Therefore, the net diffusion flux is higher in run 22, which is why we observe a higher desorption rate in the sediment than in the water column for the same type of particles. Porosity plays a dominant role in the diffusion flux, and when the fractal factor decreases the difference in desorption rates with particle size diminishes. In some extreme situations, a larger floc has a faster desorption rate than a smaller floc under the same initial conditions due to its high porosity which compensates for the disadvantages arising from the total contact area and size adjustments.

3.3.6 Model Data

Schneider *et al.* (2007) conducted Shear Turbulence Resuspension Mesocosms (STORM) tank experiments that mimic resuspension and settling of contaminated upper Hudson River sediment with realistic bottom shear stress and water column

turbulence. The maximum instantaneous bottom shear stress was about 1 dyne cm^{-2} and the volume-weighted average water column turbulence intensity and energy dissipation rate were 0.55 cm s^{-1} and $0.0032 \text{ cm}^2 \text{ s}^{-3}$, respectively. Hudson River sediment was added to a depth of 5 cm and allowed to consolidate for 10 days. The two periods of this experiment are the erosion and the free settling. During erosion, the mixing paddle continuously generated bottom shear stress for 53 hours to ensure suspended solids reached steady state. During the one hour free settling period, the paddle was then turned off to allow suspended particles to settle. Particle size distribution, dissolved and particulate PCBs, TSS, DOC, particulate carbon, nitrogen, and chlorophyll *a* were measured throughout the resuspension and settling portions of each experiment. The detailed PCBs measurement methods were reported by Schneider *et al.* (2007).

The flocculation particle sub-model was calibrated using Shear Turbulence Resuspension Mesocosm (STORM) tank experiments. Further, STORM tank experiments also provided data for the HOC sub-model, including the initial sediment PCB concentration and the K_{OC} parameter value. Later, results from the HOC sub-model are compared without further calibration with the STORM experiment measurements.

3.4 *Model Applications*

3.4.1 The Impact of Flocculation on the Steady State Contaminant Concentrations and on the Desorption Rate

This section discusses the impact of flocculation and floc transport on the water column residence time of contaminants.

Settings: In this section, the following settings were applied to the model: the fractal factor is equal to 2.0, f_{OC} is equal to 0.1 and the bottom shear stress is the same as the STORM experiment as 1 dyne cm^{-2} . Particles are assumed to be well mixed in the tank without disaggregation, deposition or erosion or any other interaction with sediment layers. The model was executed for 300 simulation days beginning with 1 ng/g of PCB 52 on 50 μm flocs without flocculation in the base case run (Run 23). Run 24 and Run 25 had the same settings except the stickiness coefficient was increased from 0 to 0.25 and 0.5 respectively. Run 26 had the same settings as the base case run except the 1 ng/g of PCB 52 began on 400 μm flocs. To examine the impact of flocculation on contaminant residence time, the model used data from the STORM free settling period for the floc submodel settings and for initial conditions in the contaminant submodel.

Results and discussion: Table 3.4 showed the detail model simulation results in all runs. The run with flocculation had a slower PCB 52 desorption rate than runs without flocculation (Figure 3.8), consistent with previous results described above that the PCB desorption rate is slower for larger flocs than smaller flocs under the same initial conditions. These results are also consistent with the concept that the number concentration or total contact area is a dominant factor controlling the

diffusive flux. Larger flocs have a faster desorption rate due to higher porosity and corresponding higher mass transfer velocity. However, as mentioned above, the diffusive flux is also controlled by the total contact area. Under the same solid mass and fractal factor value, larger flocs have a lower number concentration than smaller floc, which impacts the desorption flux. Therefore, desorption rate for run 23 was 50% faster than run 26. Desorption rates calculated from runs 24 and 25 were slower than from run 23, because large flocs were formed faster with the higher stickiness coefficient. Therefore, the total water column contaminant residence time decreases with enhancing flocculation process.

There is no sediment resuspension during the STORM free setting period, highlighting the influence of flocculation on contaminant residence time. With flocculation (Run27), the total water column HOC residence time is five times shorter than without flocculation (Run 28) (Figure 3.9). Because the particulate HOC residence time is shorter in the flocculation run larger flocs formed. This result implies that in calm floc-rich water the particulate HOCs are transported more quickly from the water column to the surficial sediment.

3.4.2 The Impact of Deposition-Erosion on the Steady-State Contaminant Concentrations and Desorption Rate

Previous flocculation models (Lick and Rapaka, 1996) did not simulate sediment erosion and settling under realistic bottom shear stress, and were designed with well-mixed sediment particles forming flocs during a long mixing period. However, during shorter events, such as tidal periods, or in dynamic environments, contaminated particles might not have enough time to desorb contaminants to the surrounding

water. This is especially true for large particles. The objective of this section is to discuss the impact of deposition and erosion on the PCB desorption rate under realistic bottom shear stresses.

Settings: In this section, the following settings were set for the model: the fractal factor is equal to 2.0, f_{OC} is equal to 0.1, and flocculation is not included in the three runs. The sediment PCB 52 concentration is 702 ng/g-carbon, consistent with the STORM measurements. The model was executed for 300 simulation days three separate times, with the bottom shear stress equal to 0, 1, and 2 dyne/m² (Run 29, Run30, and Run 31) respectively.

Results and discussion: In Figure 3.9 PCB 52 desorption rates are compared between runs with no erosion and with two different bottom shear stresses. In general, because these runs ignore flocculation, the higher bottom shear stress run erodes more flocs from the sediment to the water column. The results are consistent with runs discussed in the previous section that showed that higher TSS increases the desorption rate because more flocs are eroded from the sediment to increase the total floc contact area in the water column.

The PCB desorption rate is faster in the run without erosion and settling compared with the run under identical conditions except particles are exchanging between the sediments and overlying water (Figure 3.12). This trend was observed more clearly for the larger flocs or after including flocculation. The overall PCB desorption rate is slower in the desorption-erosion run for two reasons. As showed in the Figure 3.11, the porewater contaminant reaches equilibrium in a short time because of the higher concentration of solids in the sediment. Thereafter the

contaminant has a smoother diffusion gradient between sediment and porewater to decrease the desorption rate. Second, the time that flocs remain in the sediment due to deposition will delay the overall desorption rate than the run without deposition-erosion. In past studies, particles are often entirely suspended in the water column under a given bottom shear stress with no deposition, which may lead to errors in estimates of the desorption rate.

In the previous section, we developed a basic understanding of impacts from diffusion, flocculation, and resuspension-settling mechanisms on the HOC diffusion rate. Now we can begin to discuss the effects arising from both flocculation and resuspension-settling. When we include both flocculation and resuspension-settling into the model, the PCB desorption rates are fastest in the diffusion-only run (Run 32), followed by the flocculation-only run (Run 33), the deposition/erosion-only run (Run 34), and finally the full process run (Run 35) (Figure 3.12). With flocculation, the steady-state organic carbon concentration in the water column decreased 28%, the total HOC concentration decreased 57%, and desorption rate decreased 13% relative to the run that only included resuspension-settling and kinetic-limited HOC partitioning without flocculation.

The deviation from equilibrium was also observed for each floc cluster as a function of time, assuming that the equilibrium state is defined by K_{ow} ,

$$\text{Deviation from equilibrium} = (C_{p,t} - C_{p,t}^*)/C_{p,t}^* \quad \text{Equation 3.8}$$

where $C_{p,t}$ is the concentration of chemical contaminant in a particle phase (g contaminant/m³ water) at time t and $C_{p,t}^*$ is the corresponding particle phase

contaminant concentration in equilibrium with the surrounding dissolved phase at time t ($C_{d,t}$, g contaminant/m³-water),

$$C_{p,t}^* = K_{OW} C_{d,t} C_{dry,i} \quad \text{Equation 3.9}$$

where K_{OW} is the octanol-water partition coefficient and $C_{dry,i}$ is the flocs concentration (g organic carbon/m³-water). Note that if $C_{p,t}$ exceeds $C_{p,t}^*$, the particles are oversaturated with contaminant with respect to the dissolved phase, resulting in net desorption. In the opposite case, chemical contaminant diffuses into the particulate carbon phase. Because $C_{p,t}$, $C_{d,t}$, and F each vary somewhat independently, the deviation from equilibrium and, therefore, the diffusive flux is constantly changing for each particle type throughout the model simulation.

Under the given assumptions and initial conditions, the diffusion-only run reached sorption equilibrium faster than did the flocculation, deposition/erosion, and full process runs as discussed in the previous section. However, the deviation from equilibrium from the full process and the deposition/erosion runs are quite similar. This result hints that deposition/erosion is the dominant mechanism for slowing the approach to equilibrium. Furthermore, during the rapid transport between water and sediment, the contaminated flocs do not have enough time to fully desorb contaminants to the dissolved phase.

3.4.3 Examining the Impact of Floc Size Distribution on the Desorption Rate

Up to this point, the model has simply used a single floc size parameter to represent the entire spectrum. However, D_{50} may not well represent the behavior of the entire particle size spectrum. In the previous chapter, the field-measured particle size distribution was used as an initial condition for the floc sub-model. Therefore,

here we test the impact of including a full size distribution on the desorption rate in simulations including flocculation and deposition/erosion.

Settings: To test the model performance and desorption rate with a full floc size distribution (multi-size-flocs that D_{50} was calculated as 216 μm in the previous chapter) as an initial condition, the model combined the particle size distribution described in Chapter 2 with the PCB data measured in the STORM experiments. To compare the model results from previous runs, all runs are normalized to an initial TSS of 1 mg/L. Values for other parameters for flocculation and deposition/erosion were the same as in previous sections.

Results and discussion: The PCB 52 desorption rates among three single floc runs (50, 216, and 400 μm) (Run 23, Run 36, and Run 26) and the multi-size run (Run 37) are compared in Figure 3.13. The desorption rate for a multi-size run was intermediate between the two single floc runs (Run 23 and Run 26). This result suggests that when the fate of HOC is controlled only by diffusive exchange between flocs and the surrounding water, the floc size is the dominant parameter and that D_{50} value for multi-size run (216 μm) is between two sizes used in the single size runs (50 μm and 400 μm respectively). D_{50} may be a useful tool when diffusion is the only transport mechanism and has a normal particle size distribution because there the desorption rates are equal between the multi-size and single-size runs. However, when flocculation or deposition/erosion is included and particle size distribution is not shown as normal distribution, the multi-size run performs differently than the single-size run (D_{50}), and desorption rates are controlled by the particle size distribution. For example, when the distribution is skewed to smaller flocs, the multi-

size run desorption rate is faster than that of the single-size run. Therefore, using D_{50} to estimate desorption rate would miscalculate the short term results. Furthermore, the multi-size run results show a similar trend as the single floc run, with the full effect run having the slowest desorption rate, followed by the resuspension-settling run, the flocculation run and finally the diffusion only run (Figure 3.14). The multi-size run results were intermediate between the small and larger floc size single run results, because its medium floc size is located within these ranges.

3.4.4 Comparing Modeled PCB Behavior and Observations from STORM

Experiments

Here the fully-developed, flocculation-based contaminant model is used to predict the behavior of four PCB congeners under the STORM mesocosm conditions and to compare the results with measured data. Note that while the particle transport model was calibrated with STORM observations, the HOC submodel was not tuned to PCB measurements made during these experiments. These model results are used to explore the role of flocculation and resuspension on the fate of HOCs in the STORM experiments.

Settings: The floc sub-model used the settings and parameters calibrated from the resuspension stage of the STORM experiment, as described in Chapter Two. The initial f_{OC} , PCBs concentrations, K_{OC} values, and sediment PCBs concentrations were the same as STORM measurements and varied with corresponding PCB congener. In this section, the initial f_{OC} and PCBs concentrations are assumed the same between upper and deeper sediment layers.

Results and discussion: The model was tested with di, tri, tetra, and penta PCB congeners and the results were compared with STORM experiment measurements (Figure 3.15-3.18). The model results agreed well with the STORM measurements of particulate and dissolved phase concentrations of four PCB congeners without calibration. The larger PCB congeners required a longer time to reach steady state than the less chlorinated PCBs due to their smaller mass transfer velocity. The more chlorinated PCB congeners have a higher steady-state particulate-bound PCB concentration because of their higher K_{OC} values.

When flocculation is considered in a dynamic particle environment that includes sediment resuspension, settling, and kinetic-limited HOC partitioning, the steady state total PCB concentration in the water column is decreased by 20%, and time to reach steady state in the water column decreased by 36%, consistent with earlier findings at section 3.1.

3.4.5 The Impact of Surficial Sediment Contaminant Concentrations on Steady State Dissolved PCBs Concentrations

Settings: There are three runs in the section, runs 41-43. All are based on the same floc transport model and related parameters including bioturbation coefficient and chemical diffusion coefficient between each layer. The physical environment is the same as the resuspension period in the STORM experiment throughout 53 simulation hours for each run. The surficial sediment thickness is 0.1 cm and the deeper sediment layer thickness is 4.9 cm. All three runs assumed that top sediment layer PCB 52 concentration was 702 ng/g dry weight. However, deeper sediment layer PCB 52 concentrations were assigned values of 702, 0.702, and 7020 ng/g dry

weight respectively in the three runs to test the influence of deeper sediment PCB levels on subsequent water column concentrations.

Results and discussion: Run 41, where there was no PCB concentration gradient in the sediment, is the base case. The deeper sediments in run 42 are depleted in PCBs relative to the surface sediment layer. The PCB 52 concentration in the water column increased during the first few hours following by a smooth decline in both dissolved and particulate phases (Figure 3.19) because the surficial sediment was eroded and PCB quickly desorbed during the first few hours. When sediment was eroded to the water column, the deeper cleaner sediment was recruited to replace the lost surficial layer. At the same time, before contaminated sediment reaches equilibrium, it might cycle between the water column and sediment. As a result, the dissolved and particulate PCB concentrations did not decrease dramatically. In addition, the sediment-water diffusion flux buffers the dissolved PCB concentration, making the dissolved PCB decrease slower than the particulate PCB between five and ten simulation hours. When the model reached steady state, the PCB concentrations in the run 42 with cleaner deeper sediment was almost 10 times less than in the base case, which was similar to the ratio of deeper sediment concentrations between the two scenarios. Run 43 simulated a deeper sediment layer containing 10 times higher PCB levels compared to the surficial later. Both the dissolved and particulate PCB 52 concentrations were 10 times higher than the base case. These results indicated that deeper sediment concentration not only impacts the long term water column contaminant concentrations, but also temporally impacts the variation in desorption rate.

3.4.6 Comparing Kinetically Limited Partitioning and Equilibrium Behaviors on the Steady State Contaminant Concentrations

Objectives: Several papers suggest that the traditional radial diffusion model and equilibrium partition model overestimate the rate of desorption during resuspension events (Cornelissen *et al.*, 1997; Gong and Depinto, 1998; Wener *et al.*, 2006). In this study we apply our new multi-class flocculation-based contaminant fate model to several scenarios and compare the results predicted by the model with equilibrium and radial diffusion models under the same physical conditions. Finally, we compare the model predictions from three models with results from the STORM experiments.

Settings: The first scenario estimated the desorption rate between kinetically limited partitioning and equilibrium behaviors (run 44) with flocculation, deposition, and resuspension. In this scenario, the model was used to test two kinetically-limited partitioning behaviors including radial diffusion (run 45) and our new kinetically-limited behavior (run 47). In these runs, the model used the settings derived from the STORM experiment during the erosion period for the flocc-sub model and included flocculation, deposition, and erosion. Values for all chemical parameters for PCB 52 were based on the experimental measurements or from calculations: K_{OC} equals to 1.1 m³/g, D^* equals 2.75×10^{-6} cm²/s, and sediment C_p is 702 ng/g for both sediment layers. Further, the median particle radius D_{50} for the radial diffusion model, and TSS, dry density, and f_{OC} for the equilibrium model were all calculated outputs from the flocc-submodel. The equilibrium partitioning calculation included contaminant transport with floc movement by settling and resuspension. Further, for the radial diffusion behavior, the resuspension-settling process had been simulated using the

measured temporal varied D_{50} values (Schneider, 2007). Furthermore, for the radial diffusion behavior, the temporal varied D_{50} values had also been adjusted to fit the STORM measurements (run 46). The second scenario starts with clean water and dirty suspended particles and uses the same settings for the floc sub-model derived from the STORM experiment during the beginning of the free settling period. The first two runs (run 47 and 48) in this scenario compared the water column residence time of PCBs predicted by the equilibrium and kinetically-limited partitioning behaviors that included flocculation and settling.

Results and discussions: As described in the previous chapter, organic carbon and TSS reach steady state concentrations and size distributions during the first few hours of constant erosion in the STORM experiment. The equilibrium partitioning model predicts an instantaneous increase in the dissolved PCB concentration at the onset of resuspension. However, as showed in Figure 3.20, only about 50% of the contaminant is desorbed relative to the equilibrium value during the first 3 hours of the resuspension experiment. After 12 hours of constant resuspension, desorption reached 83% of equilibrium during the first 12.5 hours. Even during the STORM free settling period experiment, the equilibrium partitioning model predicts PCB water column residence times that are two times longer than the kinetically-limited partitioning model (Figure 3.21). One important difference between this study and previous models (DRBC, 2003) that assume partitioning equilibrium is that a contaminant does not reach sorptive steady state within a tidal period. In other words, even from a long term point of view, an equilibrium model may not be a proper choice to describe HOC behavior in an estuary or other dynamic tidal area.

The radial diffusion model also overestimates measured dissolved PCB concentrations. Most previous experiments were designed so that all particles were resuspended and well mixed in a container that lacked the interaction between sediment and water column (Lick and Rapaka, 1996). Therefore, some parameters have to be manually adjusted in order to apply the radial diffusion model to a field project. Several possible factors might explain the overestimation.

First, as shown in the previous section, the radial diffusion model only implicitly includes the feedback from the sediment layer from D_{50} values, which decreased the desorption rate for the resuspended particles in the water column. Second, the radial diffusion model does not simulate contaminant gain and loss from a floc size from flocculation. Flocculation not only changes the concentration gradient between floc and surrounding water, but also changes the boundary conditions for the floc particles. A single floc particle will not only lose contaminant by diffusion but will also gain or lose contaminant as it is coagulates with flocs with different contaminant concentrations. Third, although we applied measured D_{50} as boundary for the radial diffusion model, this strategy might not appropriate to model this variation within the given time interval, which caused numerical problems by temporally altering the boundary and initial conditions. In addition, to fit the STORM measurement, the model had to manually increased 70% of the given D_{50} values to decrease the desorption rate.

3.5 *Conclusions*

A new multi-class flocculation-based contaminant fate model was adapted to describe desorption kinetics for contaminants associated with flocculated particles

during a resuspension event. The model was effective in predicting transport of hydrophobic organic contaminants among different size flocs, water, and two sediment layers. The model also demonstrated the impact of fractal geometry, bottom shear stress, particle composition, floc size, f_{OC} , K_{OC} , and TSS on contaminant desorption rate and residence time.

Under different scenarios, this model results supports the importance of size distributed flocs, sediment-water interaction, and flocculation for the contaminant desorption rate in a water column. Both equilibrium and kinetically-limited partitioning approaches predict the same dissolved and particle contaminant concentrations at steady state. However, during the first three hours of a simulated sediment resuspension event, the equilibrium and radial diffusion behavior overestimated dissolved PCB concentrations 50% and 20% respectively. This result suggested equilibrium behavior overestimates the initial PCB release from contaminated sediments during relatively rapid events like dredging, tides, or storm water influx. The radial diffusion model, a common tool to describe desorption kinetics for a single floc, is limited by several factors, as it fails to include the contaminant exchange with surrounding flocs, it has numerical difficulties in calculating the impact of various boundary conditions, and it ignores indirect impacts from sediment-water exchange.

Flocculation could alter the contaminant residence time in many ways such as by changing the particle settling velocity and diffusion flux. Further, in a floc-rich environment, flocculation is an important mechanism redistributing contaminants among flocs. When a full flocculation is considered in a dynamic particle

environment that includes sediment resuspension, settling, and kinetic-limited HOC partitioning, the steady state total PCB concentration in the water column is decreased by 20 % and water column HOC residence time decreased by 36%.

3.6 *Figure Captions*

Figure 3.1: The conceptual diagram of the HOC fate model.

Figure 3.2: The conceptual diagram of the contaminant distribution within a single flocculation particle

Figure 3.3: Mass transfer velocity (m sec^{-1}) varied with floc size with two porosity-floc size trends. The first trend (2a) assumed porosity was a constant along all size of flocs; the second trend (2b) assumed porosity was a function of fractal factor and floc size.

Figure 3.4: Variation in desorption rates under different scenarios. Assumed all sizes of flocs have the same porosity and a constant product between number concentration and floc contact area. The detail settings are shown in Table 1.

Figure 3.5: Using Lick *et al.* (1996) experiment data to compare the impact of multi-floc-property on desorption rates as model Run 8 and Run 9 in Table 1.

Figure 3.6: Comparing desorption rates under different scenarios including base case run, smaller fractal factor run, larger fractal factor run, larger floc size run, higher f_{OC} run, lower solid density run, higher K_{OC} run, and higher D_m run. Assumed porosity is controlled by the fractal geometry and a constant total floc contact area.

Figure 3.7: Comparing desorption rates under different scenarios including base case run, smaller fractal factor run, larger floc size run, higher f_{OC} run, and higher TSS run. Assumed porosity and total floc contact area are controlled by the fractal geometry.

Figure 3.8: Comparing the impact of flocculation on desorption rates including based run, stickiness coefficient = 0.25 run, stickiness coefficient = 0.5 run, and stickiness coefficient = 0 alone with floc size = 400 μm run.

Figure 3.9: Comparing the particulate PCB 52 residence time with and without flocculation using non-equilibrium partitioning behavior based PCBs fate model during the free settling period.

Figure 3.10: Comparing the impact of bottom shear stress on the PCB 52 desorption rate. The bottom shear stress varied from 0, 1, and 2 dynes/m^2

Figure 3.11: Comparing the impact of resuspension-settling on the desorption rate in the dissolved water and sediment porewater.

Figure 3.12: Comparing the impact of diffusion, flocculation, resuspension-settling and all above processes on the desorption rate that was initiated with 50 μm and 1 mg/L flocs.

Figure 3.13: Comparing the PCB 52 desorption rate among three initial conditions: 50 μm , 400 μm , and STORM particle size distribution respectively. All runs were involved diffusion only.

Figure 3.14: Comparing the impact of diffusion, flocculation, resuspension-settling and all above mechanisms on the desorption rate when initiated with normalized STORM experiment particle size distribution with 1mg/L flocs for 300 days.

Figure 3.15-3.18: The model was tested with di (PCB 4 and 10), tri (PCB 19), tetra (PCB 52), and penta (PCB 77 and 110) PCB congeners and the results were compared with STORM experiment measurements.

Figure 3.19: Comparison of the temporal concentration variations for particulate and dissolved PCB 52 under three deeper PCB 52 concentrations: Run 27: L1=14.01 ug/g-OC (702 ng/g -dry) ; L2= 14.01 ug/g-OC; Run 28: L1=14.01 ug/g-OC (702 ng/g -dry) ; L2= 1.401 ug/g-OC; Run 29: L1=14.01 ug/g-OC (702 ng/g -dry) ; L2= 140.1 ug/g-OC.

Figure 3.20: Comparison model simulated PCB 52 desorption trends among measured, equilibrium behavior, radial diffusion model, and this study under STORM experimental conditions.

Figure 3.21: Comparing the particulate PCB 52 residence time between equilibrium and non-equilibrium partitioning behaviors using calibrated flocculation model during the STORM experiment free settling period.

Figure 3.1:

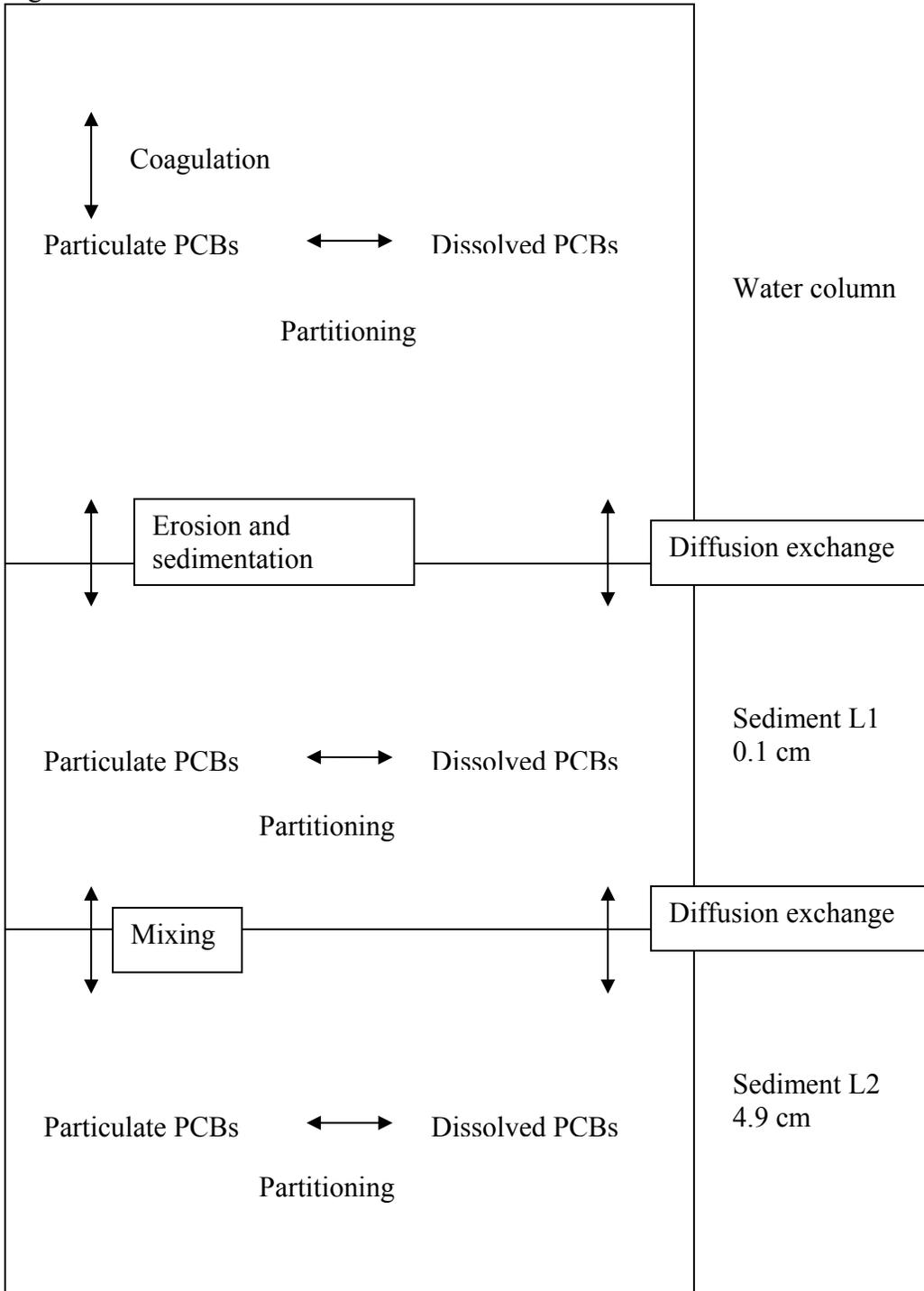


Figure 3.2

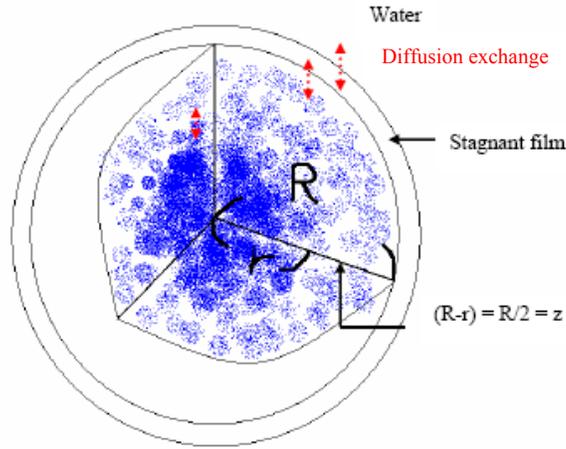


Figure 3.2a: The ideally conceptual diagram of the contaminant distribution within a single flocculation particle, where the blue dots represent the contaminants

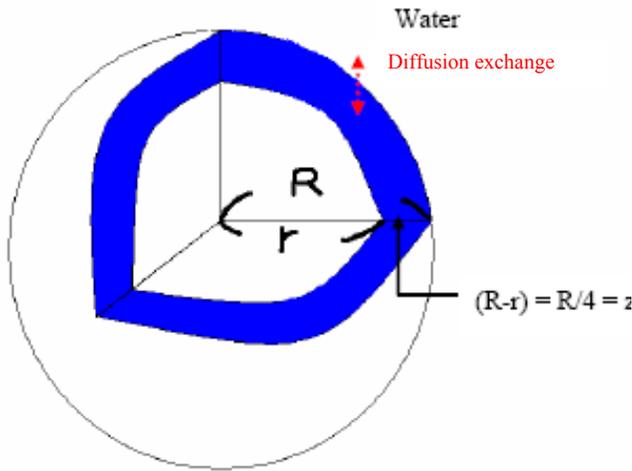


Figure 3.2b: The conceptual diagram of the contaminant distribution within a single flocculation particle in this study, where the blue dots represent the contaminants

In reality, the concentration will be high at the center and will approach $HOC_{dissolved}$ at the outer shell, where $HOC_{dissolved}$ is the dissolved concentration in the surrounding water (Figure 3.2 a). This study assumes a uniform HOC concentration distribution within the modeled floc particle sphere shell to simply the boundary condition when two flocs collided to form a new one (figure 3.2 b). Thus, writing the mass flux as may overestimate the flux because the concentration difference is too large.

$$J_i = \frac{D^*}{z} \times A \times N_i \times \left(HOC_{dissolved} - \frac{HOC_i}{OC_i \times K_{OC}} \right) \quad \text{Equation 3.5}$$

Therefore, we have to compensate by changing the length scale slightly

$$\text{Old: } (R - r) = \frac{1}{2} R = z$$

$$\text{New: } (R - r) = \frac{1}{4/3 \pi R^3} \int_0^R 4\pi r^2 (R - r) dr = \frac{\frac{1}{3} \pi R^4}{\frac{4}{3} \pi R^3} = \frac{1}{4} R = z \quad \text{Equation 3.7}$$

Figure 3.3:

Figure 3.3a

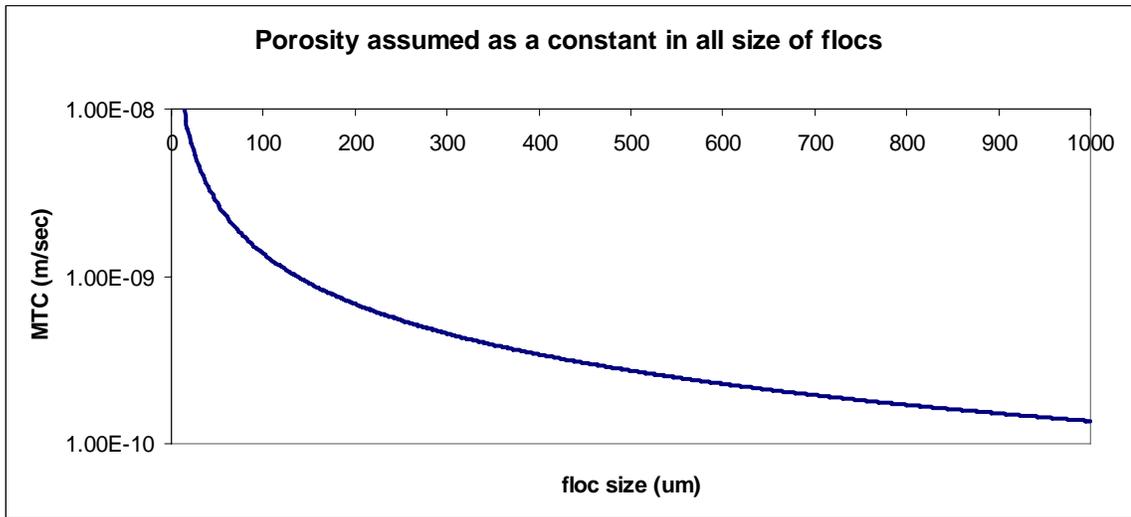


Figure 3.3b

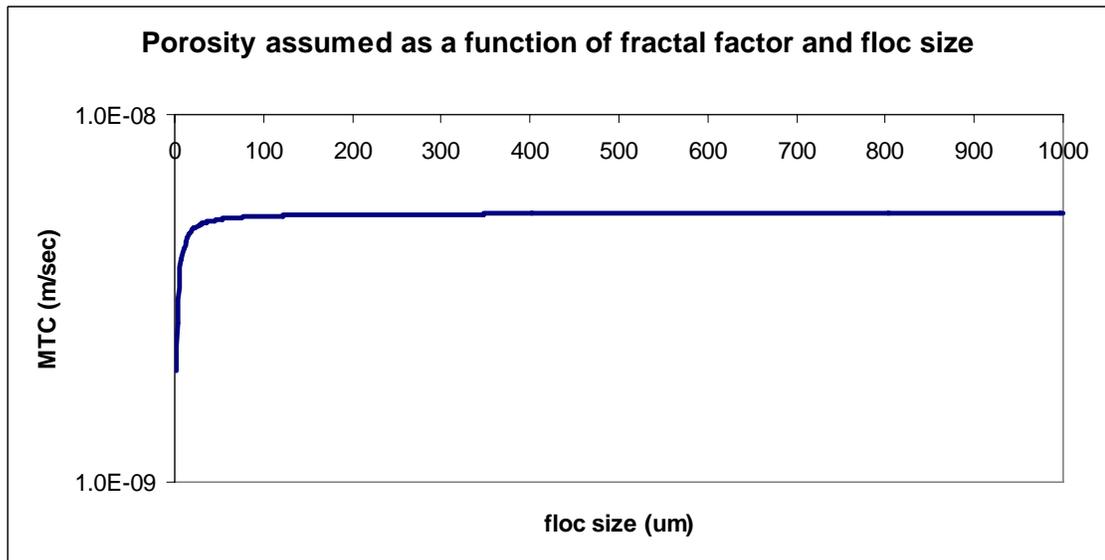


Figure 3.4:

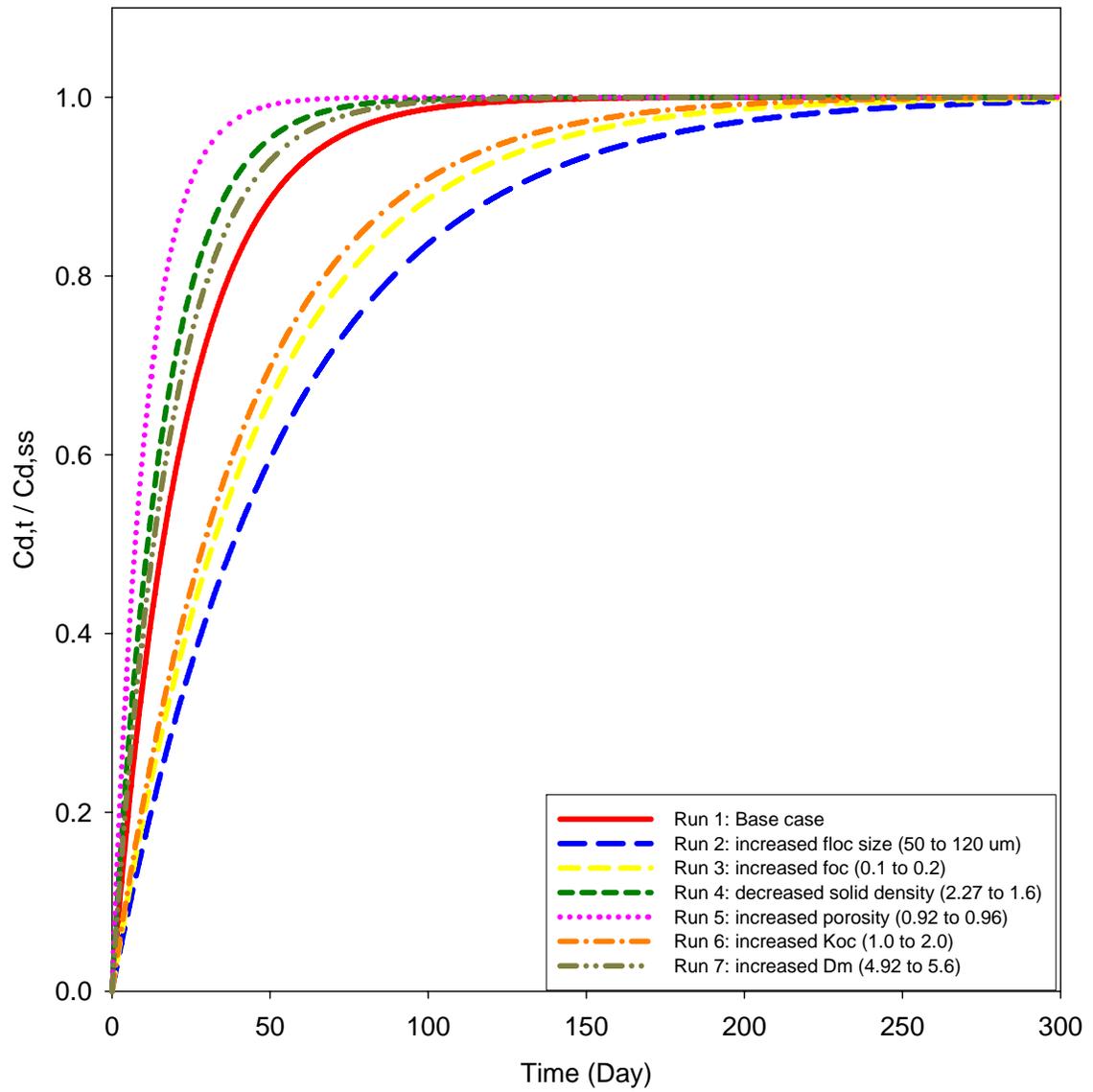


Figure 3.5

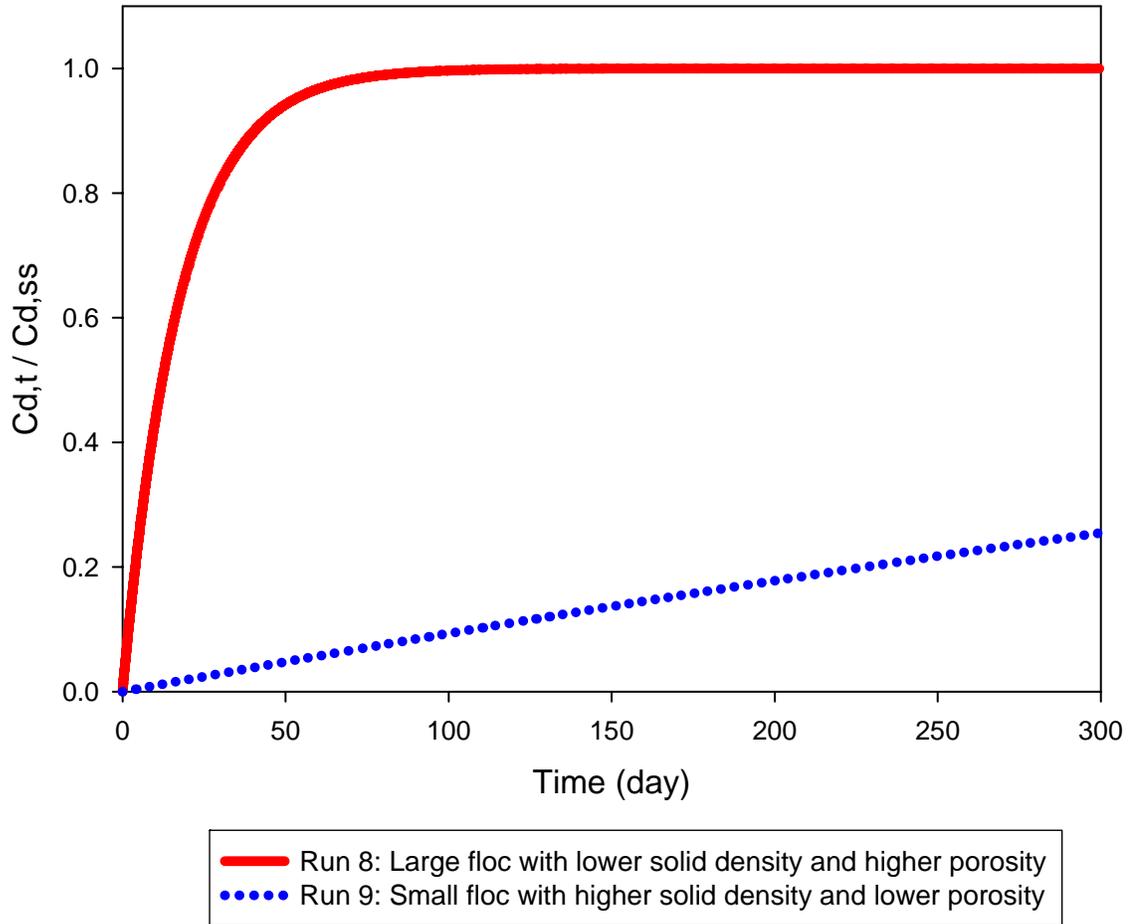


Figure 3.6:

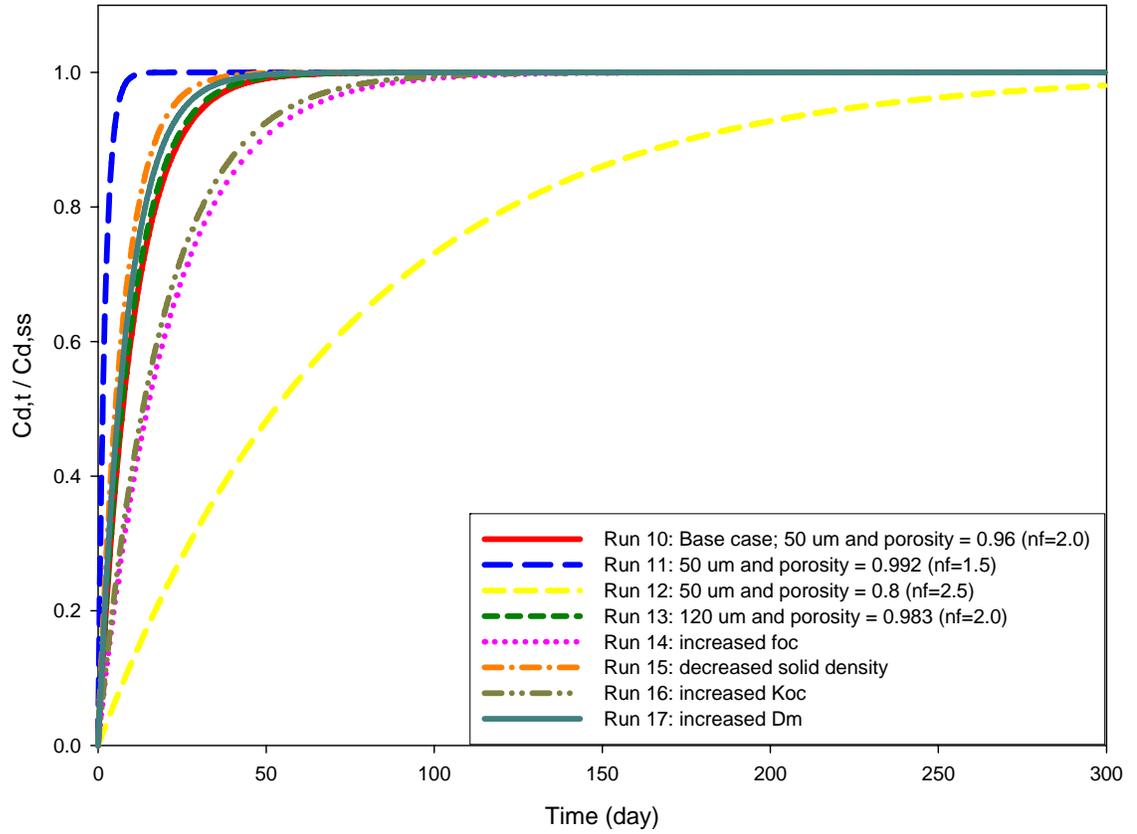


Figure 3.7:

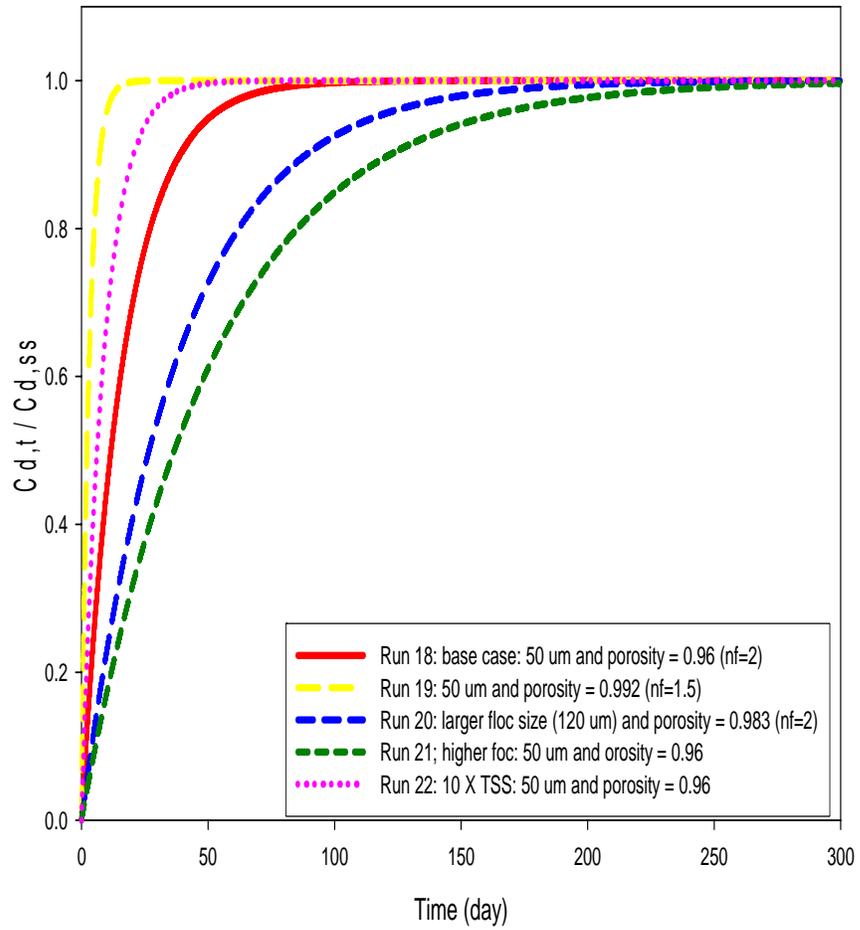


Figure 3.8:

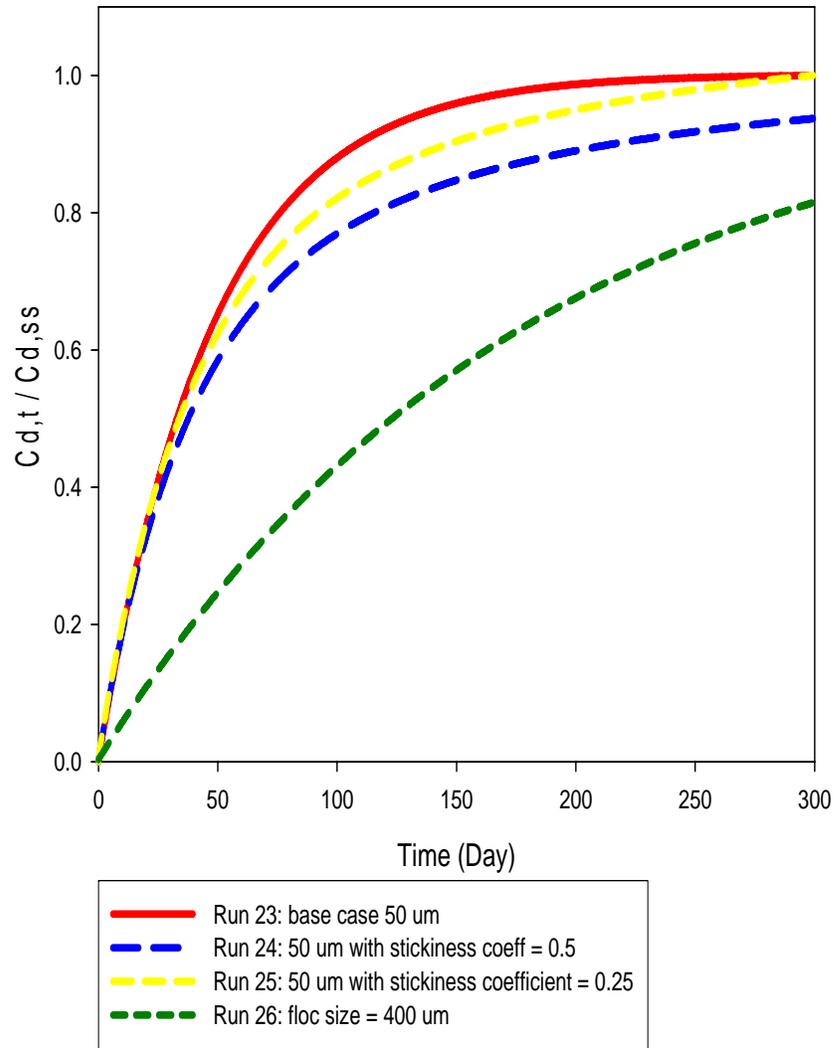


Figure 3.9:

Run 27: $R_{T, \text{with coagulation}}$ = 42 min.

Run 28: $R_{T, \text{without coagulation}}$ = 216 min.

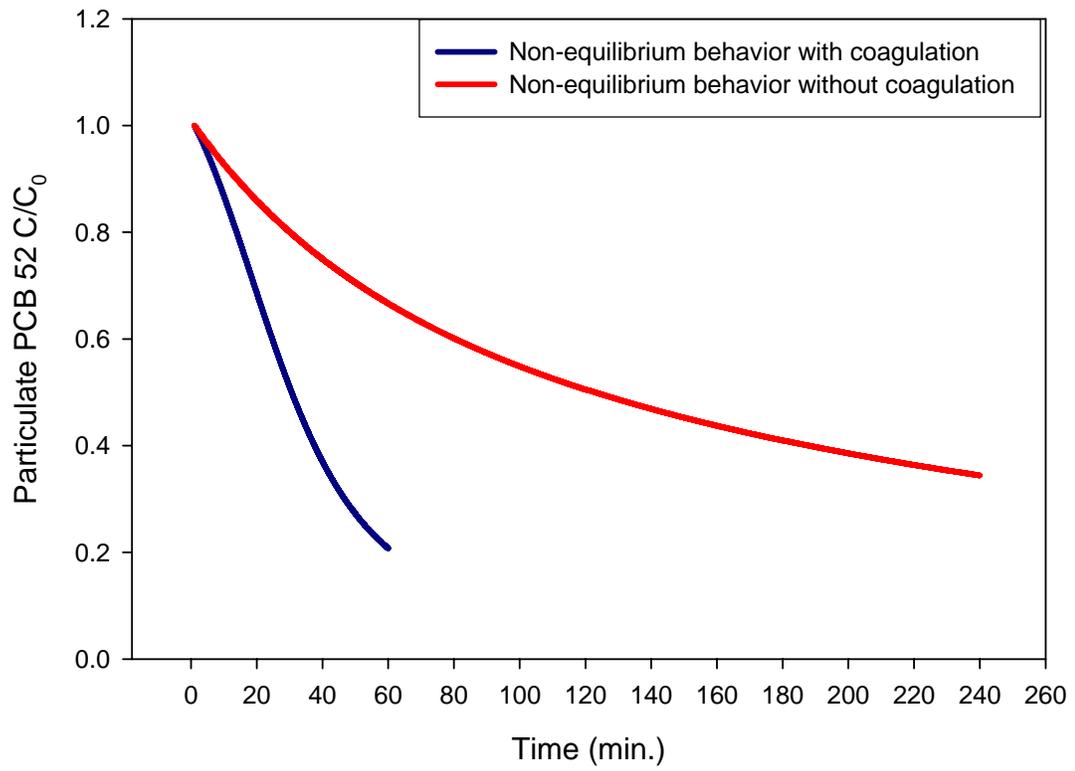


Figure 3.10:

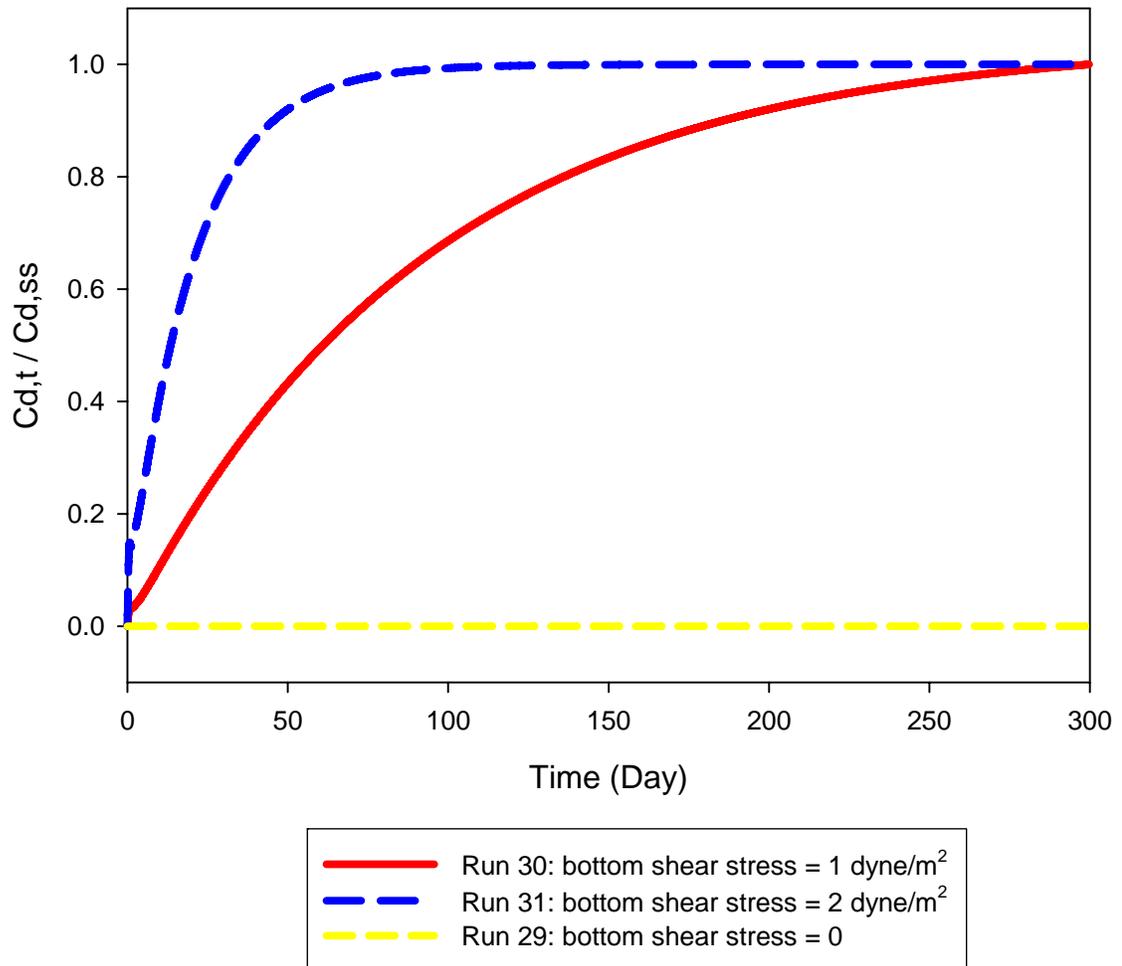


Figure 3.11

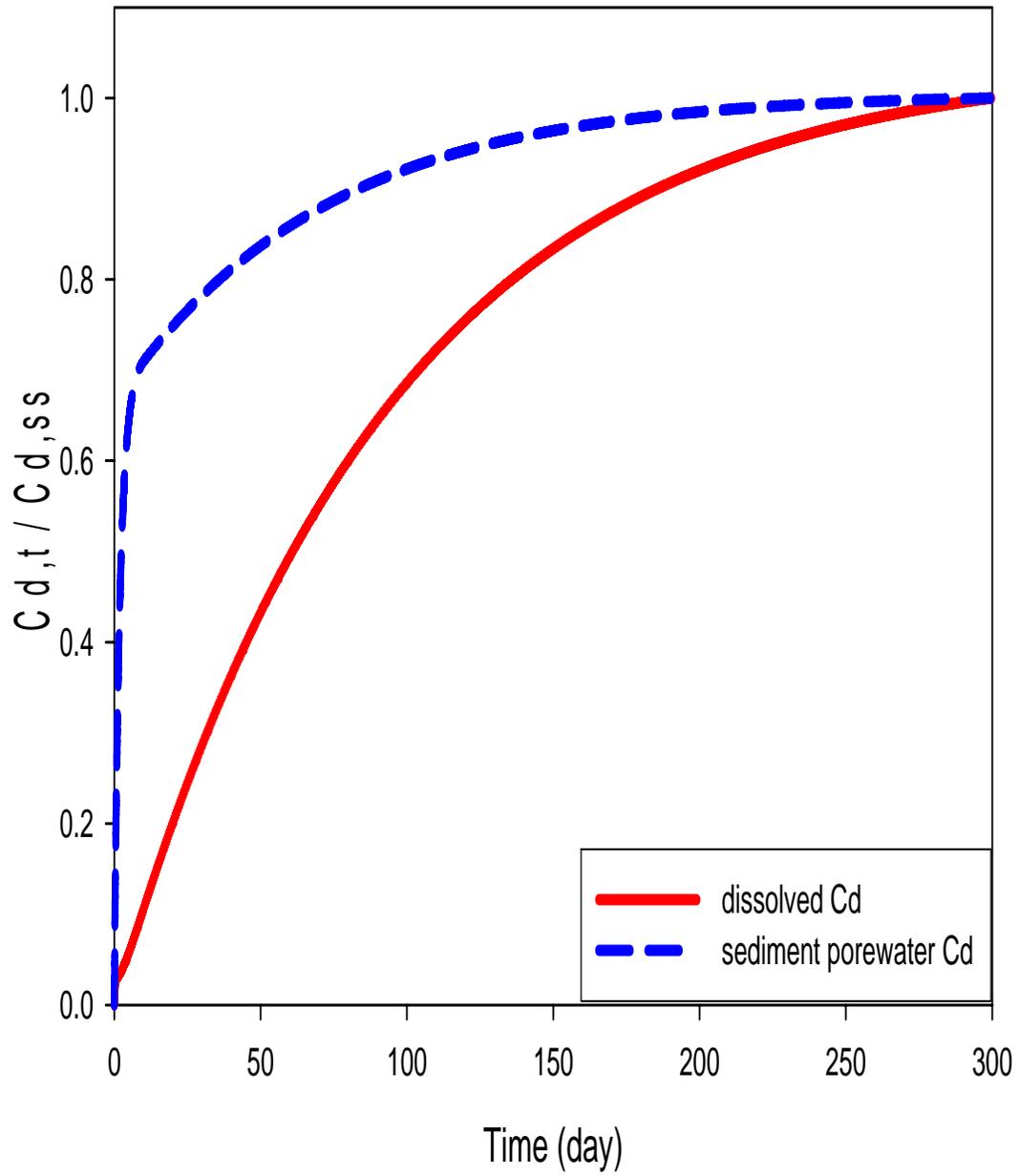


Figure 3.12:

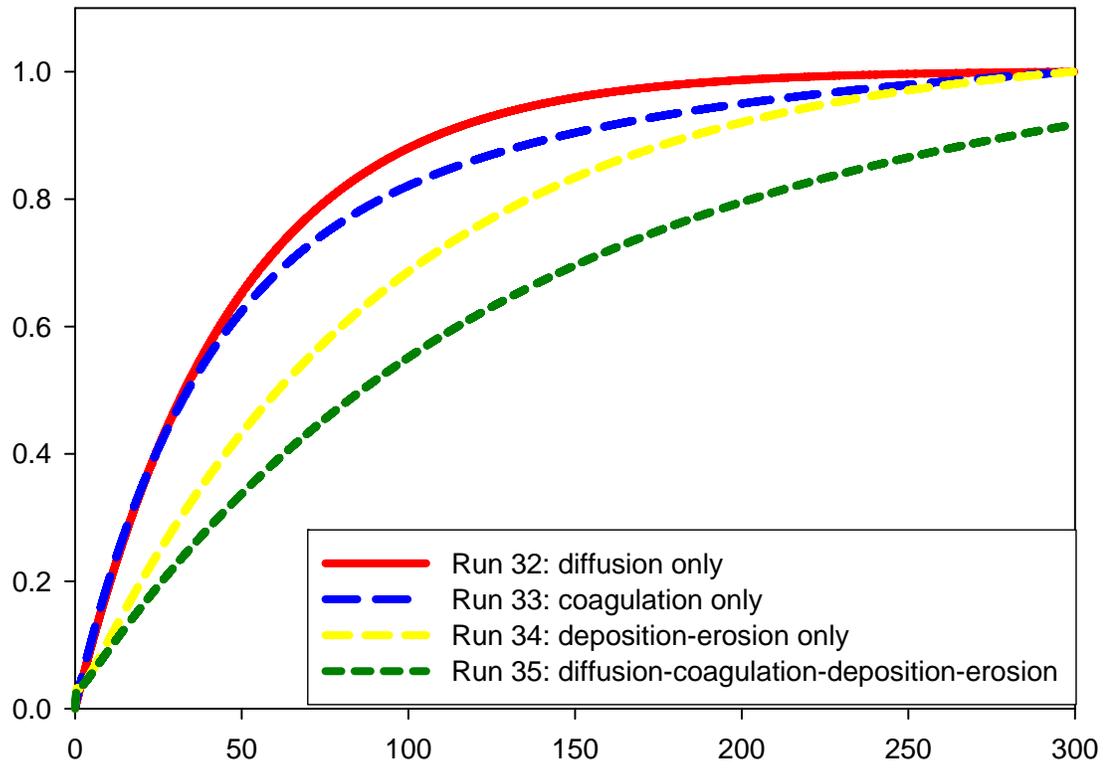


Figure 3.13:

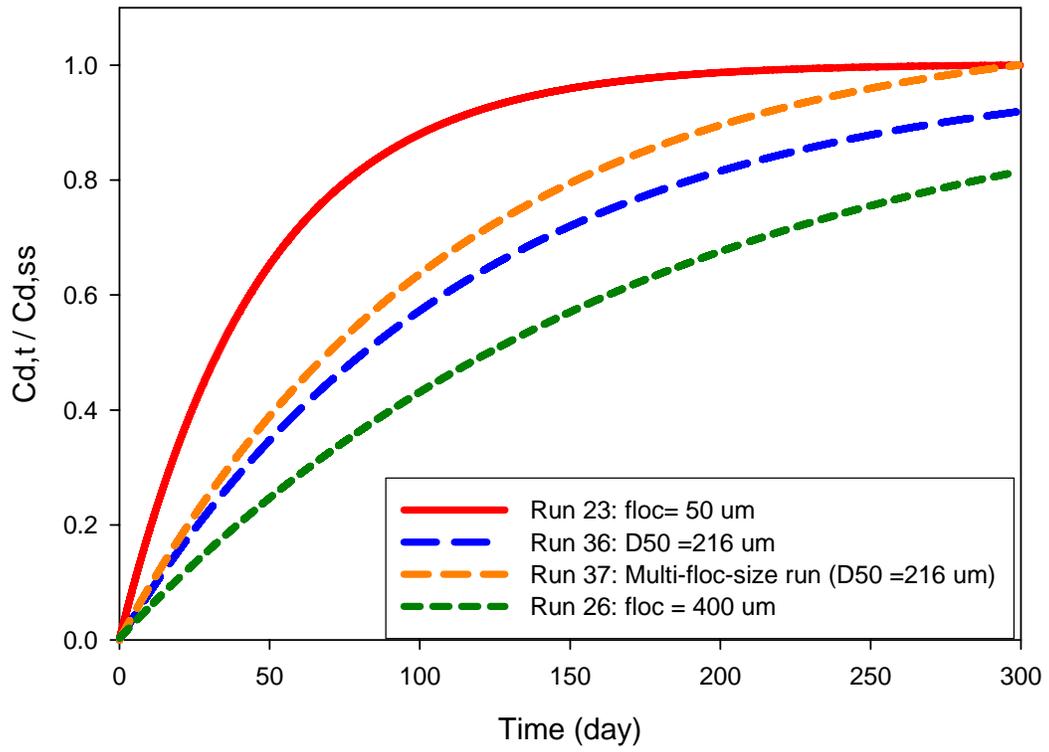


Figure 3.14:

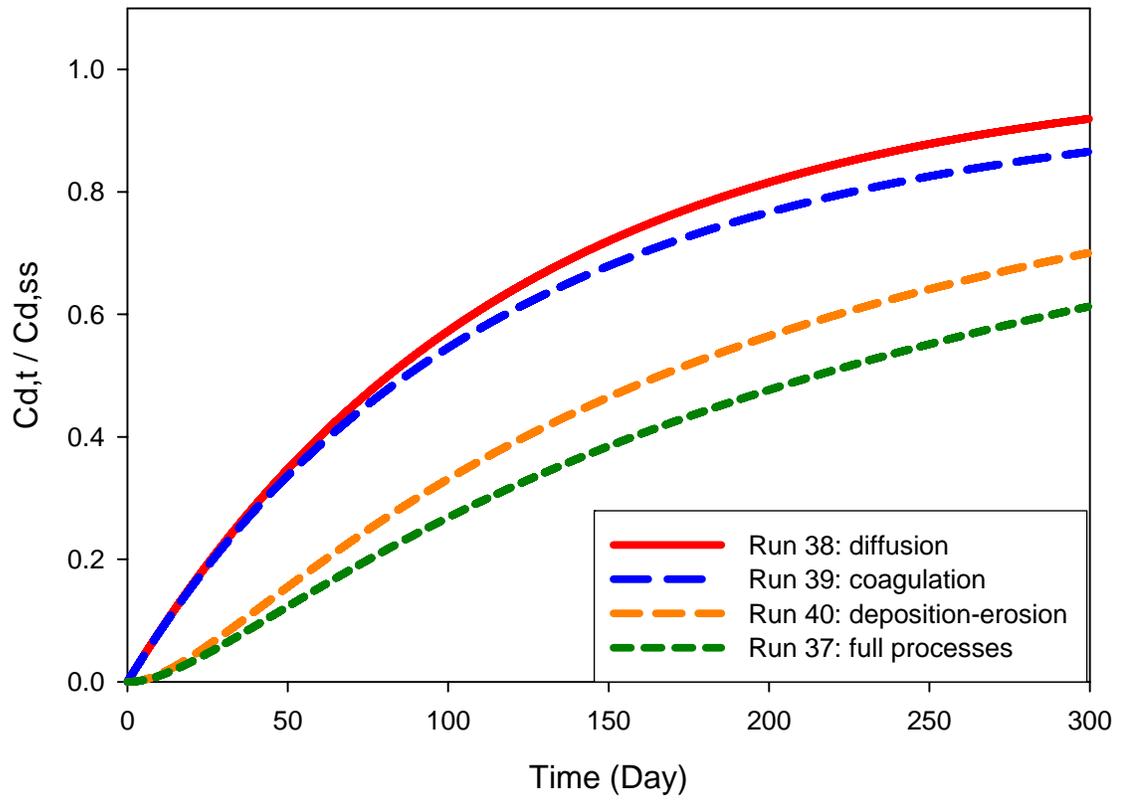


Figure 3.15:

PCB 4, 10

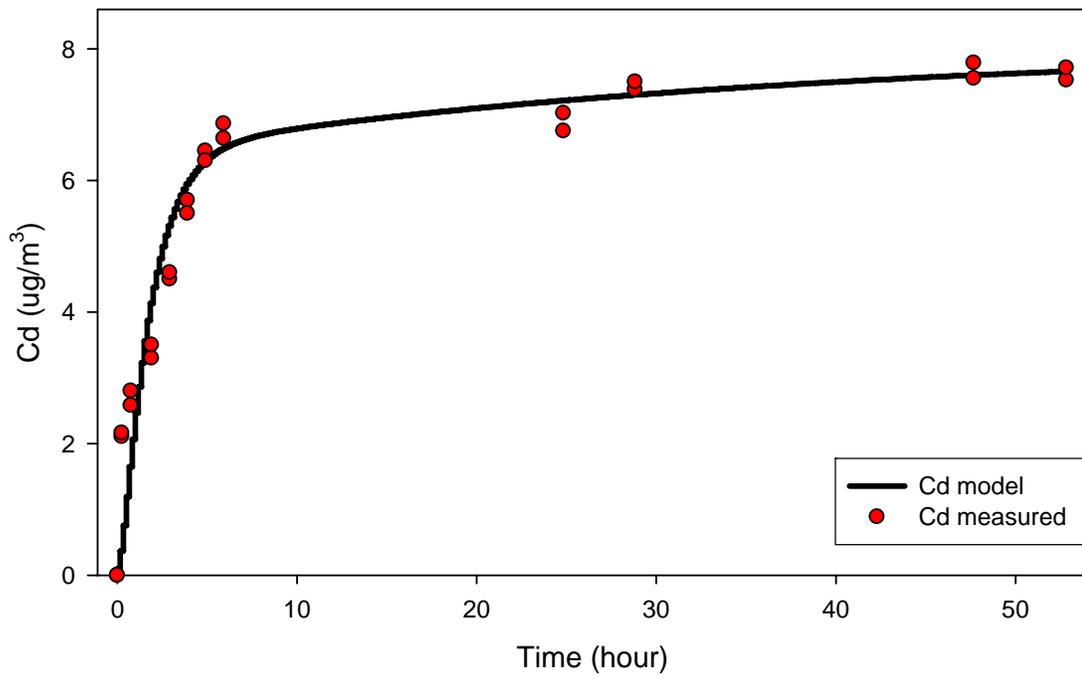
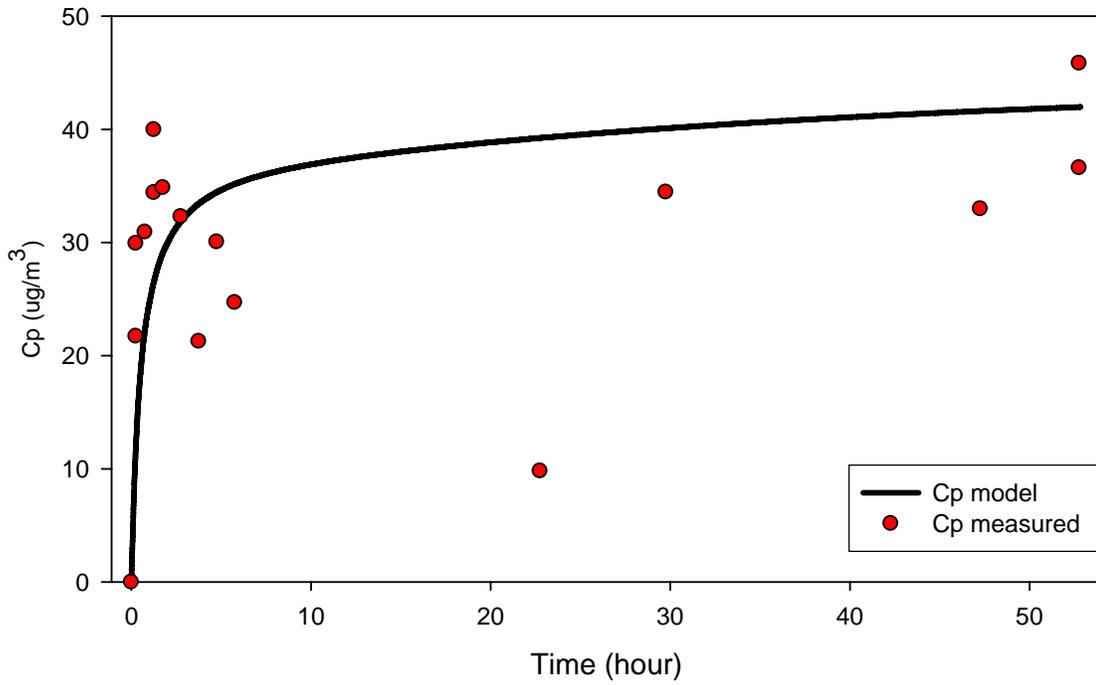


Figure 3.16:

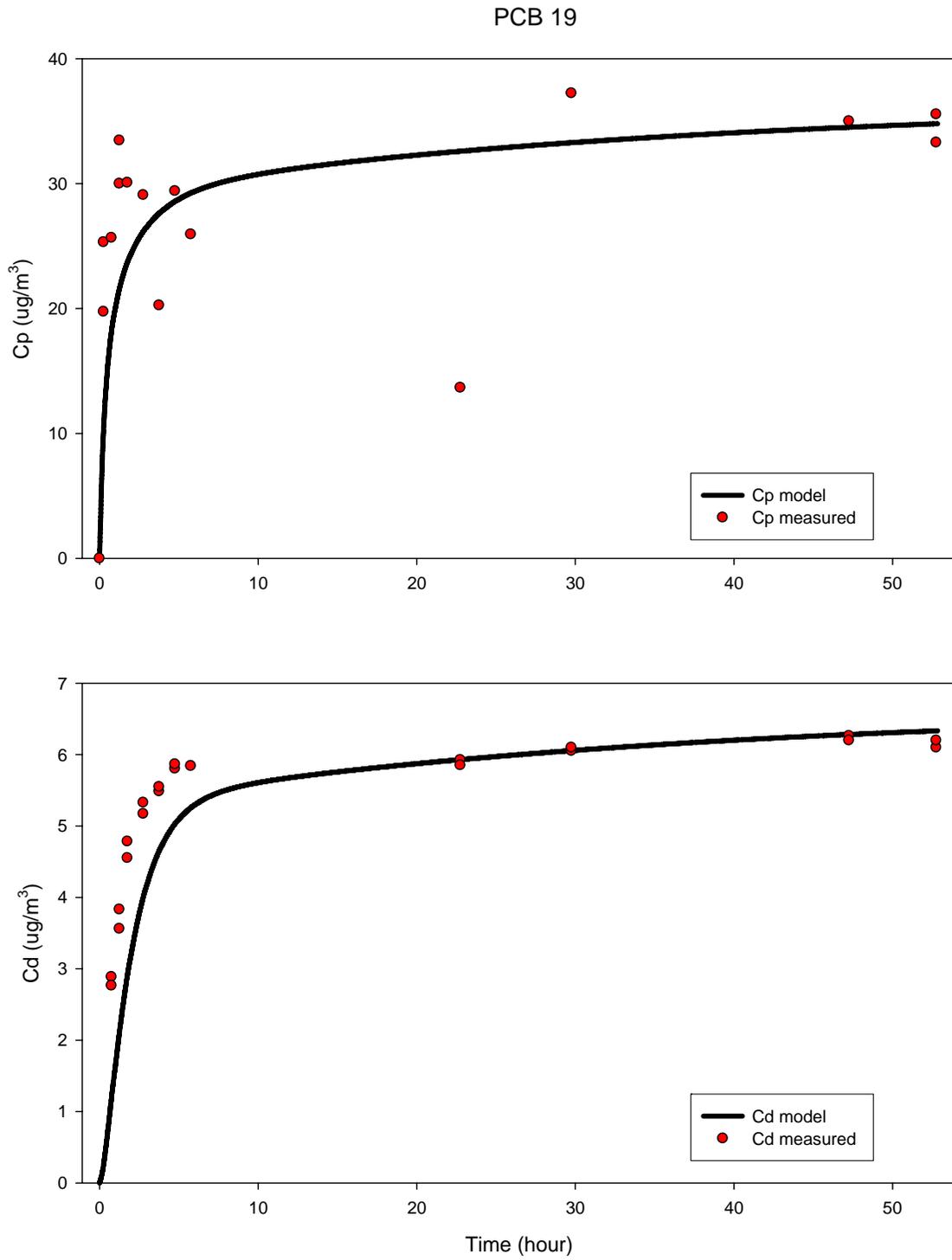


Figure 3.17:

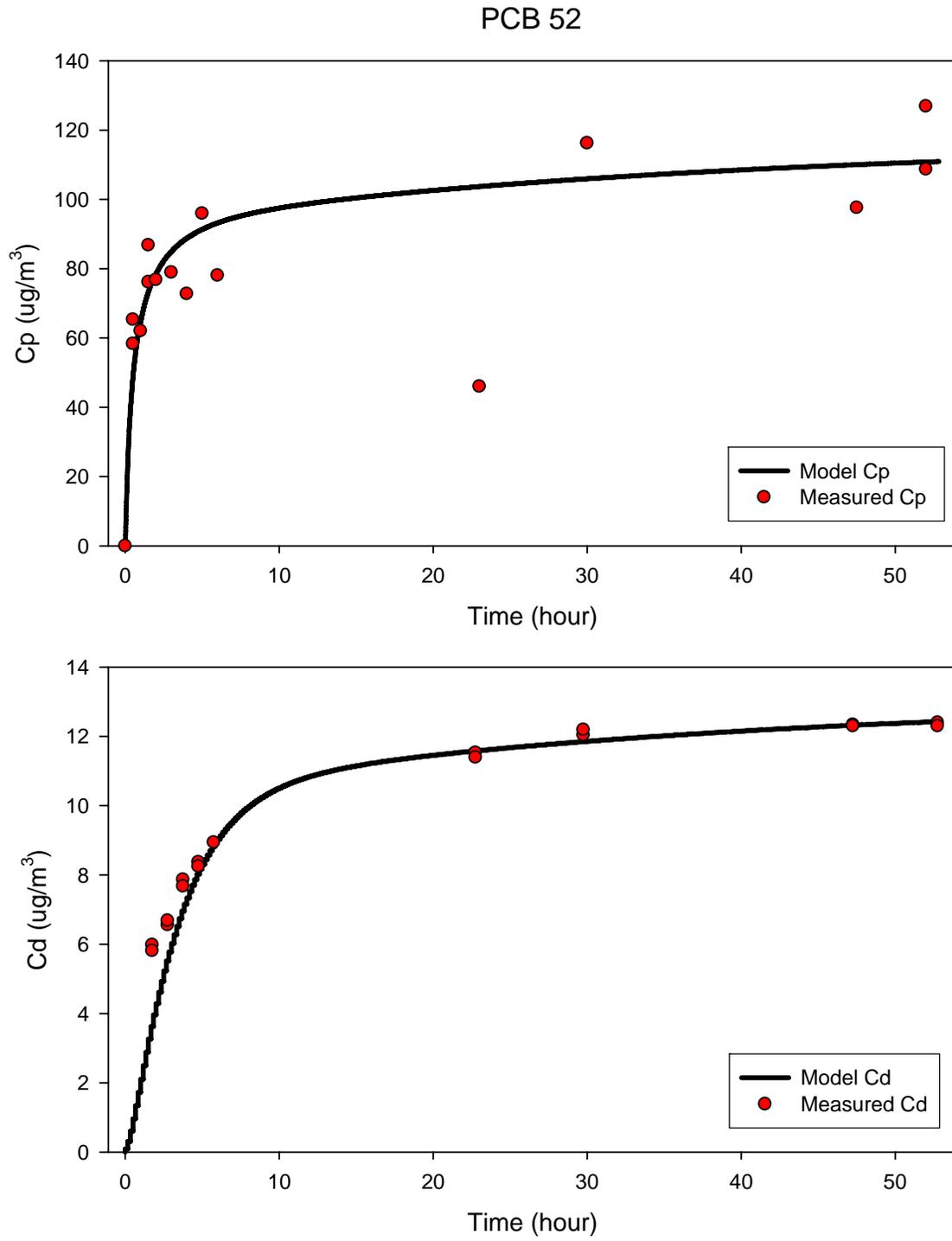


Figure 3.18:

PCB 110, 77

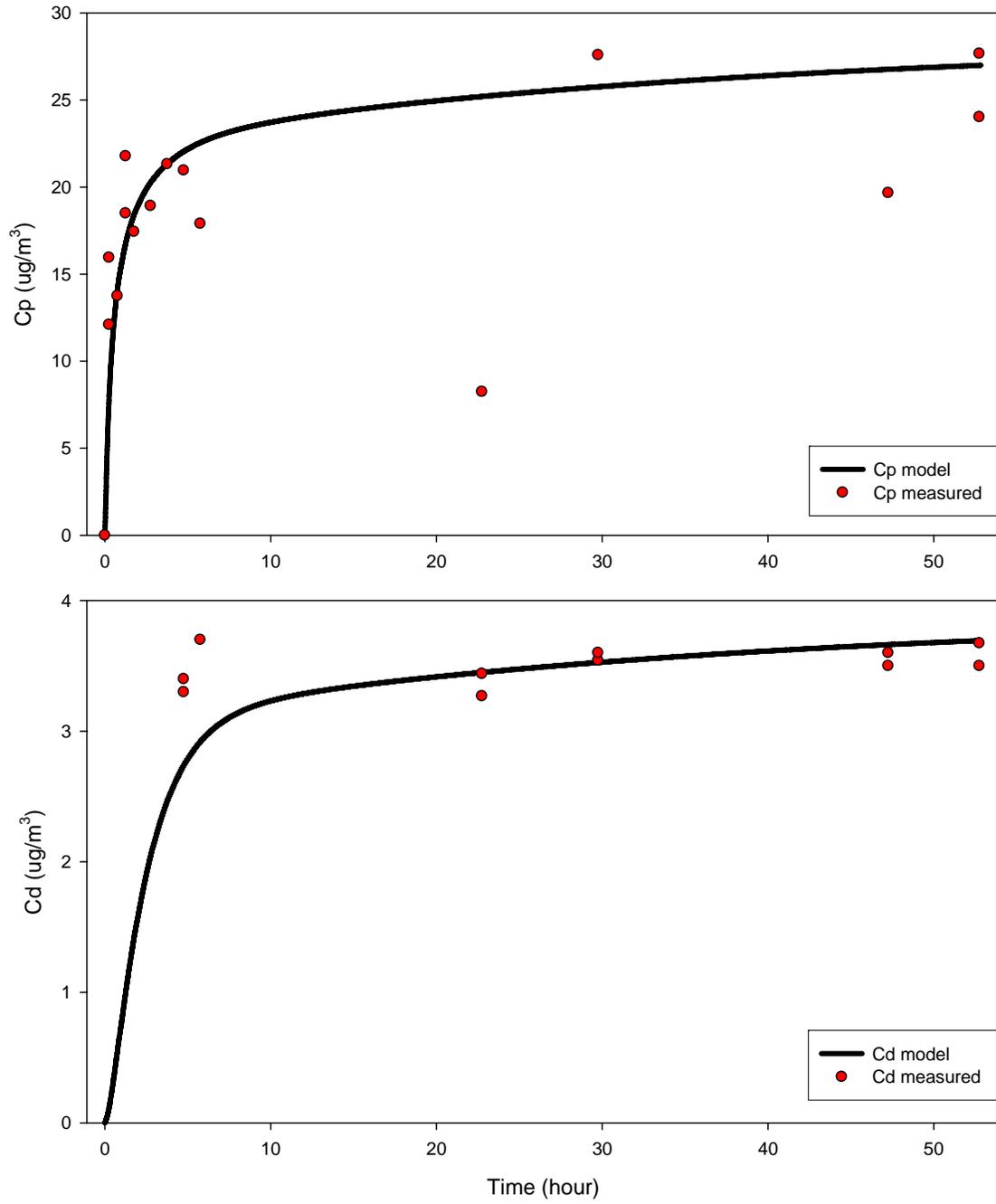


Figure 3.19:

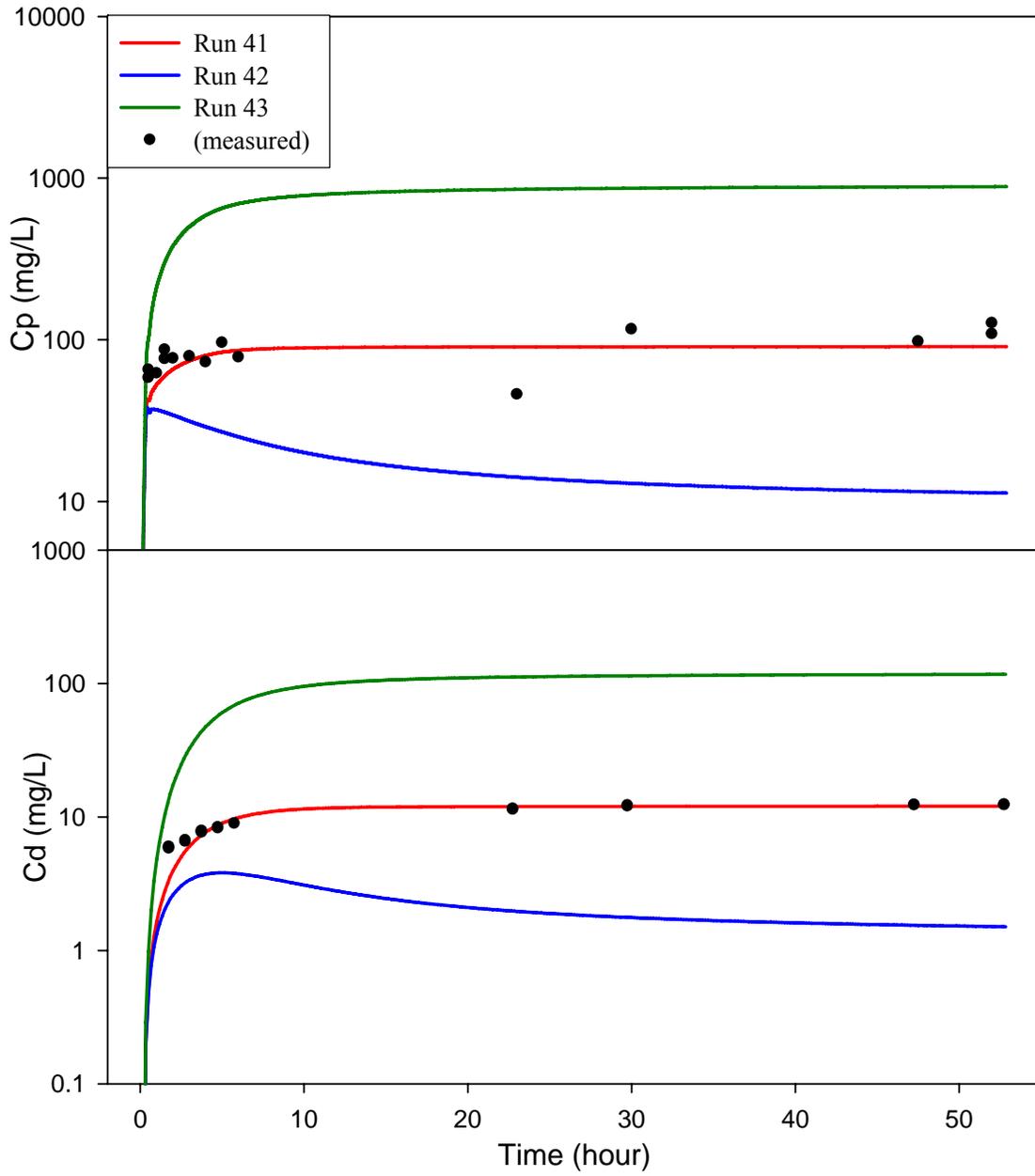
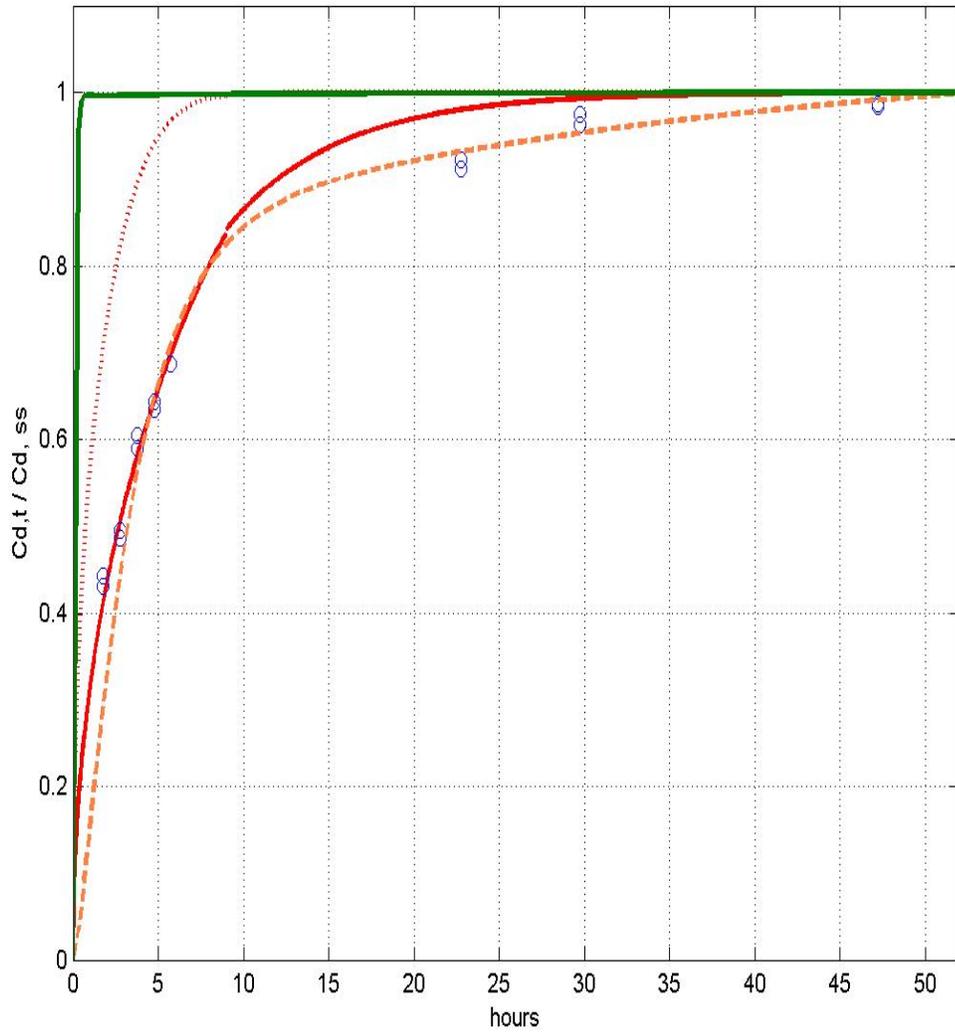


Figure 3.20:



- Run 44: Equilibrium Partition
- ⋯ Run 45: Radial diffusion behavior (before flocc size adjustment)
- Run 46: Radial diffusion behavior (after flocc size increased 70%)
- - - Run 47: New kinetic-limited behavior in this study
- STORM measurements

Figure 3.21:

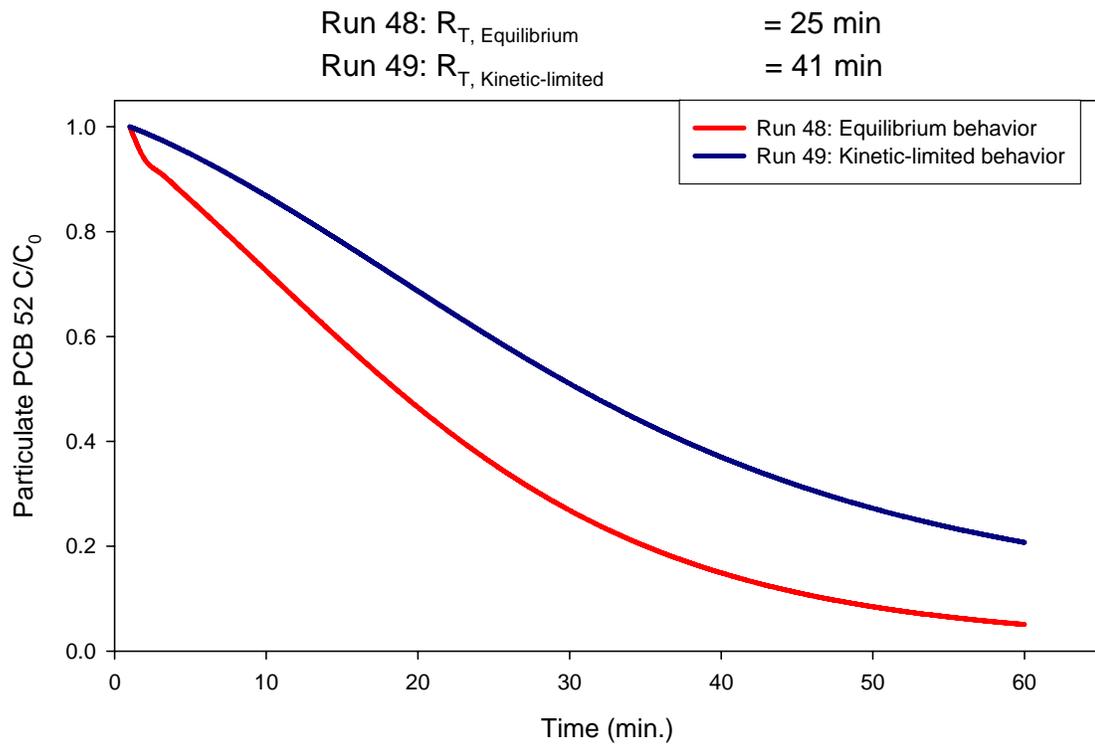


Table 3.1: The settings and the corresponding assumptions of Runs 1 to 9

Run	diameter (um)	TSS (g)	f_{oc}	solid density	porosity	$K_{oc}(m^2/g-OC)$	Dm	HOCPOC (ug/m ³)	N*A	MTV (m ³ /sec)
1	50	1	10.0%	2.27	0.92	1.1	4.92E-06	1.404	base case	3.90E-10
2	120	1	10.0%	2.27	0.92	1.1	4.92E-06	1.404	same as run1	1.62E-10
3	50	1	20.0%	2.27	0.92	1.1	4.92E-06	1.404	same as run1	1.95E-10
4	50	1	10.0%	1.6	0.92	1.1	4.92E-06	1.404	same as run1	5.53E-10
5	50	1	10.0%	2.27	0.96	1.1	4.92E-06	1.404	same as run1	8.49E-10
6	50	1	10.0%	2.27	0.92	2.0	4.92E-06	1.404	same as run1	2.14E-10
7	50	1	10.0%	2.27	0.92	1.1	5.60E-06	1.404	same as run1	4.75E-10
8	244	1	10.00%	1.05	0.97	1.1	4.92E-06	1.404	same as run1	5.12E-10
9	18	1	10.00%	2.2	0.25	1.1	4.92E-06	1.404	same as run1	8.80E-12

Table 3.2: The settings and the corresponding assumptions of Runs 10 to 17

Run	diameter (um)	TSS (g)	foc	solid density	porosity	mf	K _{oc} (m ³ /g-OC)	Dm	HOCPOC (ug/m ³)	N*A	MTV (ml/sec)
10	50	1	10.0%	2.27	0.960	2	1.1	4.92E-06	1.404	base case	8.49E-10
11	50	1	10.0%	2.27	0.992	1.5	1.1	4.92E-06	1.404	same as run10	4.53E-09
12	50	1	10.0%	2.27	0.800	2.5	1.1	4.92E-06	1.404	same as run10	1.18E-10
13	120	1	10.0%	2.27	0.983	2	1.1	4.92E-06	1.404	same as run10	8.90E-10
14	50	1	20.0%	2.27	0.960	2	1.1	4.92E-06	1.404	same as run10	4.24E-10
15	50	1	10.0%	1.6	0.960	2	1.1	4.92E-06	1.404	same as run10	1.20E-09
16	50	1	10.0%	2.27	0.960	2	2.0	4.92E-06	1.404	same as run10	4.67E-10
17	50	1	10.0%	2.27	0.960	2	1.1	5.60E-06	1.404	same as run10	1.03E-09

Figure 3.3: The settings and the corresponding assumptions of Runs 18 to 22

Run	diameter (μm)	TSS (g)	f_{oc}	solid density	porosity	nt	K_{oc} ($\text{m}^3/\text{g}\cdot\text{oc}$)	Dm	HOCPOC ($\mu\text{g}/\text{m}^3$)	NI	At	MTV (minsec)	MTV*NAF
18	50	1	10.0%	2.27	0.980	2	1.1	4.92E-06	1.404	4.10E+12	3.14E-10	8.27E-10	1.08E-06
19	50	1	10.0%	2.27	0.982	1.5	1.1	4.92E-06	1.404	2.05E+13	6.28E-11	4.41E-09	5.68E-06
20	120	1	10.0%	2.27	0.983	2	1.1	4.92E-06	1.404	7.12E+11	7.54E-10	8.68E-10	4.65E-07
21	50	1	20.0%	2.01	0.960	2	1.1	4.92E-06	1.404	4.10E+12	3.14E-10	4.79E-10	6.17E-07
22	50	10	10.0%	2.27	0.980	2	1.1	4.92E-06	1.404	4.10E+13	3.14E-10	8.27E-10	1.08E-05

Table 3.4: Model simulation results in Runs 23 to 26

run	OC _{ss} (mg/L)	TSS _{ss} (mg/L)	Cp OC (ng/L)	Cd (ng/L)	tHOCwater (ng/L)	OCSS (hr.)	Cp OC (hr.)
23	9.07	89.48	118.00	11.78	129.8	12.0	13.3
24	8.41	77.90	110.94	11.86	122.8	10.0	8.5
25	5.58	38.58	90.17	14.26	104.4	9.5	2.2
26	23.32	230.01	306.34	11.90	318.2	15.0	13.2

Chapter 4: Predicting the Behavior of PCBs in Activated Carbon-Amended Sediments

4.1 Abstract

Traditional contaminated sediment clean-up methods are often complex and demonstrate a poor understanding of the relationship between sediment geochemistry and bio-availability (Ghosh *et al.*, 2003). As an alternative, several investigators have proposed adding granular activated carbon (AC) *in situ* (e.g., Kosian *et al.*, 1999; Lebo *et al.*, 2003; Zimmerman *et al.*, 2005). AC has a stronger sorption capacity for hydrophobic organic contaminants (HOCs) than do flocs of the same size and organic carbon content. However, organic contaminants sorb more slowly to AC (i.e., have smaller mass transfer velocity) due to the relatively low porosity and higher partitioning coefficient (K_{AC}). Over long time scales, a majority of HOCs will be sequestered in the activated carbon. However, in the absence of flocculation in a dynamic erosional and depositional environment, the initial HOC distribution in the sediment and the relatively slower sorption rate for activated carbon are the dominant factors controlling the distribution of HOCs in the water column. The interaction time becomes a very important issue under conditions when activated carbon does not have enough time to adsorb or compete with other organic particles. The above analysis considers the case where amended activated carbon does not interact with natural sediments or other aquatic particles to form flocs. We found that the lower porosity and higher settling velocities of activated carbon

relative to natural organic carbon result in lower steady-state activated carbon concentrations in the water column and longer times to equilibrium.

In AC-amended sediment, the total water column HOC concentration is significantly lower compared to the AC free run. Further, when AC aggregates with sediment particles, this flocculated AC has the same physical properties as the floc, resulting in a slower settling velocity, longer residence time, and higher AC-associated HOC concentration in the water column. When activated carbon is added to contaminated sediments, the total PCB concentration in the water column decreases by 90% (123.4 to 11.4 ng/L). If the activated carbon coagulates with the resuspended sediment, this decrease is partially offset by some activated carbon being entrained in slowly-settling flocs, and the steady-state PCB concentration is 61 ng/L

4.2 Introduction

Hydrophobic organic contaminants (HOCs) are important pollutants in urban estuarine sediments. HOCs include polycyclic aromatic hydrocarbon (PAH) and polychlorinated biphenyls (PCBs). Concerns arise due to their toxicity and potential carcinogenicity to humans. Therefore, PCBs production was banned in the mid 1970s (NYS DH, 1998). The distribution and fate of HOCs are highly correlated to their sorptive behavior (Karickhoff *et al.*, 1979; Wu and Gschwend, 1986). Sorption to resuspended particles and sediments plays an important role controlling the water column residence times and spatial distributions of HOC in aquatic environments. HOCs residence times and the time required to reach sorptive equilibrium are highly dependent on the chemical character, the surrounding environment, and particle types and compositions (Schnoor, 1996).

Therefore, in 1993 EPA evaluated *in situ* remediation techniques for contaminated sediment (EPA, 1993). Two types of *in situ* contaminated sediment treatments have been conducted in the field, namely biological/chemical treatment methods and solidification/stabilization treatment methods. *In situ* biological/chemical treatment involves the addition of microorganisms and/or chemicals to the sediments to initiate or enhance bioremediation. *In situ* solidification/stabilization treatment involves the addition of chemicals or cements, such as Portland cement, to encapsulate the contaminated sediments and convert them into less soluble or mobile forms (e.g., Murphy *et al.*, 1995; Chowdhury *et al.*, 1996).

However, traditional *in situ* contaminated sediment clean-up methods are often complex and demonstrate a poor understanding of the relationship between sediment geochemistry and bio-availability (Ghosh *et al.*, 2003). As an alternative to traditional methods of cleaning up sites with contaminated sediments, several investigators have proposed adding granular activated carbon *in situ* (e.g., Kosian *et al.*, 1999; Lebo *et al.*, 2003; Zimmerman *et al.*, 2005). This added carbon effectively sequesters PCBs, reducing bioavailability and therefore the level of risk associated with the contaminated sediments (Zimmerman *et al.*, 2004). Zimmerman and colleagues (2005) mixed 3.4% activated carbon with contaminated sediment at Hunters Point, San Francisco Bay, reducing the bio-availability of PCBs by 82% for worms and 70% for amphipods.

However, these methods have been demonstrated mainly in laboratory experiments. They assume AC either was well mixed with sediment suspending in the experimental chamber, or simply remained on the chamber bottom (Zimmerman *et al.*, 2005; Sun and Ghosh, 2007). These experiment designs left significant operational questions remain

prior to their application in the field on a large-scale. Of primary importance is the fate of the activated carbon particles in the sediments. Are they stable? How do we tell the contribution between AC and OC on HOC remediation, which is difficult to measure in a dynamic environment (Simpson and Hatcher, 2004)? Do they aggregate with sediment particles? Are they susceptible to resuspension? If they will resuspension, do HOCs concentrations higher in the resuspended AC than other resuspended organic carbon (OC)? Do they have a different residence time and steady state concentration than OC in the water column? Further, during the limited interaction time, do AC have the same remediation effect?

To address these questions, the model developed in the previous chapters is modified to include activated carbon as a state variable. Scenarios are examined in which the activated carbon particles are resuspended and aggregate with eroded sediments to explore the long-term performance of the *in situ* remediation.

4.3 Objectives

The overall objective of this chapter is to use the model to explore the behavior of activated carbon when added to sediments as an *in situ* remediation technology.

4.4 Strategies

In this application, the previously developed dynamic sediment-water exchange/ flocculation/ HOC partitioning model was modified to include activated carbon as a state variable (Figure 4.1 and 4.2). To simplify model parameters and environmental settings, the same STORM experiment conditions were used as in prior model development (Chapters 2 and 3). Activated carbon was added to the model in two stages, first insuring

that the additional state variable behaves as expected in isolation and does not interfere with the previously calibrated parameters, then allowing for interactions (flocculation and sorption competition) among activated carbon and sediments. This enhanced model is then used to explore the influence of AC on the fate of HOC in the water column, when activated carbon interacts with natural sediments to form flocculated particles.

First, activated carbon particles with specific properties were added as a state variable to the model. Model performance was verified by first focusing on the behavior of activated carbon (AC) and organic carbon (OC) independently, including deposition to and resuspension from the sediment bed and PCB diffusion into and from the particles. In this stage, the model was used to examine the competitive interactions between AC and OC for PCB partitioning, without allowing for flocculation.

Second, after demonstrating model performance with the additional AC state variable, we then allowed AC to coagulate with resuspended sediment particles, forming mixed AC-OC aggregates. The model explicitly calculates the size-specific physical properties (bulk density, porosity, settling velocity, and stickiness coefficient) of these mixed aggregates. The impact of these aggregates on PCB concentrations and speciation in the water column was examined.

4.5 Properties of Suspended Solids, Organic Carbon, and Activated Carbon Used in These Simulations

The organic carbon (OC) and bulk suspended solids (TSS) have the same properties as in the previous chapter, which are controlled by the fractal factor, the nature of the primary particle, f_{OC} , and floc diameter. Properties of the activated carbon (AC) are D_{50} is

150 μm , porosity is 0.55, f_{OC} is 0.44, and K_{AC} is 1000 times larger than K_{OC} (Zimmerman *et al.*, 2005).

4.6 *Stage one: Model Evaluation with Activated Carbon in the Absence of Flocculation*

4.6.1 Model Scenarios and Settings

Eight model runs were conducted to systematically add and evaluate activated carbon to the model (Table 4.1). Runs 1.1 to 1.3 compared rates and times to equilibrium of PCB diffusion into AC and OC for a suspension (*i.e.*, no erosion or settling). Runs 1.4 and 1.5 tested the impact of deposition-resuspension on the steady state solid concentration and dissolved and particle HOC concentrations when either OC or AC is present in the sediments and water column. Here, AC and OC of equal particle size are assumed to be resuspended at the same rate when under the same bottom shear stress.

Runs 1.6 to 1.8 examined the HOC sorption competition between AC and OC and the impact of the initial distribution of HOC between OC and AC in the sediment.

4.6.2 Results and Discussions

Before adding activated carbon to the model, we verified that the existing model gave identical results as previous runs (Chapter 2) when using the same initial conditions.

Activated carbon was added as a state variable to the model, and a run was made where all of the solids in the slurry were AC, at the same concentration as OC in the base case run. This run reached the same steady state HOC concentration in the dissolved and particle phases when the AC and OC partition coefficients were set to be equal. The only difference between these two runs was the time to reach steady state. As expected, the

run with AC took longer to reach steady state than the OC-only run because AC is less porous than the same size OC flocs. The diffusion rate is controlled by the total contact area and mass transfer velocity. In run 1.3, the AC partition coefficient was increased to 1000 times that of OC. As a result, the steady state HOC concentration increased and the time to reach steady state was much longer because increasing K_{AC} significantly decreases the mass transfer velocity. In short, AC has a stronger HOC capacity than the same size OC because of the difference of partition coefficient. However, AC requires a longer time to reach steady state or equilibrium than the same size of flocs. These results could affect the HOC fate in a dynamic environment.

Run 1.4 to 1.8 expanded the simulations above to include settling and resuspension, but still did not include flocculation. AC and OC were resuspended at equal rates in these runs. Runs 1.4 and 1.5 differ in how the carbon and HOC are initially distributed in the sediment (100% in OC in run 1.4 *versus* 100% in AC in run 1.5). Under these conditions with settling and resuspension, AC reached a lower steady state mass concentration in the water more quickly than OC because it has a faster settling velocity (higher bulk density than OC flocs for a given particle size). Further, AC has a stronger sorption capacity and a slower sorption rate than the same size of OC. As a result, the resuspended carbon concentration decreases 99.3%, and total water column HOC concentration decreases 99.5% in run 1.5 compared to run 1.4.

Run 1.6 assumed AC and OC were each 50% of the initial sediment total carbon content and sediment HOC reached equilibrium status between AC and OC as the initial condition. The steady state mass concentration of AC and OC are the same whether beginning with only one or the other in the sediments (runs 1.4 and 1.5) or both (run 1.6),

verifying the particles are acting independently in this version of the model. Furthermore, although AC was responsible for 0.7% of total steady state water column carbon in run 1.6, the total water column HOC in run 1.6 decreased 98.7% compared with run 1.4.

Runs 1.6 to 1.8 compared total water column HOC concentration under three different initial sediment HOC distributions. All runs were simulated based on the same solids transport. As shown in Table 1, with limited contact time between the HOC and the AC, the initial sediment HOC in AC was lower than the equilibrium status. Under these initial conditions, the total water column HOC concentrations are 76 times higher than resuspension events where the HOCs had enough time to adsorb. These results suggest that the interaction time is a very important issue under conditions when AC does not have enough time to adsorb or compete with other organic particles.

The steady state concentrations of particulate (C_p) and dissolved (C_d) PCB 52 are much higher when OC sediments are resuspended (run 1.4) than when the PCBs begin sorbed to activated carbon in the sediment (run 1.5). Total (dissolved plus particulate) PCB congener 52 concentrations are higher when OC is resuspended compared to AC resuspension (225 to 1.2 ng/m^3) for two reasons. First, the steady-state OC concentration in the water column is higher due to slower settling velocities. Both runs start with the same initial sediment HOC concentration and resuspension rate, and therefore the PCB gross erosion fluxes are the same. However, the lower settling velocity of OC relative to AC results in more HOC in the water column. The dissolved HOC concentration, which is proportional to the particulate HOC concentration, is also higher in the OC run. Finally, since the AC mass transfer velocity is substantially smaller than that of OC, less

HOC is released to the dissolved phase as a result of AC resuspension (more time is required to desorb into the water column).

The objective of runs 1.6 to 1.8 is to examine the impact of the initial sediment HOC concentration distribution on the fate of dissolved HOC. All three runs were based on the same solid transport. Run 6 assumed sediment HOC ($\text{ng-PCB } 52/\text{m}^3$) was evenly distributed between AC and OC in the sediment layer. In contrast, the sediment HOC was allowed to reach equilibrium between sediment porewater, sediment C_{pOC} , and sediment C_{pAC} (run 1.8). An intermediate initial condition, in which the 75% of the HOC was initially associated with the AC, was also explored (run 1.7). The total HOC mass and the particle transport rates are the same among these runs, as verified by identical trends in AC and OC concentrations. However, because AC has a higher partition coefficient than OC, most HOC is associated with sediment AC rather than sediment OC. Combining the effects of the initial sediment HOC distribution and concentrations of resuspended solids, the dissolved and OC-associated PCB concentrations are much lower in run 1.8 than run 1.6, as the bulk of the PCB remains bound to activated carbon. In run 7, in which the HOC initial distribution was intermediate, the C_d , C_{pOC} , C_{pAC} values are between those of run 1.6 and run 1.8.

AC has a stronger sorption capacity than flocs of the same size and organic carbon content. However, AC has a slower sorption rate (smaller mass transfer velocity) due to its relatively lower porosity and higher K_{AC} value. Over long time scales, a majority of HOCs will be sequestered in the activated carbon. However, in the absence of flocculation in a dynamic erosional and depositional environment, the initial HOC distribution and the relatively slower sorption rate for AC are the dominant factors

controlling the distribution of HOCs in the water column. The interaction time becomes a very important issue under conditions when AC doesn't have enough time to adsorb or compete with other organic particles.

The above analysis considers the case where amended activated carbon does not interact with natural sediments or other aquatic particles to form flocs. We found that the lower porosity and higher settling velocities of AC relative to natural OC result in lower steady-state AC concentrations in the water column and longer times to sorptive equilibrium. In the next section, we will examine what happens when these AC particles become part of larger, more porous, less dense flocs.

4.7 Stage Two: Model Evaluation with Added Activated Carbon and Flocculation

4.7.1 Stage Two Descriptions and Settings

Beginning with the model described above, we now allow flocculation to occur between activated carbon particles and the natural OC. We assume that activated carbon (AC), organic carbon (OC), and inorganic solids (INS) are components of flocs with settling velocities and porosity determined by the relative contribution of the three particle types. Floc porosity is a function of the fractal factor, the nature of the primary particle, and the floc diameter, which means that the AC content has no impact on the floc porosity for any given size of flocs. The total suspended solids concentration (TSS) is the sum of INS, OC and AC, and the weighted fractional organic carbon ($f_{OC, total}$) equals $(OC+AC)/TSS$. Both the floc solid density and the stickiness coefficient are calculated as linear functions of f_{OC} and f_{AC} (Equation 4.1-4.3).

$$\rho_{dry,i} = 2 \times f_{OC,i} \times \rho_{biotic} + (1 - 2 \times f_{OC,i} - f_{AC,i}) \times \rho_{clay} + f_{AC,i} \times \rho_{AC} \quad \text{Equation 4.1}$$

$$\rho_{bulk,i} = \rho_{dry,i} \times (1 - \phi_i) + \rho_{water} \times \phi_i \quad \text{Equation 4.2}$$

$$\alpha_i = 2 \times f_{OC,i} \times \alpha_{biotic} + (1 - 2 \times f_{OC,i} - f_{AC,i}) \times \alpha_{clay} + f_{AC,i} \times \alpha_{AC} \quad \text{Equation 4.3}$$

where $\rho_{dry,i}$ is dry floc density (g/m^3), $\rho_{bulk,i}$ is bulk floc density (g/m^3), α_i is floc stickiness coefficient for floc at size i , f_{OC} and f_{AC} are fraction of organic carbon and activated carbon for floc at size i , ρ_{biotic} , ρ_{clay} , ρ_{AC} are organic carbon, inorganic carbon, and activated carbon dry density (g/m^3), ρ_{water} is water density (g/m^3), ϕ_i is floc porosity for floc at size i , and α_{biotic} , α_{clay} , α_{AC} are organic carbon, inorganic carbon, and activated carbon stickiness coefficient respectively.

4.7.2 Model Scenarios

Seven model scenarios were created to successively explore the interaction of AC and OC and its impact of HOC partitioning (Table 4.2). The model began with the same settings and initial conditions as described above, and PCB 52 was used as the target contaminant. In run 2.2 to 2.7, the model assumed natural organic carbon and AC were each responsible for 50% of total organic carbon.

1. Run 2.1: In this ‘base case’ run, the model started with the same settings as described above and did not contain activated carbon. The model included flocculation and resuspension-deposition processes.
2. Run 2.2 and 2.3: AC is added as a model variable, and the model had the same resuspension flux rate as run 1 for total organic carbon (TOC = sum of OC and AC) and TSS. As an initial condition, the AC has a diameter of 150 μm and a mass concentration equal to the sum of organic substrates for all size of flocs. The floc properties then vary temporally with the composition of AC, OC, and INS.

The model is initialized with HOCs in the sediment reaching equilibrium among AC, OC, and porewater, which was the result after a 100 year simulation (Appendix 1). Runs 2.2 and 2.3 compare two different settling velocity calculations (Stokes law equation *versus* fractal geometry adjusted settling velocity equation) for activated carbon in the absence of flocculation.

3. Run 2.4: Model starts with the same physical and chemical conditions as run 2.3. However, in this run OC is allowed to coagulate while flocculation of AC is prevented.
4. Run 2.5: The model begins with the same physical and chemical conditions as run 2.3 but AC is allowed to coagulate, which affects the floc density, mass transfer velocity, stickiness coefficient, and settling velocity.
5. Run 2.6: This run explores the possible impact of selective resuspension of AC relative to OC. The model started with the same physical and chemical settings as run 2.3 except the AC erosion rate was doubled. To simplify interpretation of the simulation results, AC did not flocculation in this run.
6. Run 2.7: Earlier we found the initial sediment HOC distribution played an important role in the fate of HOCs in the water column. In this run, the model uses the same physical setting as run 2.5, AC is allowed to coagulate, and sediment HOCs are evenly distributed between OC and AC

4.7.3 Results and Discussions

The temporally varying results for the AC, OC, TSS, and related HOC variables for all scenarios are shown in Figures 4.3 to 4.7. Similar to run 1.6, although AC only responded with 1% of TSS, the water column HOC decreased by 50% in run 2.2

compared to run 2.1 to emphasize the importance of AC on the fate of HOCs in the water column.

The settling velocity is an important factor determining the fate of AC and related HOC in the water column. Run 2.2 assumed that AC did not aggregate with flocs and that its shape did not follow principles of fractal geometry. Therefore, Stokes' law was used to calculate the AC settling velocity in run 2.2 instead of the fractal geometry adjusted settling equation used in run 2.3. Under the same resuspension flux, fractal settling was slower and the steady state AC concentration in the water column was 5.5 times higher in the run that assumed AC became a part of the floc (run 2.3) rather acting as an independent solid (run 2.2). Although the dissolved and organic carbon-associated HOC concentration did not significantly change between runs 2.2 and 2.3, the AC-related HOC concentration was 5.3 times higher when the AC particles followed fractal rather than Stokes settling.

After including the flocculation effect (run 2.5), the resuspended AC concentration and total water column HOC concentration were approximately 5 times higher compared to the run without flocculation (run 2.2). Further, runs 2.4 and 2.5 compared the effect of flocculation on OC and AC and their related HOC concentrations (Figures 4.11 and 4.12). The impact of AC as a part of the flocs was discussed in the previous paragraph. After including flocculation of AC, the suspended AC concentration decreased because of the formation of larger AC contained flocs (Figure 13). Thereafter, the total HOC (ng/L) in the water column decreased by 10%, especially AC-related HOC, in run 2.5 compared to run 2.4. However, the total water column HOC in run 2.5 was still 4.8 times higher than in run 2.2. When full flocculation is considered in a dynamic particle

environment that includes all three types of solids, sediment resuspension, settling, and kinetic-limited HOC partitioning, the steady state total suspension solid concentration decreased by 34%, the total PCB-52 concentration in the water column increased nearly four-fold, and the water column HOC residence time increased by 37%.

The importance of suspended AC on the water column HOC concentration was demonstrated again in run 2.6 (Figure 14). With the slow mass transfer velocity from AC, the sediment AC resuspension rate was doubled, and yet the dissolved HOC concentration only increased by less than 1% after a 53 hour simulation, even though the total water column HOC and AC-related HOC concentrations were almost double those used in run 2.4. Therefore, during a strong resuspension event (such as run 2.6), the dissolved HOC concentrations remained as low as those prior to the resuspension event due to the very strong sorptive capacity of the activated carbon.

In stage one of the model developments discussed earlier in this chapter, we observed that the interaction time is a very important issue under conditions when AC has insufficient time to adsorb or compete for HOCs with other organic particles. When flocculation is added, the interaction time also demonstrated a strong influence on the fate of water column HOC (Figure 4.15). Without sufficient interaction time, the total HOC concentration in run 2.7 increased nearly eight times compared to run 2.2.

4.7 *Summary*

AC has a stronger sorption capacity than flocs of the same size and with the same organic carbon content. However, AC has a slower sorption rate (smaller mass transfer velocity) due to its relatively lower porosity and higher K_{AC} value. Over long time scales, the majority of HOCs will be sequestered in the activated carbon. However, in the

absence of flocculation in a dynamic erosional and depositional environment, the initial HOC distribution and the relatively slower sorption rate for AC are the dominant factors controlling the distribution of HOCs in the water column. The interaction time becomes a very important issue under conditions when AC does not have sufficient time to adsorb or compete with other organic particles. The above analysis considers the scenario where the amended activated carbon does not interact with the natural sediments or other aquatic particles to form flocs. We found that the lower porosity and higher settling velocities of AC relative to natural OC result in lower steady-state AC concentrations in the water column, and longer times to reaching sorptive equilibrium.

With the enrichment of AC, the total water column HOC concentration significantly decreases. Further, when AC aggregates with sediment particles, AC is going to have the same physical properties as particles of similar floc size, which results in a slower settling velocity compared to single AC particles, and will coherently have a longer residence time and higher AC related HOC concentration in the water column. When activated carbon is added to contaminated sediments, the total PCB concentration in the water column decreases by 90% (123.4 to 11.4 ng/L). If the activated carbon coagulates with the resuspended sediment, this decrease is partially offset by some activated carbon being entrained in slowly-settling flocs, and the steady-state PCB concentration is 61 ng/L.

4.8 *Figure Captions*

Figure 4.1 and 4.2: Flow diagrams for refining the PCB fate and the flocs transport model to include (1) flocculation kinetics, (2) PCB partitioning kinetics, and (3) activated carbon as a state variable.

Figure 4.3: Comparison of the predicted temporally varying resuspended organic carbon, activated carbon, and inorganic solids among different scenarios at stage two

Figure 4.4: Comparison of the predicted temporally varying water column PCB 52 (ng/L) in the organic carbon, activated carbon, and dissolved water among different scenarios at stage two

Figure 4.5: Comparison of the predicted temporally varying total water column PCB 52 (ng/L) among different scenarios at stage two

Figure 4.6: Comparison of the predicted steady state resuspended organic carbon, activated carbon, and inorganic solids among different scenarios at stage two

Figure 4.7: Comparison of the predicted steady state water column PCB 52 (ng/L) in the organic carbon, activated carbon, and dissolved phases among different scenarios at stage two

Figure 4.8: Predicted behavior of PCB 52 and solids in carbon-amended sediments. This is a reference run without activated carbon (Run 2.1). The model starts with the equilibrium sediment PCB 52 between organic carbon and porewater and includes flocculation, resuspended, and deposition processes.

Figure 4.9: Predicted behavior of PCB 52 and solids in activated carbon-amended sediments. This is a run without flocculation that includes activated carbon (Run 2.2). The model starts with the equilibrium sediment PCB 52 among activated carbon, organic carbon and porewater and includes resuspension and deposition processes. The activated carbon settling is calculated using the Stokes' law settling velocity equation.

Figure 4.10: Predicted behavior of PCB 52 and solids in activated carbon-amended sediments. This is a run without flocculation that includes activated carbon (Run 2.3). The model starts with the equilibrium sediment PCB 52 among activated carbon, organic carbon and porewater and includes resuspension and deposition. The activated carbon settling is calculated using the fractal geometry adjusted settling velocity equation.

Figure 4.11: Predicted behavior of PCB 52 and solids in activated carbon-amended sediments. This is a run with flocculation of organic carbon but not activated carbon (Run 2.4). The model starts with the equilibrium sediment PCB 52 among activated carbon, organic carbon and porewater and with resuspended, and deposition

processes. The activated carbon is adapted fractal geometry adjusted settling velocity equation.

Figure 4.12: Predict behavior of PCB 52 and solids in activated carbon-amended sediments: this is a run with flocculation on both carbon solids and activated carbon is involved (Run 2.5). The model starts with the equilibrium sediment PCB 52 among activated carbon, organic carbon and porewater and with resuspended, and deposition processes. The activated carbon is adapted fractal geometry adjusted settling velocity equation.

Figure 4.13: Predict steady state fraction of activated carbon size distribution at run 2.5

Figure 4.14: Predict behavior of PCB 52 and solids in activated carbon-amended sediments: this is a run with double activated carbon erosion flux, flocculation on organic carbon and activated carbon is involved (Run 2.6). The model starts with the equilibrium sediment PCB 52 among activated carbon, organic carbon and porewater and with resuspended, and deposition processes. The activated carbon is adapted fractal geometry adjusted settling velocity equation.

Figure 4.15: Predict behavior of PCB 52 and solids in activated carbon-amended sediments: this is a run with flocculation on both carbon solids and activated carbon is involved (Run 2.7). The model was initialized with 50 % of sediment PCB 52 in activated carbon and organic carbon respectively. The activated carbon is adapted fractal geometry adjusted settling velocity equation

Figure 4.1:

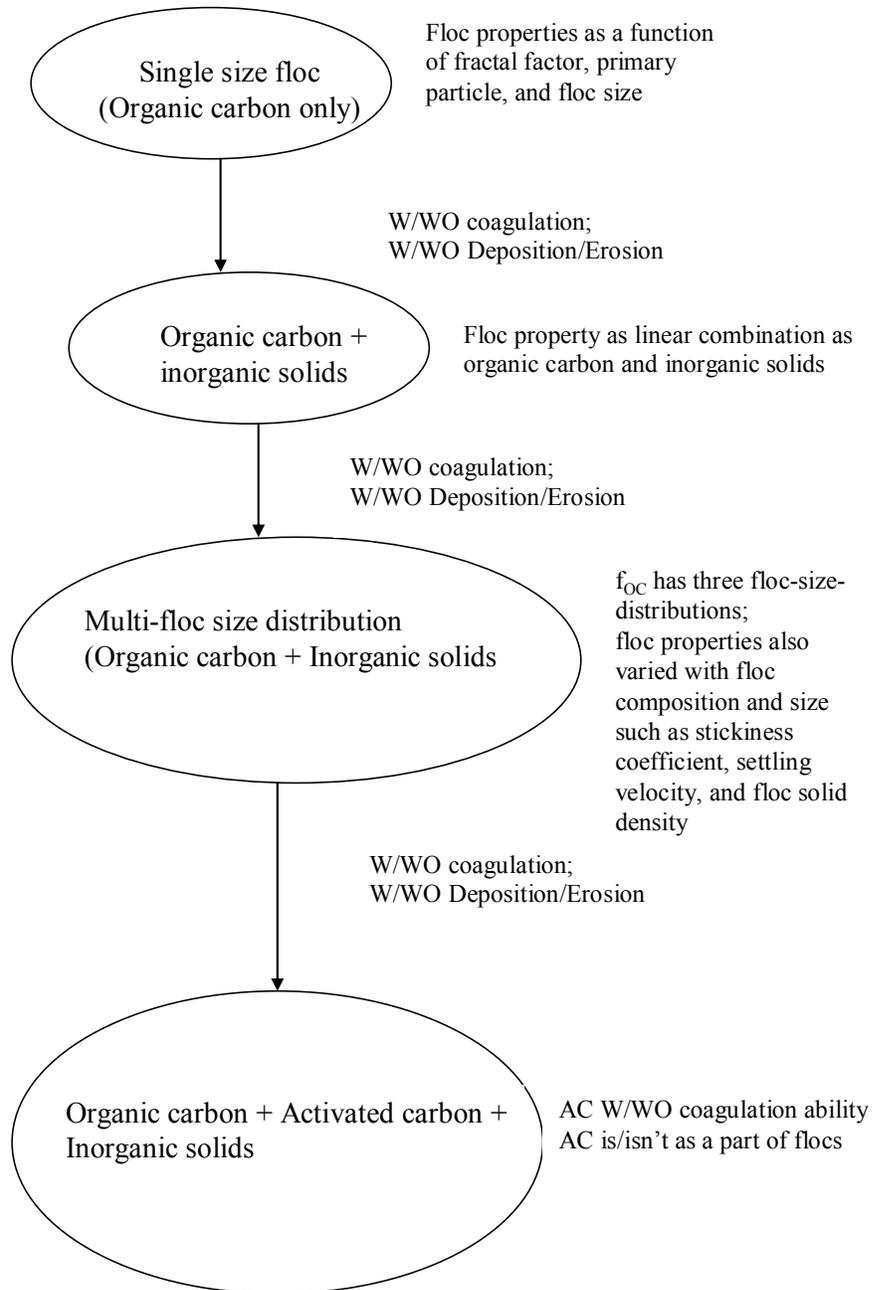


Figure 4.2:

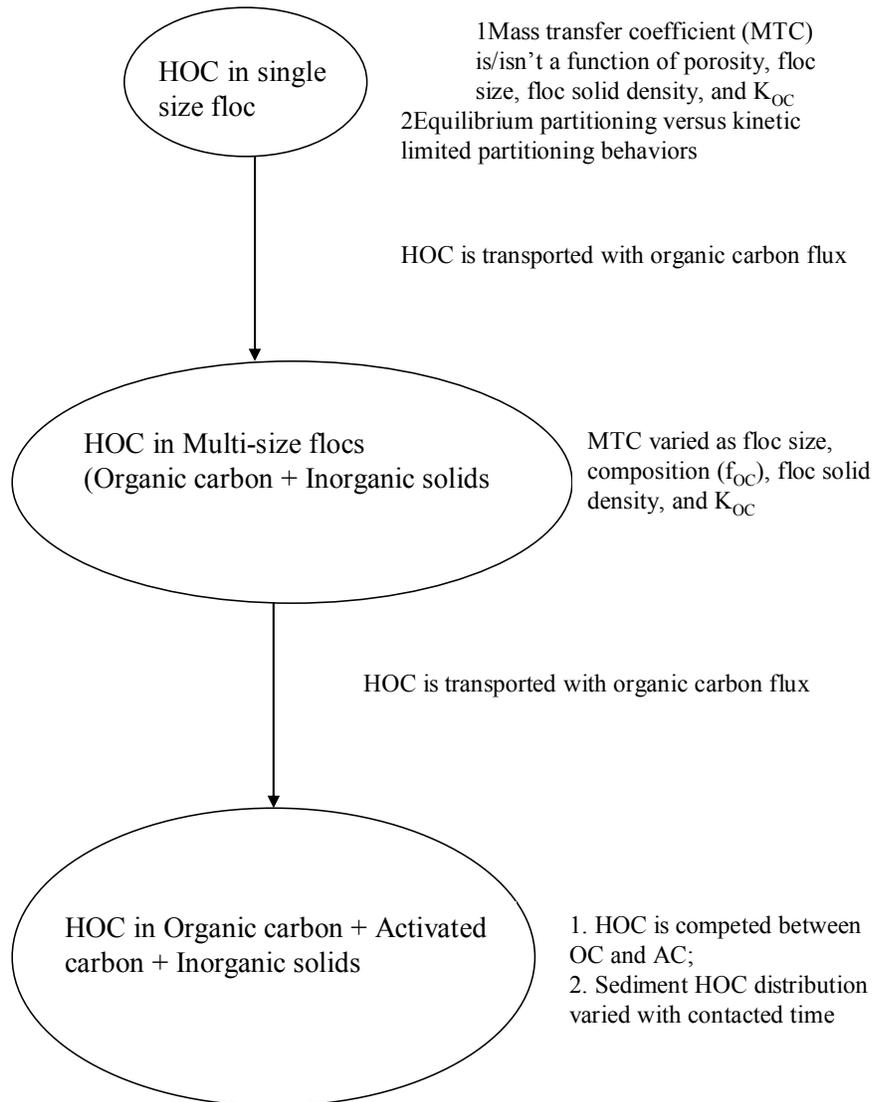


Figure 4.3:

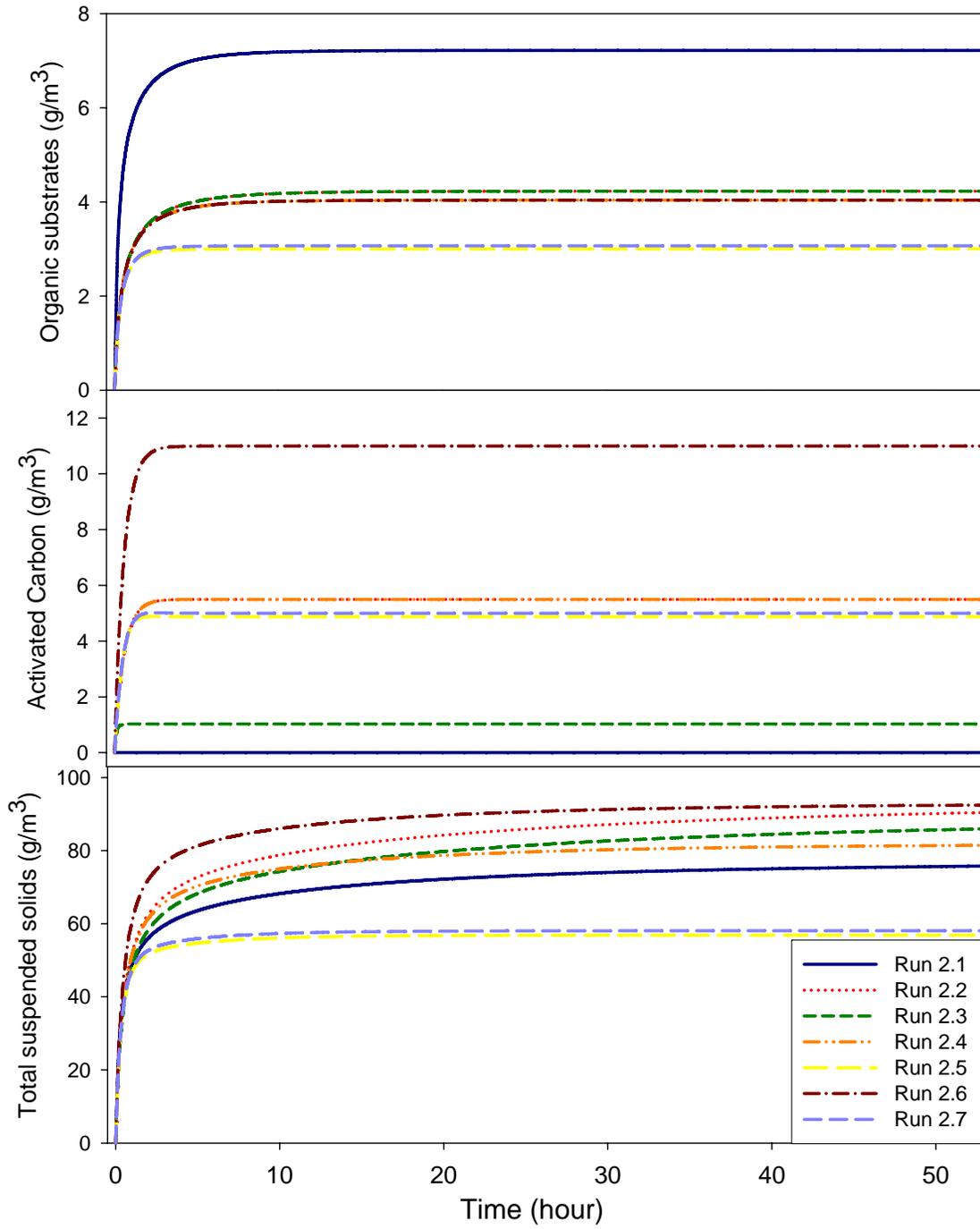


Figure 4.4:

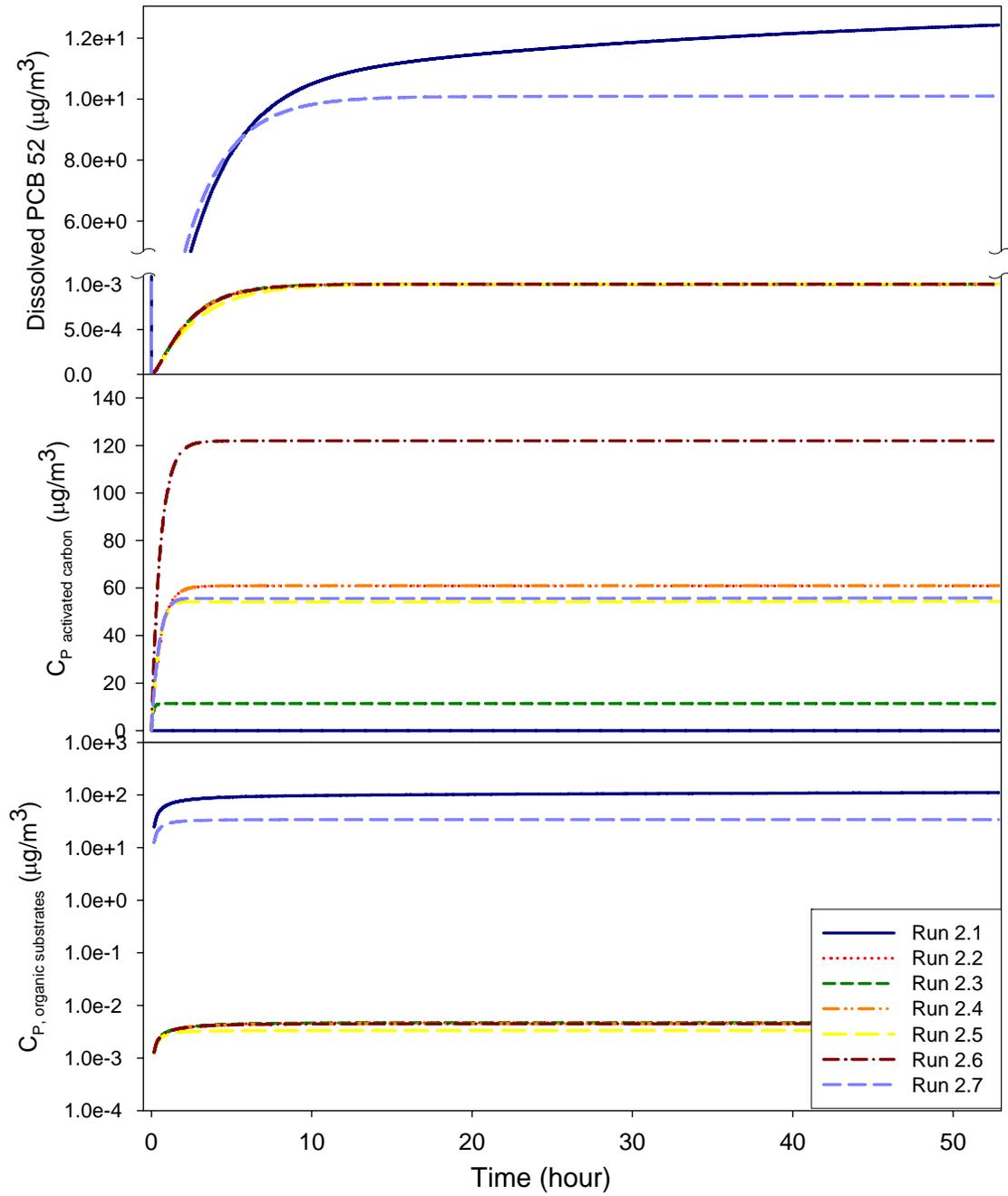


Figure 4.5:

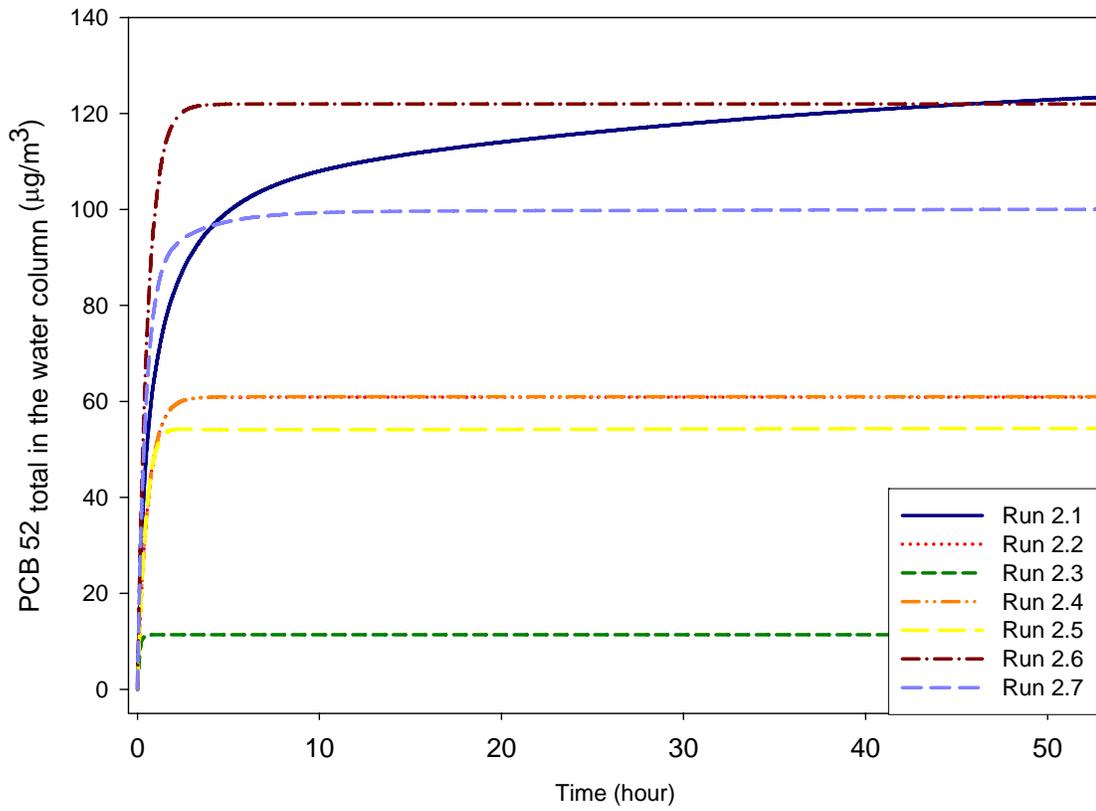


Figure 4.6

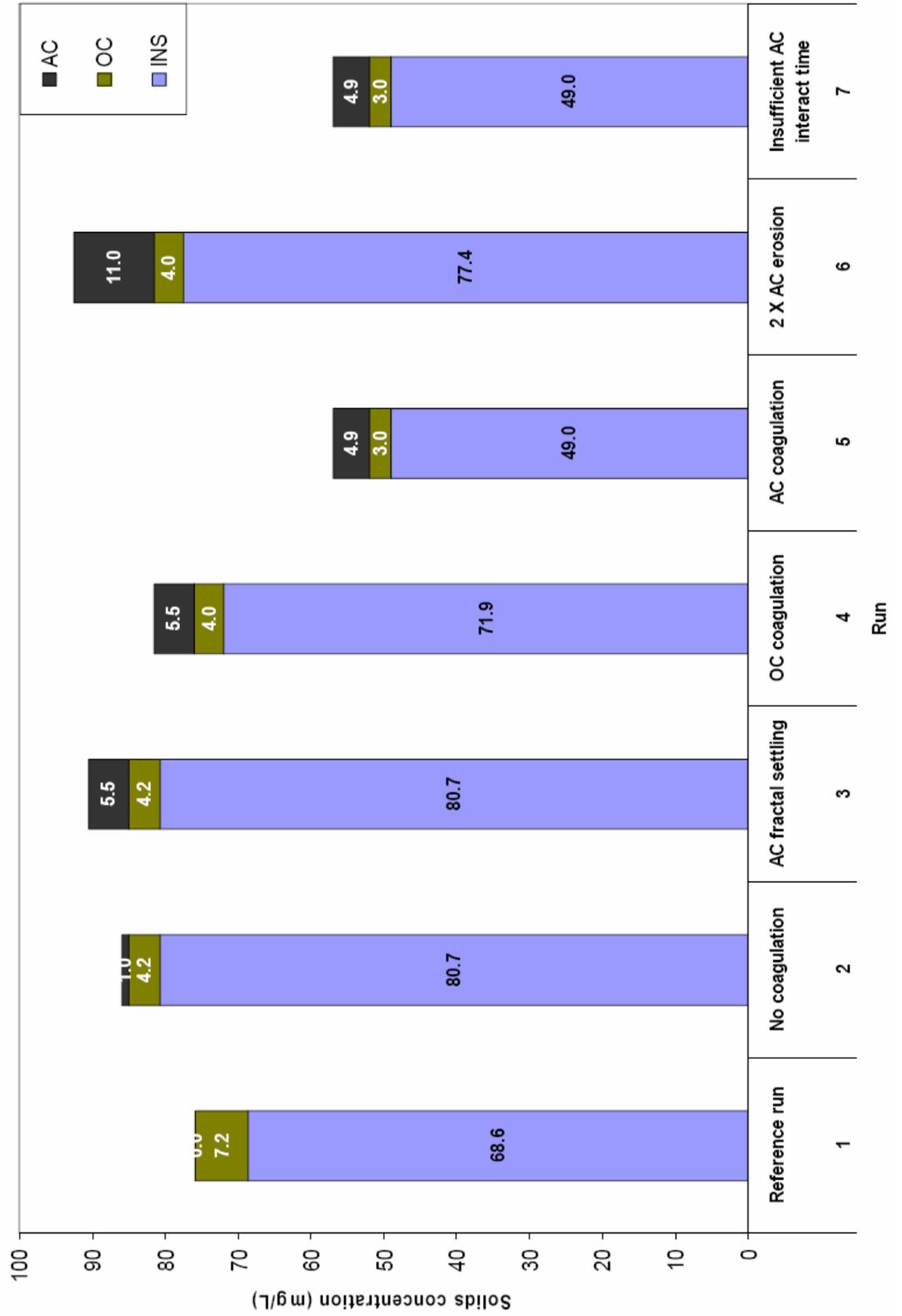


Figure 4.7



Figure 4.8:

Run 2.1

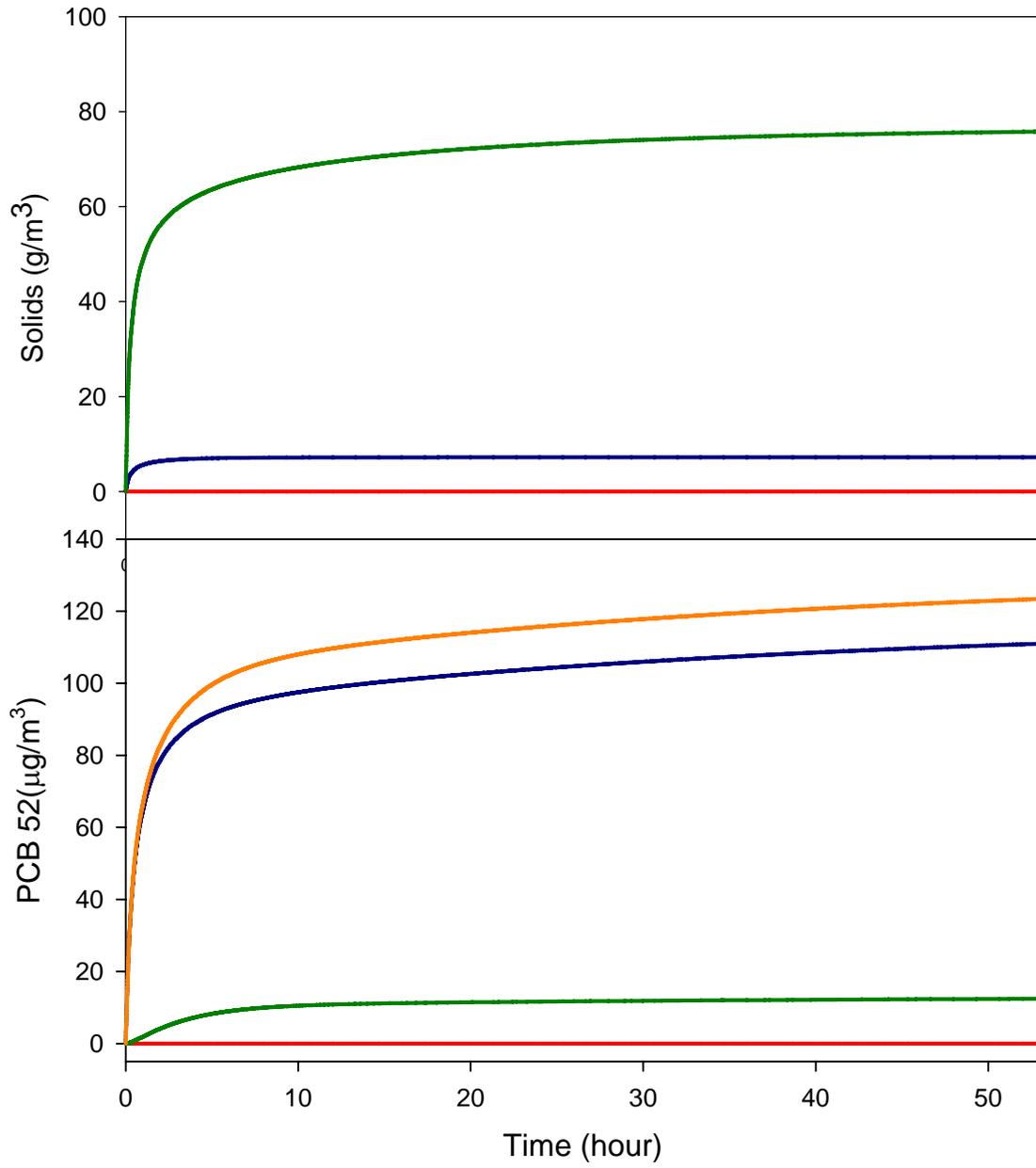


Figure 4.9:

Run 2.2

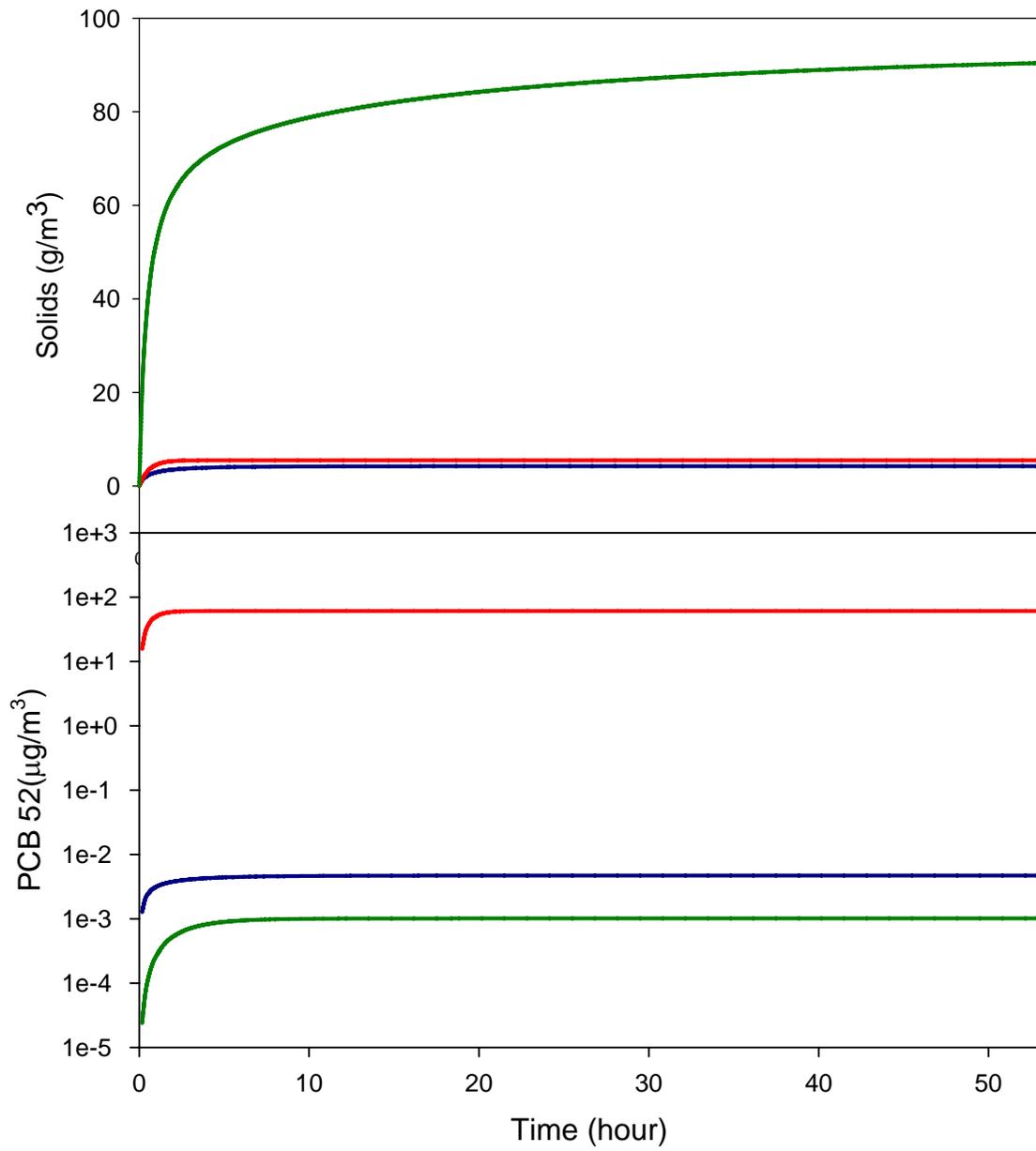


Figure 4.10:

Run 2.3

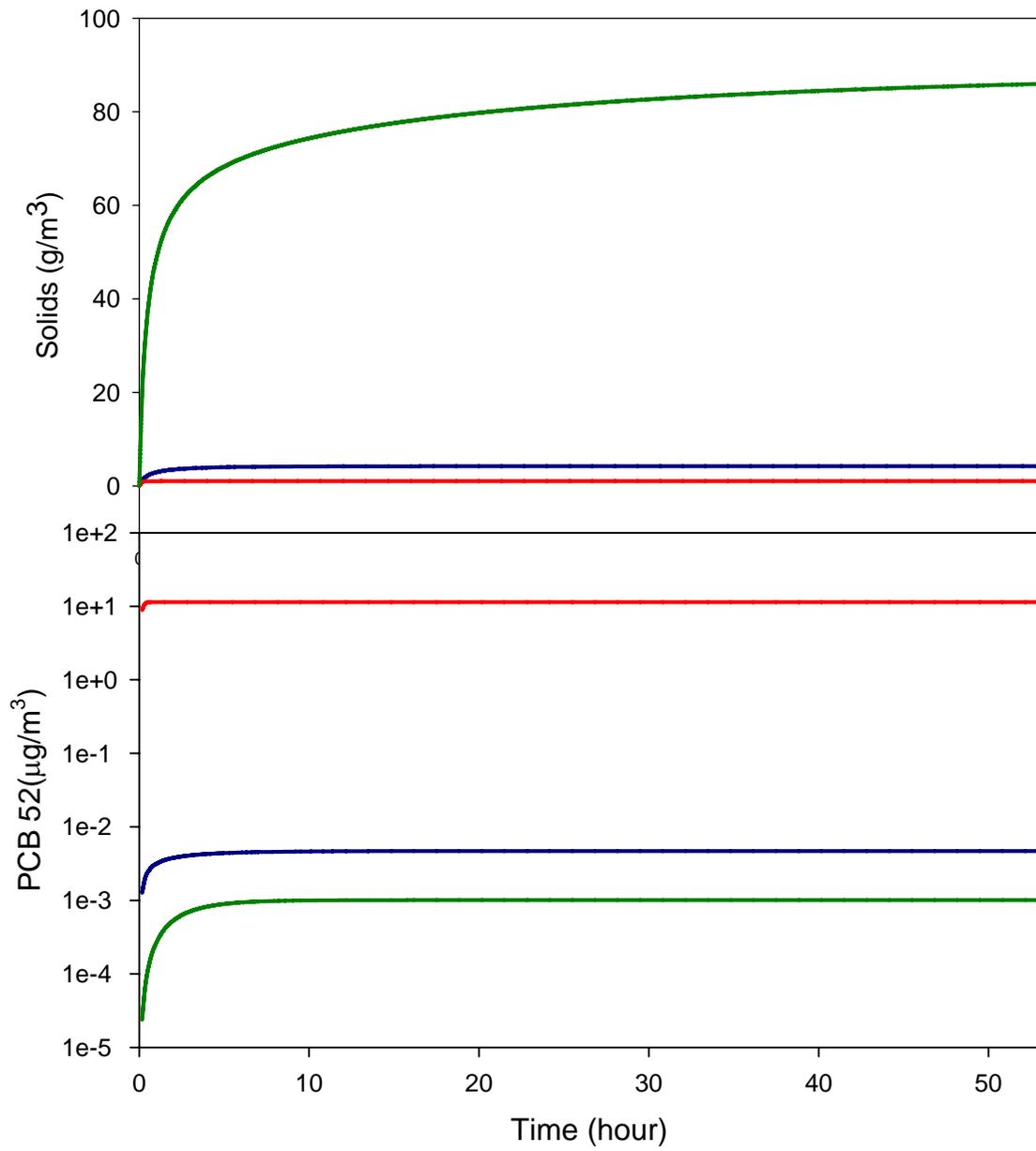


Figure 4.11:

Run 2.4

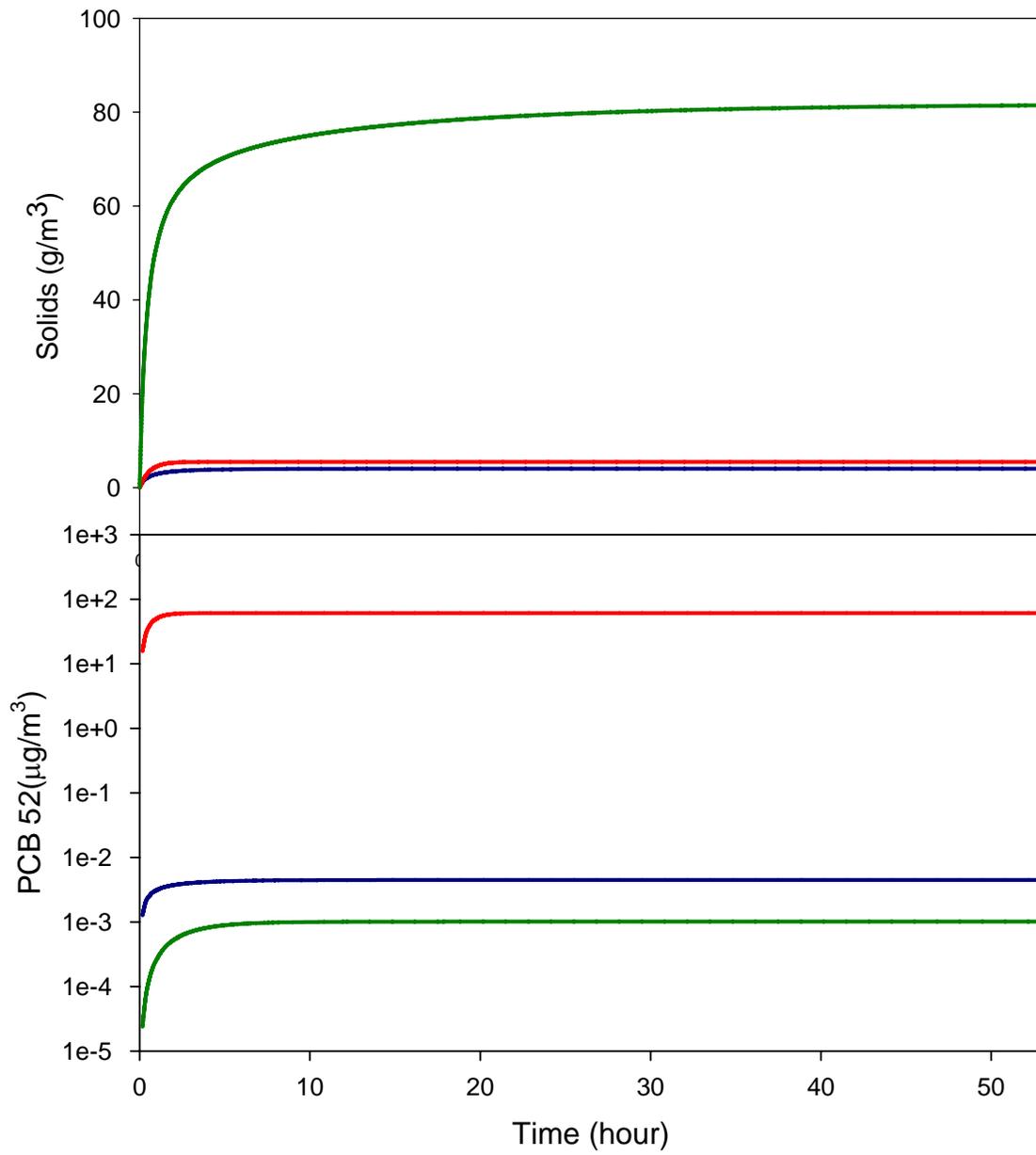


Figure 4.12:

Run 2.5

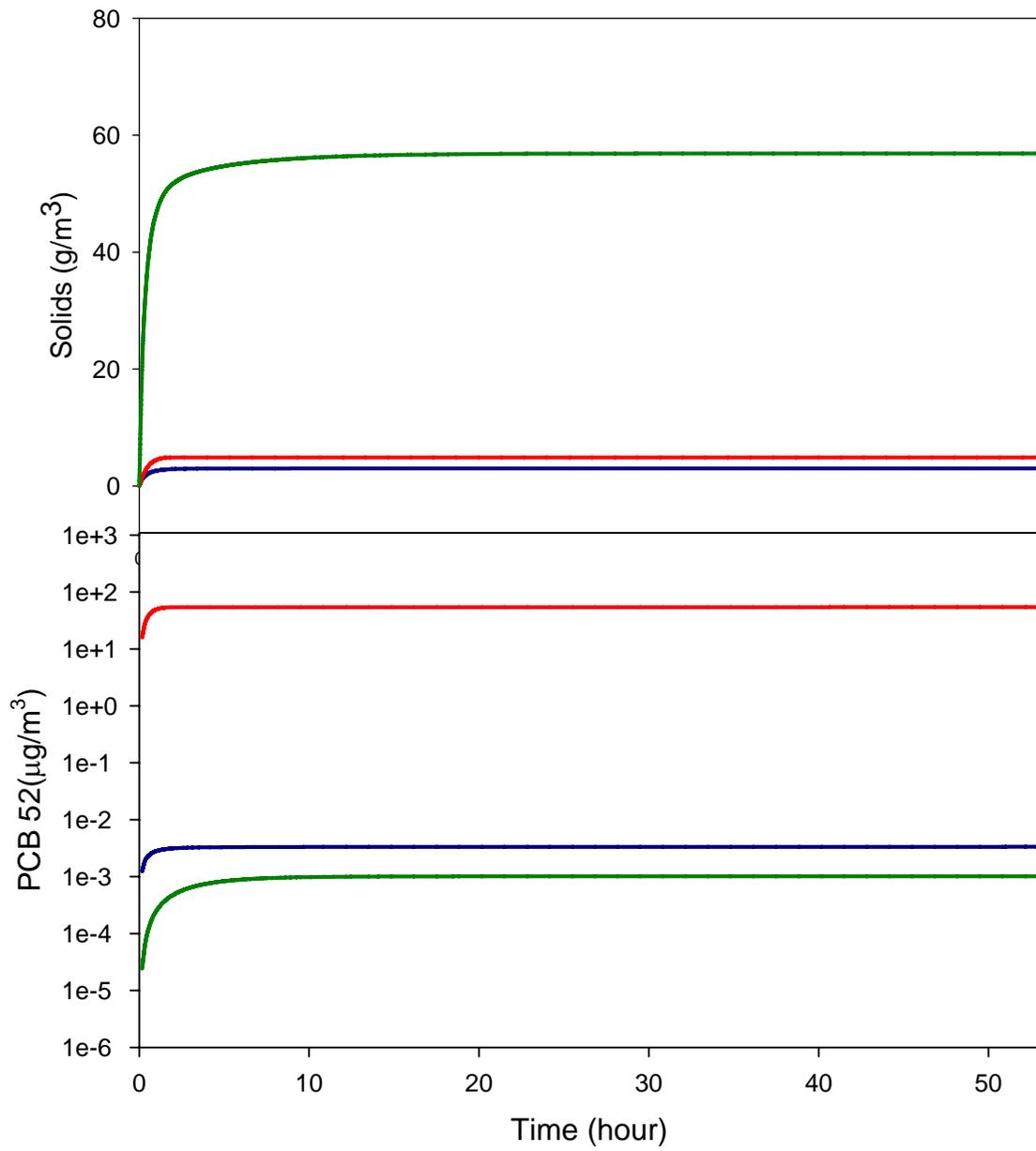


Figure 4.13:

f_{AC} size distribution for run 2.5

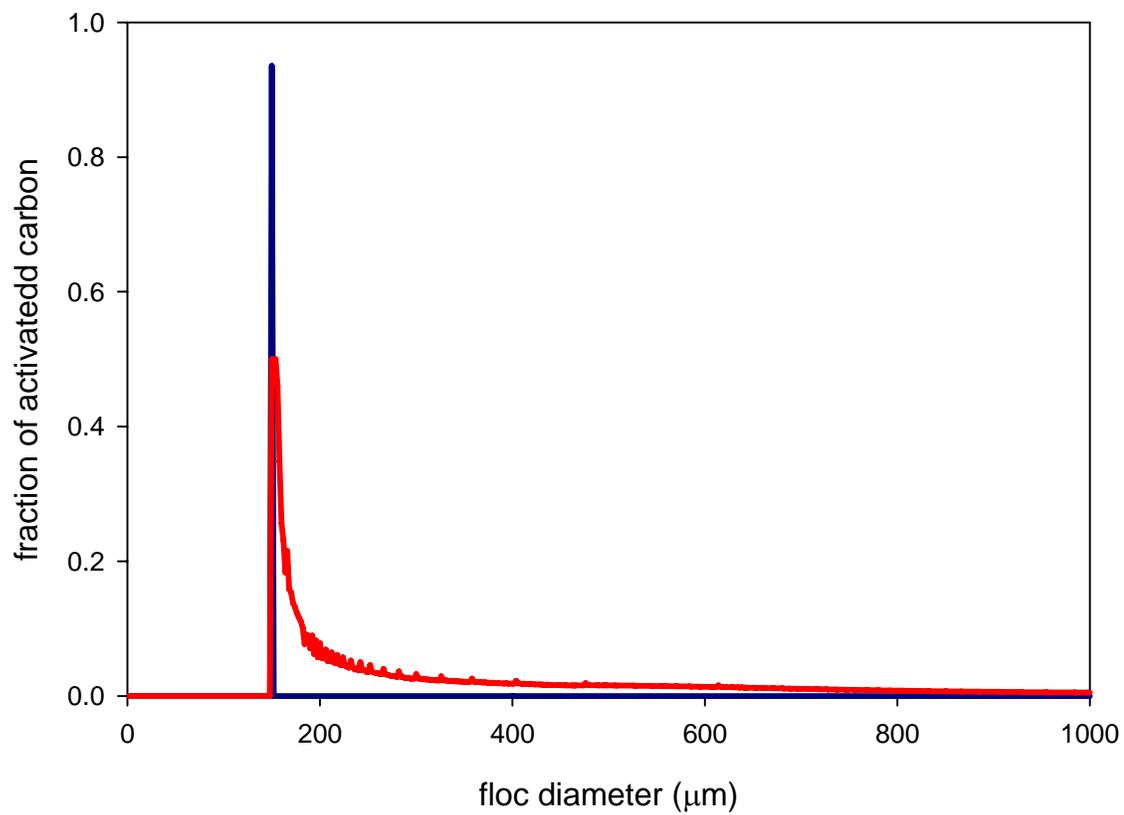


Figure 4.14:

Run 2.6

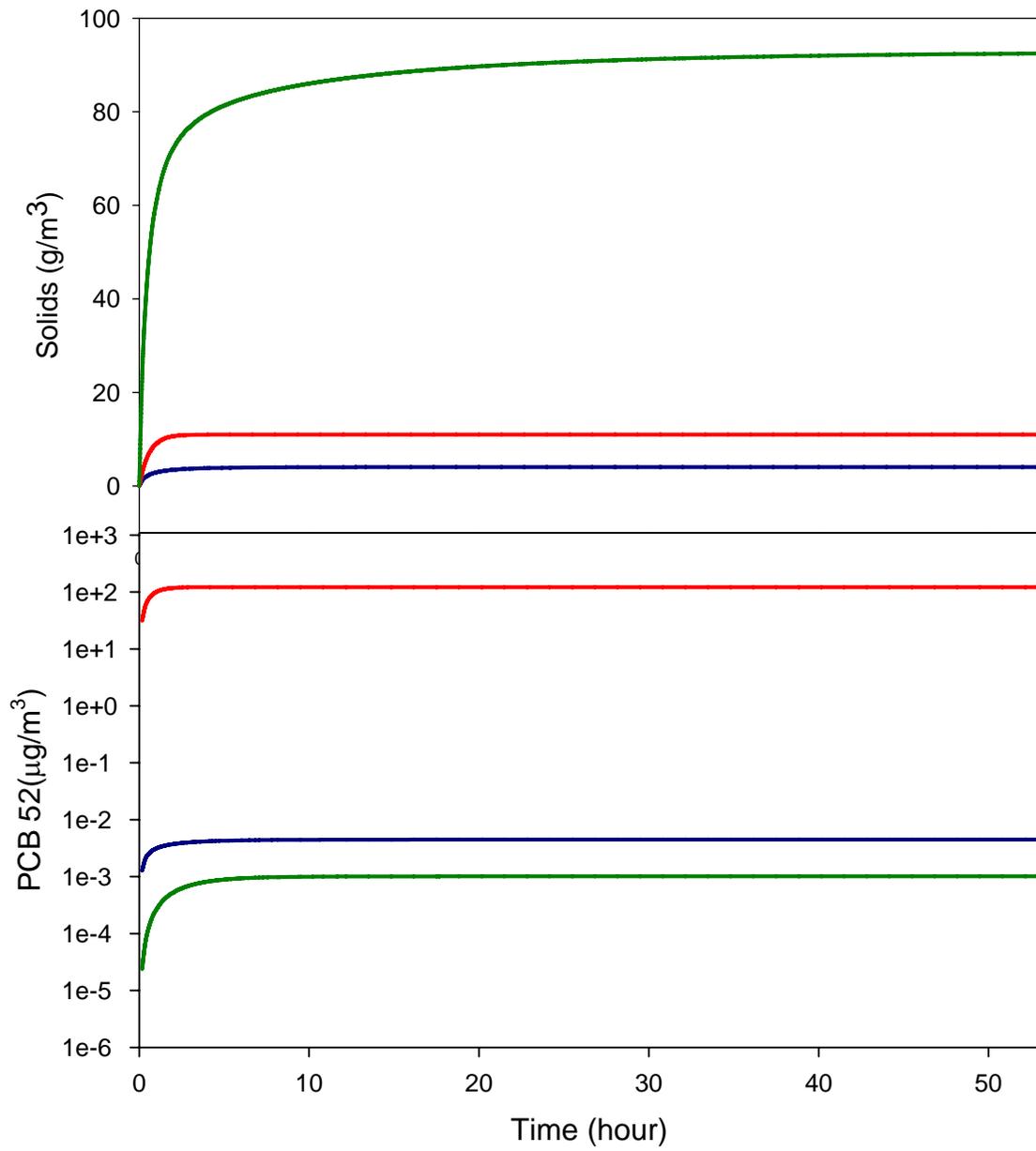


Figure 4.15:

Run 2.7

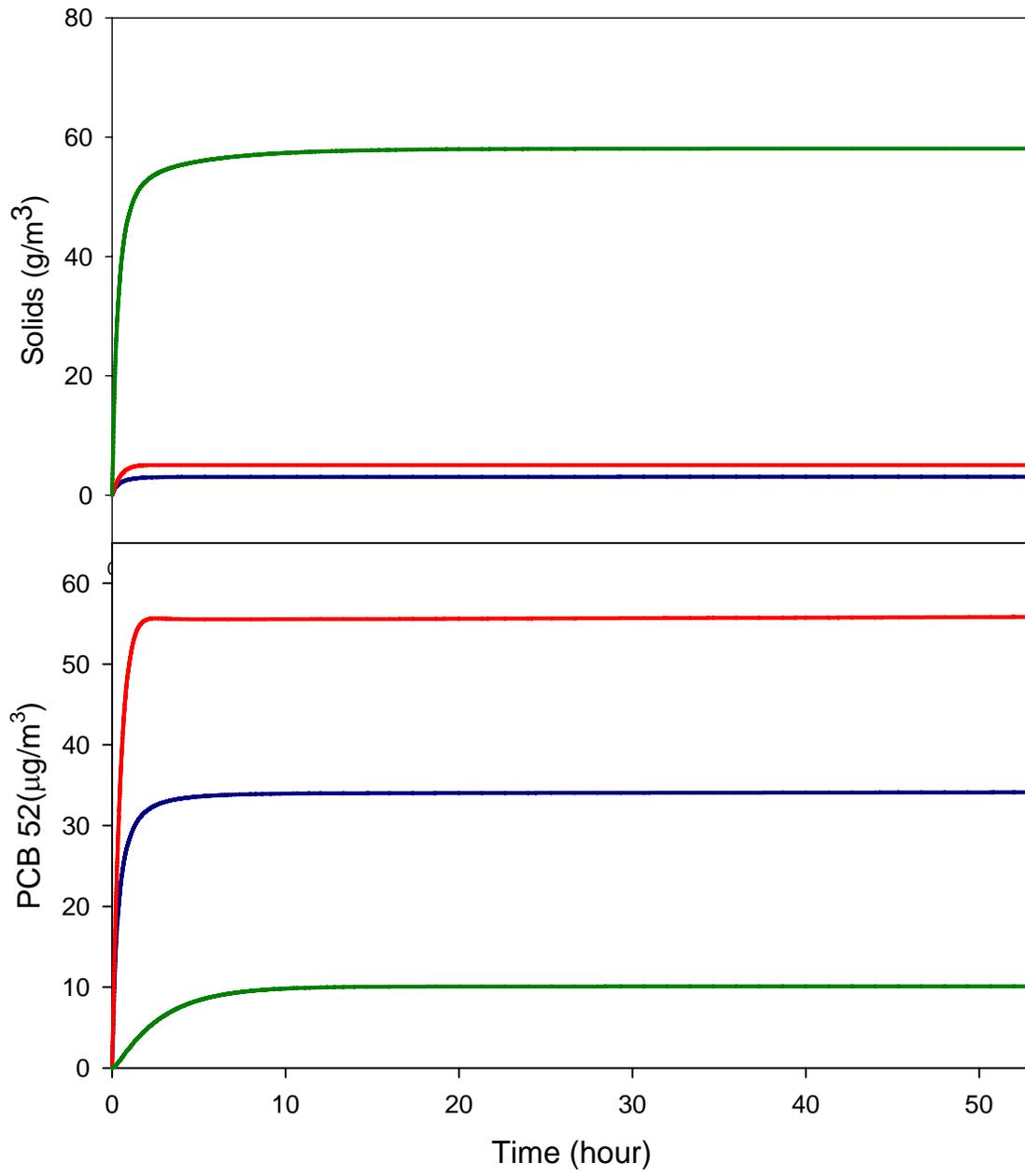


Table 4.2: Seven model scenarios were created to successively explore the interaction of AC and OC and its impact of HOC partitioning at stage two

run	K _{oc} , K _{ac}	coagulation effect: AC	coagulation effect: OC	e-d AC flux	e-d OC flux	initial condition and special settings
2.1	1	na	on	na	on	same as chapter 3
2.2	1, 1000	off	off	on	on	Ws=Stokes law equation for AC; sediment HOC reaching equilibrium
2.3	1, 1000	off	off	on	on	Ws=Winterwerp equation for AC and OC sediment HOC reaching equilibrium
2.4	1, 1000	off	on	on	on	sediment HOC reaching equilibrium
2.5	1, 1000	on	on	on	on	sediment HOC reaching equilibrium
2.6	1, 1000	off	on	X2	on	sediment HOC reaching equilibrium
2.7	1, 1000	on	on	on	on	AC and OC took 50% of sediment HOC OC+AC=OC _{run1} ; and sum(OC)=AC _{D150um}

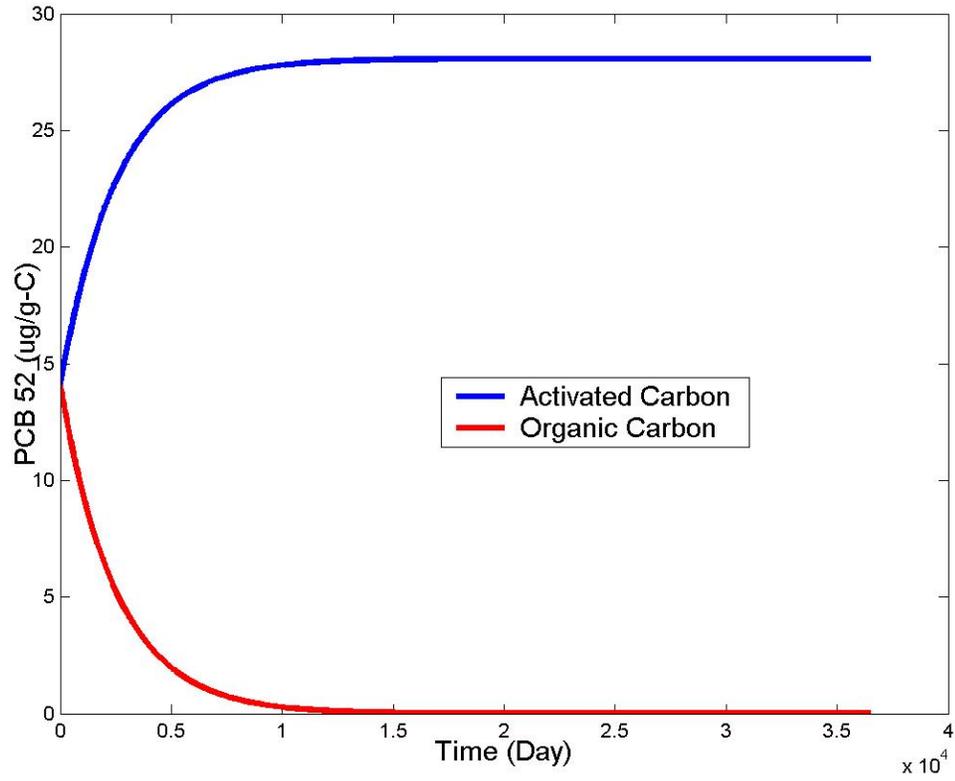
Appendix 1 Steady State Sediment PCB 52 Distribution

In chapter four, two important implications are made from the stage one runs. First, the initial sediment PCB distribution and the relatively slower sorption rate for activated carbon are the dominant factors controlling the distribution of PCBs in the water column in a dynamic erosional and depositional environment. Second, the interaction time becomes a very important issue when the activated carbon does not have sufficient time to adsorb or compete with other organic particles.

In stage two, the STORM data is adapted by the model scenarios. However, STORM only measured the steady state sediment PCB 52 equaling $14.04 \mu\text{g/g-OC}$ in the absence of activated carbon. Therefore, the model needs to make certain assumptions regarding the initial sediment PCB distribution between organic carbon and activated carbon in order to include the activated carbon in stage two. First, the initial gross sediment f_{OC} has the same values in all the runs. Further, the initial sediment organic carbon is divided into two groups, namely organic carbon and activated carbon. Each group shares a 50% gross organic carbon concentration. Second, to minimize the impact of initial sediment PCBs on the model predictions, the model starts with an equilibrium status for PCBs among organic carbon, activated carbon, and porewater in the sediment. To have a better assumption of the initial sediment PCB distribution, a 100-year simulation was conducted. At time zero, both activated carbon and organic carbon had the same PCB 52 concentration of $14.04 \mu\text{g/gC}_{(\text{AC}+\text{OC})}$. The model disconnected all physical and chemical transport mechanisms except for diffusion exchange among pore water, organic carbon, and activated carbon in the top sediment layer. At a steady state,

99.99 % of PCB 52 was absorbed by activated carbon (28.0764 $\mu\text{g/g-AC}$), and only 0.01% was absorbed by organic carbon (0.0034 $\mu\text{g/g-OC}$) (Figure A1).

Figure A-1: The model predicted temporal varied sediment PCB 52 concentrations between activated carbon and organic carbon during a 100 year simulation



Appendix 2 Model Equation and Parameter Definition

$$\frac{\partial N_i}{\partial t} = F_i + \left(\frac{W_{S,i}}{x} \times N_i \right) + E_{b,i} \quad \text{Equation 2.1}$$

$$F_k = \frac{1}{2} \sum_{i+j=k} \alpha_{i,j} \beta_{i,j} N_i N_j - N_k \sum_{i=1}^{\infty} \alpha_{i,k} \beta_{i,k} N_i \quad \text{Equation 2.2}$$

$$\beta_{i,j} = \beta_{shear_i,j} + \beta_{settling_i,j} + \beta_{Brownian_i,j} \quad \text{Equation 2.3}$$

$$\beta_{recrilinea\ r_shear} = \frac{G}{6} (D_i + D_j)^3 \quad \text{Equation 2.4}$$

$$\beta_{recrilinea\ r_settling} = \frac{\pi}{4} (D_i + D_j)^2 |W_{S,i} - W_{S,j}| \quad \text{Equation 2.5}$$

$$\beta_{Brownian} = \pi (Diff_{floc,i} + Diff_{floc,j}) (D_i + D_j) \quad \text{Equation 2.6}$$

$$\beta_{curvilinea\ r} = \left(\frac{(D_i - 2D_j)^2}{(D_i + D_j)^2} \right) \times \beta_{recrilinea\ r} \quad \text{Equation 2.7}$$

$$E_{b,i} = \sigma \left(\frac{\tau_b}{\tau_c} - 1 \right) \quad \text{Equation 2.8}$$

$$W_s = \frac{a}{18b} \frac{(\rho_{floc,dry} - \rho_w)}{\mu} D_p^{3-nf} \frac{D_f^{nf-1}}{1 + 0.15 Re^{0.687}} \frac{1}{1 + 2.5\phi} \quad \text{Equation 2.9}$$

$$V_{dry,i} = N_i \times V_{spherical_equivant_cluster,i} \quad \text{Equation 2.10}$$

$$C_{dry,i} = V_{dry,i} \times \rho_{dry,i} \times (1 - \phi_i) \quad \text{Equation 2.11}$$

$$\phi_f = \left[1 - \left(\frac{D_f}{D_p} \right)^{nf} \left(\frac{D_p}{D_f} \right)^3 \right] \quad \text{Equation 2.12}$$

$$\rho_{dry,i} = 2 \times f_{OC,i} \times \rho_{biotic} + (1 - 2 \times f_{OC,i}) \times \rho_{clay} \quad \text{Equation 2.13}$$

$$\rho_{bulk,i} = \rho_{dry,i} \times (1 - \phi_i) + \rho_{water} \times \phi_i \quad \text{Equation 2.14}$$

$$\alpha_i = 2 \times f_{OC,i} \times \alpha_{biotic} + (1 - 2 \times f_{OC,i}) \times \alpha_{clay} \quad \text{Equation 2.15}$$

$$Residence_Time = \frac{\sum [C_{dry,i} \times \left(\frac{V_{tank}}{Area \times W_s} \right)]}{\sum C_{dry,i}} \quad \text{Equation 2.16}$$

$$\frac{\partial TSS}{\partial t} = -\frac{w_{S,TSS}}{x} TSS \quad \text{Equation 2.17}$$

$$\frac{\partial N_{i,surface}}{\partial t} = -\frac{m_{1,2}}{(sz_1 \times sz_2)} (N_{i,surface} - N_{i,deeper}) - E_{b,i} \quad \text{Equation 3.2}$$

$$\frac{\partial N_{i,deeper}}{\partial t} = \frac{m_{1,2}}{(sz_1 \times sz_2)} (N_{i,surface} - N_{i,deeper}) \quad \text{Equation 3.3}$$

$$\frac{\partial HOC_i}{\partial t} = HOC_flux_{F_N, flocc} + J_{i,j} \quad \text{Equation 3.4}$$

$$J_i = k \times A \times N_i \times \left(HOC_{dissolved} - \frac{HOC_i}{OC_i \times K_{OC}} \right) \quad \text{Equation 3.5}$$

$$k = \frac{D^*}{z} = \frac{\phi_f f D m}{z \times (K_{OC} f_{OC} \times (1 - \phi_f) \times \rho_{f,solid} + \phi_f)} \quad \text{Equation 3.6}$$

$$(R - r) = \frac{1}{4/3 \pi R^3} \int_0^R 4\pi r^2 (R - r) dr = \frac{\frac{1}{3} \pi R^4}{\frac{4}{3} \pi R^3} = \frac{1}{4} R = z \quad \text{Equation 3.7}$$

$$\text{Deviation from equilibrium} = (C_{p,t} - C_{p,t}^*) / C_{p,t}^* \quad \text{Equation 3.8}$$

$$C_{p,t}^* = K_{OW} C_{d,t} F \quad \text{Equation 3.9}$$

$$\rho_{dry,i} = 2 \times f_{OC,i} \times \rho_{biotic} + (1 - 2 \times f_{OC,i} - f_{AC,i}) \times \rho_{clay} + f_{AC,i} \times \rho_{AC} \quad \text{Equation 4.1}$$

$$\rho_{bulk,i} = \rho_{dry,i} \times (1 - \phi_i) + \rho_{water} \times \phi_i \quad \text{Equation 4.2}$$

$$\alpha_i = 2 \times f_{OC,i} \times \alpha_{biotic} + (1 - 2 \times f_{OC,i} - f_{AC,i}) \times \alpha_{clay} + f_{AC,i} \times \alpha_{AC} \quad \text{Equation 4.3}$$

N_i	the floc number concentration ($\# \text{ m}^{-3}$)
$W_{S,i}$	the floc settling velocity (m sec^{-1})
$W_{S \text{ TSS}}$	the gross settling velocity for TSS (m sec^{-1})
F_i	the flocculation effect on particle number balance ($\# \text{ m}^{-3} \text{ sec}^{-1}$)
x	the water column depth (m)
$E_{b,i}$	the pure sediment resuspension flux ($\# \text{ m}^{-3} \text{ sec}^{-1}$)
α	the stickiness coefficient (dimensionless)
α_{biotic}	the organic carbon stickiness coefficient (dimensionless)
α_{clay}	the inorganic carbon stickiness coefficient (dimensionless)
α_{AC}	the activated carbon stickiness coefficient (dimensionless)
$\beta_{i,j}$	the cohesive collision frequency ($\text{m}^3 \text{ sec}^{-1}$)
β_{shear}	the collision frequency from shear ($\text{m}^3 \text{ sec}^{-1}$)
β_{settling}	the collision frequency from differential settling ($\text{m}^3 \text{ sec}^{-1}$)
β_{Brownian}	the collision frequencies from Brownian motion ($\text{m}^3 \text{ sec}^{-1}$)
G	the shear stress gradient (sec^{-1})
Diff	the ideal floc diffusion coefficient ($\text{m}^2 \text{ sec}^{-1}$)
σ	a calibration coefficient (dimensionless)
τ_b	bottom shear stress (dyne cm^{-2})
τ_c	critical shear stress (dyne cm^{-2})
a and b	shape parameters (dimensionless)
nf	the fractal dimension (dimensionless)
Re	the Reynolds number (dimensionless)
μ	the dynamic viscosity of fluid ($\text{g m}^{-1} \text{ sec}^{-1}$)
$D_{i,j}$	the floc diameter (m)
D_p	the diameter of primary particle (m)
ϕ	the floc porosity ($\text{m}^3 \text{ m}^{-3}$)
V_{dry}	the solid volume concentration ($\text{m}^3 \text{ m}^{-3}$)
$C_{\text{dry},i}$	the solid floc mass concentration (g m^{-3})
ρ_{floc}	floc solid density (g m^{-3})
ρ_w	water density (g m^{-3})
$\rho_{\text{dry},i}$	the dry floc density (g m^{-3})
$\rho_{\text{bulk},i}$	the bulk floc density (g m^{-3})
ρ_{biotic}	the organic carbon dry density (g m^{-3})
ρ_{clay}	the inorganic carbon dry density (g m^{-3})
ρ_{AC}	the activated carbon dry density (g m^{-3})
V_{tank}	the experimental tank volume (m^3)
Area	the sediment surface area (m^2)
SZ_1	the top sediment thickness (m)
SZ_2	the deeper sediment thickness (m)
HOC_i	the HOC concentration in the floc cluster i (ng L^{-1})
$\text{HOC_Flux}_{F_n, \text{floc}}$	the HOC flux due to particle transport by flocculation, resuspension, settling, or bioturbation ($\text{ng L}^{-1} \text{ sec}^{-1}$)
$J_{i,j}$	the HOC being transported between particles variable i and surrounding water variable j ($\text{ng L}^{-1} \text{ sec}^{-1}$)
k	the mass transfer velocity (m/sec),
A	the specific interfacial exchange area (m^2/floc),

OC_i	the organic carbon state variable concentration (g-carbon m^{-3})
$HOC_{dissolved}$	the dissolved HOC concentration (ng-HOC m^{-3})
K_{OC}	the organic carbon normalized partition coefficient (m^3 g-carbon $^{-1}$)
z	the volume-weighted effective interior diffusion distance (m)
f	the tortuosity factor (dimensionless)
D^*	the adjusted diffusion coefficient (m^2 sec $^{-1}$)
D_m	the chemical molecular diffusion coefficient in the water (m^2 sec $^{-1}$)
r	the distance from the center of the sphere is defined as r (m)
R	the radius of the sphere (m)
K_{OW}	the octanol-water partition coefficient (m^3 g $^{-1}$)
f_{AC}	the fraction of activated carbon (g-AC g-TSS $^{-1}$)
f_{OC}	the fraction of organic carbon for floc at size i (g-OC g-TSS $^{-1}$)

Appendix 3 Model Program Codes

```
C      Last change:  AC      8 May 2008      6:58 pm
!AC simulation stage 2 design === full run with flocculation and with
50% of OC eroded flux and no coagulation
!Sediment HOC assumed reaching equilibrium among porewater, Oc, and AC

!check shear stress, MTC, sz, and fractal factor value later it's not
the same as original design

!      Last change:  AC      20 Apr 2008      10:32 pm
!correct the diffusion flux bug between HPOC and POC
!add and expend area to At, Ain, and Aext
!using new MTC equation: MTCtotal=Ain/At*MTCin+Aext/At*MTCext
!adjust diffusion equation from A(i)*kk... At(i)*MTCt(i)*.....
!adjusted L1-L2 mixing equation to avoid numerical and mass balance bug

!Added prewater HOC and sediment HOC in L1 and L2
!Included pw-pw2, pw-soc, pw2-soc2, and pw-w diffusion exchange
!Included [HSOCL1-HSOCL2(mixed)] Di Toro's scenario

!Eroded SHOCPOC flux is not a constant and varies with SOC and HSOCPOC
!Eroded SPOCHOC flux is proportional to the product of OCeroded flux
/SOC *SHOCPOC
!SOC and SHOCPOC's size distribution is the same as organic carbon size
distribution
!Add SHOCPOC, SPOC

!continuous IC
!C=lof with (1-pi); Wf=df^nf-1*df^3-nf; alphaC=0.1;

!STORM tank 2 event 3 data also cal D50
!linear Dp domain volume setting
!Add diffusion equation
!Add HOC with non-equilibrium method (Baltimore Harbor PCB Model)
!add OC
!add C from C(i,j) to C(i,j,k)
!correct lofb and lof
!correct the proportionality between TSS and TVC
!normal rk4 Nt+1=Nt+(D1+2D2+2D3+D4)/6
program simpleflocculation
implicit none

INTEGER :: i,j,k,l,t,t1
INTEGER :: mfc=503,mxt,hr,lp,rk,rkt
REAL(KIND=8), ALLOCATABLE ::
N(:,:,:),C(:,:,:),OC(:,:,:),HOCPOC(:,:,:),HOCdissolved(:,:)
REAL(KIND=8), ALLOCATABLE ::
AC(:,:,:),SAC1(:,:,:),SAC2(:,:,:),ACD(:,:),SAC1D(:,:),SAC2D(:,:),addAC(
:),addAC2(:)
REAL(KIND=8), ALLOCATABLE ::
SACmixedS1(:),SACmixedS2(:),ACSink(:),ACSink2(:),ACCheck(:),ACWs(:)
```

```

REAL(KIND=8), ALLOCATABLE ::
HOCAC(:, :, :), HOCSAC1(:, :, :), HOCSAC2(:, :, :), HOCACD(:, :), HOCSAC1D(:, :), HO
CSAC2D(:, :),
REAL(KIND=8), ALLOCATABLE ::
diffsac1(:), diffsac2(:), addHOCAC(:), HSACmixedS1(:), HSACmixedS2(:), ACMTCT(:)
REAL(KIND=8), ALLOCATABLE ::
HOCACsink(:), HOCACsink2(:), diffHOCAC(:), HOCACcheck(:), HOCSACcheck(:)
REAL(KIND=8), ALLOCATABLE ::
tempgAC(:), ACgain(:), tempgHAC(:), HACgain(:)
REAL(KIND=8), ALLOCATABLE ::
SHOCPOC(:, :, :), SHOCPOCSink(:), SHOCPOCD(:, :), SHOCPOCcheck(:)
REAL(KIND=8), ALLOCATABLE ::
SOC(:, :, :), SOCflux(:), SOCD(:, :), diffspoc1(:)
REAL(KIND=8), ALLOCATABLE ::
SOC2(:, :, :), SOC2D(:, :), SOCmixedS1(:), SOCmixedS2(:)
REAL(KIND=8), ALLOCATABLE ::
HSPOC2(:, :, :), HSPOC2D(:, :), diffspoc2(:), HSOCmixedS1(:), HSOCmixedS2(:)
REAL(KIND=8), ALLOCATABLE :: pw1D(:, :), pw2D(:, :), pw1(:, :), pw2(:, :),
REAL(KIND=8), ALLOCATABLE ::
foc(:, :, :), lof(:, :), Ws(:, :), alfa(:, :, :), D(:, :), ocD(:, :),
REAL(KIND=8), ALLOCATABLE ::
vcf(:), pro(:), Df(:), beta(:, :), Rf(:), tempgc(:), lofb(:, :), V(:)
REAL(KIND=8), ALLOCATABLE :: At(:), Ain(:), Aext(:), MTCT(:), pro2(:)
REAL(KIND=8), ALLOCATABLE ::
gain(:), lost(:), OCgain(:), tempgain(:), tempOClost(:)
REAL(KIND=8), ALLOCATABLE ::
tempgHOC(:), HOCPOCgain(:), HOCPOClost(:), HOCPOCD(:, :), HOCPOCcheck(:)
REAL(KIND=8), ALLOCATABLE ::
diffHOCPOC(:), HOCwD(:, :), totalHOC(:), addN(:), addOC(:), addHOC(:)
REAL(KIND=8), ALLOCATABLE ::
Masscheck(:), Ncheck(:), OCcheck(:), ndw(:, :), nfvc(:, :),
REAL(KIND=8), ALLOCATABLE :: D50(:), SOCcheck(:), addOC2(:)
REAL(KIND=8) :: loAC, ACKOC,
diffgainsac1, diffgainsac2, diffgainac, proAC=5.5d-1
REAL(KIND=8) :: W12b, ACarea=9.96d2, ACarea2=1d-2, ACfoc=0.44!0.88
REAL(KIND=8) :: alfad, alfac, los, low, lod, pi, fs, dt, nf=2.416!2.416d0! for
free settling and 2.545 for erode
REAL(KIND=8) :: g, z, sz, sz2, spro, mu, shear, Dp, Dc, dDf, diffsw, diffPWPW
REAL(KIND=8) ::
temp4, temp7, temp10, temp11, temp12, temp13, temp14, temp15, temp16, temp22
REAL(KIND=8) ::
MTCext, MTCin, skk, kksw, KOC, SKOC, HOCtemp, diffgain, diffgainS1, diffgainS2, A
AVCF, AVCF(0:600)
REAL(KIND=8) :: W12, addHOCadj=7.9d-1, Madj=1.41d-3!1.3d-3*1d0!Madj= 1/
time to the first LISST reading (25 min)
REAL(KIND=8) ::
rkct1=1d0, rkct2=1d0, rkct3=1d0, rkct4=1d0, rkct5=1d0, rkct6=1d0, rkct7=1d0, r
kct8=1d0, rkct9=1d0
REAL(KIND=8) :: Dm=4.92e-10, L0=8d1, hfc=5.2d-1, nf2=1.15d0
!REAL(KIND=8) :: Dm=4.92e-7!PCB 52=4.92e-6; PCB 4,10=5.439e-6; PCB
19=5.1634; PCB 110=4.7096
!rkct1 controls the flocs sinking, rkct2 controls the POC-
water diffusion
!rkct3 controls the sediment-pw diffusion (L1), rkct4 controls the pw1
and pw2 diffusion

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```

!rkct5 controls the sediment-pw diffusion (L2), rkct6 controls the pw-
water diffusion
!rkct7 controls the SAC L1-L2 exchange          , rkct8 controls the
HOCAC to others connection
!rkct9 controls the sediment-water interaction for particles
interface
function lalfa(a,b,c,d)
implicit none
REAL(KIND=8), INTENT(IN) :: a,b,c,d
REAL(KIND=8) :: lalfa
END function lalfa

function temp5(a)
implicit none
REAL(KIND=8), INTENT(IN) :: a
REAL(KIND=8) :: temp5
END function temp5

function temp6(h,a,b,c,d,e,f,g)
implicit none
REAL(KIND=8) :: temp6
REAL(KIND=8), INTENT(IN) :: h,a,b,c,d,f,g
INTEGER, INTENT(IN) :: e
END function temp6

function temp6b(h,a,b,e)
implicit none
REAL(KIND=8), INTENT(IN) :: h,a,b
INTEGER, INTENT(IN) :: e
Real(kind=8) :: temp6b

END function temp6b

function temp6c(d,a,b,c,e)
implicit none
REAL(KIND=8), INTENT(IN) :: a,b,c,d
INTEGER, INTENT(IN) :: e
REAL(KIND=8) :: temp6c
END function temp6c

function temp6d(a,b,e)
implicit none
REAL(KIND=8) :: temp6d
REAL(KIND=8), INTENT(IN) :: a,b
INTEGER, INTENT(IN) :: e
END function temp6d

function temp6e(a,b,c,e)
implicit none
REAL(KIND=8), INTENT(IN) :: a,b,c
INTEGER, INTENT(IN) :: e
REAL(KIND=8) :: temp6e
END function temp6e

function temp6f(a,b,c,d,e)
implicit none
REAL(KIND=8), INTENT(IN) :: a,b,c,d

```

```

INTEGER, INTENT(IN) :: e
REAL(KIND=8) :: temp6f
END function temp6f

function temp8(a)
implicit none
REAL(KIND=8), INTENT(IN) :: a
REAL(KIND=8) :: temp8
END function temp8

function temp20(a)
implicit none
REAL(KIND=8), INTENT(IN) :: a
REAL(KIND=8) :: temp20
END function temp20

function temp21(a,b)
implicit none
REAL(KIND=8), INTENT(IN) :: a,b
REAL(KIND=8) :: temp21
END function temp21
END interface

INTEGER :: temp1,temp2

hr=53
! hr run
WRITE(*,*) hr,'hours. simulation in this test'
WRITE(*,*) ' It is running now '

!OPEN
(30,file='T2E3IC.csv',status='old',access='sequential',form='formatted'
)
OPEN
(30,file='T2E3ICerodeP2.csv',status='old',access='sequential',form='for
matted')
!OPEN (31,FILE='T2E3ErodeFungchikPCB52evenAP2.out',
STATUS='unknown',ACCESS='sequential',FORM= 'formatted')
OPEN (31,FILE='AC1run6.out', STATUS='unknown',ACCESS='sequential',FORM=
'formatted')

!dt=2.5d-1 !dt(sec) = 60 sec. x 60=1 hr
!mxt=240 !for one min cycle (if dt=60 and mxt=60)
!dt=3.6d2
!mxt=10
!dt=2.5d-1

!dt=3.6d2
!mxt=24

dt=6d1
mxt=60

!60 units-corresponding to 3600 sec with dt=60 sec: total simulate
3600sec*hr== ?hr
!basic time unit = sec

```

```

ALLOCATE
(N(0:mfc,0:mxt,0:4),c(0:mfc,0:mxt,0:4),OC(0:mfc,0:mxt,0:4),HOCPOC(0:mfc,0:mxt,0:4))
ALLOCATE
(AC(0:mfc,0:mxt,0:4),SAC1(0:mfc,0:mxt,0:4),SAC2(0:mfc,0:mxt,0:4),ACMTCT(0:mfc))
ALLOCATE
(ACD(0:4,0:mfc),SAC1D(0:4,0:mfc),SAC2D(0:4,0:mfc),ADDAC(0:mfc),ADDAC2(0:mfc))
ALLOCATE
(SACmixedS1(0:mfc),SACmixedS2(0:mfc),ACSink(0:mfc),ACCheck(0:mxt))
!ALLOCATE (ACWs(0:mfc))
ALLOCATE
(ACSink2(0:mfc),HOCACD(0:4,0:mfc),HOCSAC1D(0:4,0:mfc),HOCSAC2D(0:4,0:mfc))
ALLOCATE
(HOCAC(0:mfc,0:mxt,0:4),HOCSAC1(0:mfc,0:mxt,0:4),HOCSAC2(0:mfc,0:mxt,0:4))
ALLOCATE
(diffsacl(0:mfc),diffsac2(0:mfc),addHOCAC(0:mfc),HSACmixedS1(0:mfc),HSACmixedS2(0:mfc))
ALLOCATE
(HOCACsink(0:mfc),HOCACsink2(0:mfc),diffHOCAC(0:mfc),HOCACcheck(0:mxt),HOCSACcheck(0:mxt))
ALLOCATE (tempgAC(0:mfc),ACgain(0:mfc),tempgHAC(0:mfc),HACgain(0:mfc))
ALLOCATE
(D(0:4,0:mfc),ocD(0:4,0:mfc),HOCPOCD(0:4,0:mfc),HOCwD(0:4,600))
ALLOCATE (At(0:mfc),Ain(0:mfc),Aext(0:mfc),MTCT(0:mfc))
allocate (addN(0:mfc),addOC(0:mfc),addOC2(0:mfc),addHOC(0:mfc))
allocate
(Masscheck(0:mxt),OCcheck(0:mxt),Ncheck(0:mxt),totalHOC(0:mxt))
ALLOCATE (D50(0:mxt),SOCcheck(0:mxt),SHOCPOCcheck(0:mxt),pro2(0:mfc))
ALLOCATE
(SOCD(0:4,0:mfc),SOC(0:mfc,0:mxt,0:4),SOCflux(0:mfc),diffspoc1(0:mfc))
allocate
(SHOCPOCD(0:4,0:mfc),SHOCPOC(0:mfc,0:mxt,0:4),SHOCPOCSink(0:mfc))
ALLOCATE (pw1D(0:4,600),pw2D(0:4,600),pw1(0:mxt,0:4),pw2(0:mxt,0:4))
ALLOCATE
(tempgHOC(0:mfc),HOCPOCgain(0:mfc),HOCPOClost(0:mfc),HOCPOCcheck(0:mxt))
)
ALLOCATE
(diffHOCPOC(0:mfc),HOCdissolved(0:mxt,0:4),ndw(0:mfc,0:mfc),nfvc(0:mfc,0:mfc))
allocate
(foc(0:mfc,0:mxt,0:4),lof(0:mfc,0:mxt),Ws(0:mfc,0:mxt),alfa(0:mfc,0:mxt,0:3))
ALLOCATE
(vcf(0:mfc),pro(0:mfc),Df(0:mfc),beta(0:mfc,0:mfc),Rf(0:mfc),V(0:mfc),lofb(0:mfc,0:mxt))
ALLOCATE
(tempOClost(0:mfc),gain(0:mfc),tempgain(0:mfc),lost(0:mfc),tempgc(0:mfc),OCgain(0:mfc))
ALLOCATE
(SOC2D(0:4,0:mfc),SOC2(0:mfc,0:mxt,0:4),SOCmixedS1(0:mfc),SOCmixedS2(0:mfc))

```

```

ALLOCATE
(HSPOC2D(0:4,0:mfc),HSPOC2(0:mfc,0:mxt,0:4),diffspoc2(0:mfc),HSOCmixedS
1(0:mfc),HSOCmixedS2(0:mfc))
!clean all variables
do i=1,mfc
  do j=1,mxt
    do k=0,3
      N(i,j,k)=0      !#/m^3
      C(i,j,k)=0      !g/m^3
      OC(i,j,k)=0     !g-OC/m^3
      SOC(i,j,k)=0
      SOC2(i,j,k)=0
      AC(i,j,k)=0     !g-AC/m^3
      SAC1(i,j,k)=0
      SAC2(i,j,k)=0
      HOCAC(i,j,k)=0  !ug-PCB/m^3 = ng/L
      HOCSAC1(i,j,k)=0
      HOCSAC2(i,j,k)=0
      HOCPOC(i,j,k)=0 !ug-PCB/m^3 = ng/L
      SHOCPOC(i,j,k)=0!ug-PCB/m^3
      HSPOC2(i,j,k)=0
      alfa(i,j,k)=0  ![0,1]
      foc(i,j,k)=0  ![0,0.5]
    end do
    lof(i,j)=0      !g/m^3 bulk dry density
    lofb(i,j)=0     !g/m^3 bulk density
    Ws(i,j)=0       !m/sec
  end do
  do j=1,4
    D(j,i)=0
    ocD(j,i)=0
    ACD(j,i)=0
    SAC1D(j,i)=0
    SAC2D(j,i)=0
    HOCPOCD(j,i)=0
    SHOCPOCD(j,i)=0
    SOCD(j,i)=0
    HOCwD(j,i)=0
  end do
  vcf(i)=0         !Vfs/V dry volume concentrations
  !vcf2(i)=0       !Vfb/V bulk volume concentrations = vcf/(1-pro)
  pro(i)=0         ![0,1]
  Df(i)=0          !um, 1e-6 m
  rf(i)=0
  V(i)=0           !m^3
  At(i)=0          !m^2 total surface area
  Ain(i)=0         !m^2 internal surface area = At-Aext
  Aext(i)=0        !m^3 external surface area
  MTct(i)=0        !mass transfer velocity
  gain(i)=0
  lost(i)=0
  tempgc(i)=0
  OCgain(i)=0
  tempgain(i)=0
  Ocgain(i)=0
  tempgHOC(i)=0
  HOCPOCgain(i)=0

```

```

HOCPOClost(i)=0
tempgAC(i)=0
ACgain(i)=0
tempgHAC(i)=0
HACgain(i)=0
SHOCPOCSink(i)=0
diffHOCPOC(i)=0
do j=1,mfc
  beta(i,j)=0
end do
end do

do i=0,mxt
  Ncheck(i)=0
  Masscheck(i)=0
  Occheck(i)=0
  HOCPOCcheck(i)=0
  SHOCPOCcheck(i)=0
  SOCcheck(i)=0
! diffgain(i)=0
  do j=0,3
    HOCdissolved(i,j)=0
    pw1(i,j)=0
    pw2(i,j)=0
  end do
end do

!setup initial constant values
pi=3.1415925d0
spro=1.2d-1 ! sediment porosity
fs=pi/6d0
alfad=0!4d-1!8d-1!4d-1
alfac=0!1d-1!1d-1
g=9.8d0 !m/s^2
z=1d0 !m
sz=1d-2!1d-2!m: the top sediment layer thickness
sz2=(5d-2 - sz) !m: the second sediment layer thickness
mu= 1d0 !g/m-sec
los=2.65d6 !g/m^3
low=1d6 !g/m^3
lod=1.05d6 !g/m^3
loAC=1.96d6!g/m^3
shear=6.68d-2!1/sec
W12=0!1.4e-9/sz2 !m/sec: particle mixed rate[sediment flux] Ditoro
W12b=0!1.4e-9*sz2
templ6=1d0 - proAC
!logKoc=5.02 for PCB 19, 0.104713
!Koc=10^(-0.46) for PCB 52 as the unit of m^3/g: 0.34673685
!Koc=10^(0.3) for PCB 85 as the unit of m^3/g : 1.99526
!Koc=10^(0.74) for PCB 118 as the unit of m^3/g: 5.49541
!kk=1.5d2
MTCin=6.25d-8! watch not -5 to compare with single model run 5
!fungchi fast
MTCext=MTCin*(1d-2)
!kk=3.6d-6
!kk=3.6d-7 !fungchi slow
!kk=1.472d-6!abby

```

```

skk=MTCin*(1d-5)!*(1d-5)
kksw=2d-10 !m^2/sec Lake Michigan PCB model p.189
KOC=1.1d0!PCB 52
!KOC=1.315955
!KOC=3.4673685d-1
!KOC=1.99526d0
!KOC=5.49541d0
SKOC=KOC*spro
ACKOC=KOC*1d3
!set up initial values for time 0 -----
-----
do i=1,mfc
READ(30,*) vcf(i),foc(i,0,0),HOCPOC(i,0,0) !switch to HOCPOC in the
free settling scenario
end do

do i=1,mfc
HOCPOC(i,0,0)=0
SOC(i,0,0)=1d0*foc(i,0,0)/SUM(foc)*1d4/sz!*0.5
SHOCPOC(i,0,0)=1.404d1*foc(i,0,0)/SUM(foc)*1d4/sz*0.0001!(SKOC/(SKOC+AC
KOC))!because of we assumed 0.5AC+0.5OC=TOC *0.00115
SOC2(i,0,0)=1d0*foc(i,0,0)/SUM(foc)*1d4/sz2!*0.5
HSPOC2(i,0,0)=1.404d1*foc(i,0,0)/SUM(foc)*1d4/sz2!/4D0!*0.5
!based on [SOC] g-C/m^3, foc=5%, and measured sediment PCB 52 =702 ng/g
and 14.04 ug/g-OC
SAC1(i,0,0)=0!1d0*foc(i,0,0)/SUM(foc)*1d4/sz!/4D0!*0.5
HOCSAC1(i,0,0)=0!1.404d1*foc(i,0,0)/SUM(foc)*1d4/sz*0.9999!(ACKOC/(SKOC
+ACKOC))*rkct8!*0.99885
SAC2(i,0,0)=0!1d0*foc(i,0,0)/SUM(foc)*1d4/sz2!/4D0!*0.5
HOCSAC2(i,0,0)=0!1.404d1*foc(i,0,0)/SUM(foc)*1d4/sz2!/4D0!*0.5*rkct8
end do

SAC1(75,0,0)=1d4/sz
SAC2(75,0,0)=1d4/sz2
HOCSAC1(75,0,0)=1.404d1*1d4/sz*0.9999
HOCSAC2(75,0,0)=1.404d1*1d4/sz!*0.9999

WRITE(*,*) SHOCPOC(75,0,0),HOCSAC1(75,0,0)
HOCdissolved(0,0)=0 !ug/m^3-water
pw1(0,0)=0!1d0*SUM(SHOCPOC)/SUM(SOC)/SKOC!ug/m^3-sed
pw2(0,0)=0!1d0*SUM(HSPOC2)/SUM(SOC2)/SKOC !ug/m^3-sed
!-----
-----
!setup Dp,dDf
Dp=(2d-6) ! 2 um dimaeter
dDf=(2d-6) ! 2 um diameter increasement

Df(1)=Dp
V(1)=fs*(1d-18)*(Df(1)*(1d6))**(3d0) !micrometer^3
pro(1)=0
pro2(1)=0
Aext(1)=4d0*pi*(Df(1)**(2d0))
At(1)=Aext(1)
Ain(1)=0
MTCT(1)=MTCext
! the interval will increase with linear scale

```

```

do i=2,mfc
Df(i)=Dp+(i-1)*dDf          !linear interval
V(i)=fs*(1d-18)*(Df(i)*(1d6))**(3d0) !micrometer^3
pro(i)=(1d0 -((Df(i)/Dp)**nf)*((Dp/Df(i))**(3d0)))
pro2(i)=(1d0 -((Df(i)/Dp)**nf2)*((Dp/Df(i))**(3d0)))
Aext(i)=4d0*pi*(Df(i)**(2d0))
Ain(i)=4d0*pi*(Df(i)**(nf-1d0))*(Dp**(3d0-nf))
At(i)=Ain(i)*0+Aext(i)
end do
!pro2(75)=0.55
!WRITE(*,'(10E15.7)') At(500),Ain(500),Aext(500)
!set up Df's range index
Rf(1)=Df(1)
do i=2,mfc
j=i-1
Rf(i)=Df(j)+(Df(i)-Df(j))*5d-1
end do

!precalculate the new floc volume range to which class of flocs
do i =1,mfc
do j=1,mfc
!V1+V2=V3 volume conservation rule
Dc=0
Dc=((Df(i)**3d0)+(Df(j)**3d0))**(1d0/3d0)

if (Dc >= Rf(mfc)) then          !set up upper boundary limit
nfvc(i,j)=mfc
ndw(i,j)=1d0!(V(i)+V(j))/V(mfc)
else
do l=1,mfc-1
temp2=l+1
if (Dc >= Rf(l) .and. Dc < Rf(temp2)) then
nfvc(i,j)=l
ndw(i,j)=1d0
end if
end do
end if
end do
end do

!-----
do i=1,mfc
foc(i,0,0)=foc(i,0,0)*6d-1*5d-1!(6.95d-1)!0.56 for free settling and
0.78 for erode
alfa(i,0,0)=2d0*foc(i,0,0)*alfad+(1d0 -2d0*foc(i,0,0))*alfac
lof(i,0)=2d0*foc(i,0,0)*lod+(1d0 -foc(i,0,0)*2d0)*los
lofb(i,0)=lof(i,0)*(1d0 -pro(i))+low*pro(i)
!Ws(i,0)=(lof(i,0)-low)*(Df(i)*Df(i))*9.8d0/1.8d1      !Stokes's law eqn
Ws(i,0)=(lof(i,0)-low)*((Df(i)**(nf-1d0))*(Df(1)**(3d0 -
nf)))*g/1.8d1*hfc/mu !Winterwerp's Ws eqn
!ACWs(i)= (loAC-low)*((Df(i)**(nfAC-1d0))*(Df(1)**(3d0 -
nfAC)))*g/1.8d1*hfc/mu
N(i,0,0)=vcf(i)/V(i)
N(i,0,0)=temp5(N(i,0,0))
C(i,0,0)=(lof(i,0)*Vcf(i))*(1d0-pro(i))
C(i,0,0)=temp5(C(i,0,0))
OC(i,0,0)=C(i,0,0)*foc(i,0,0)

```

```

OC(i,0,0)=temp5(OC(i,0,0))
AC(i,0,0)=OC(i,0,0) !!!temp assumption that AC=OC for erosion rate
! or name AC value later
AC(i,0,0)=temp5(AC(i,0,0))
MTct(i)=Dm/(Df(i)/L0)*(pro2(i)*pro2(i)/((1d0-
pro2(i))*Koc*foc(i,0,0)*lof(i,0)+pro2(i)))
!
addN(i)=N(i,0,0)*(6.77d-2*1d0/2.459d-2 -1d0)*(Madj)*rkct9 !lisst T1
data #/m^3/sec
addOC(i)=OC(i,0,0)*(6.77d-2*1d0/2.459d-2 -1d0)*(Madj)*rkct9 !Lisst T1
data #/m^3/sec
addOC2(i)=addOC(i)
addAC(i)=0!AC(i,0,0)*(6.77d-2*1d0/2.459d-2 -1d0)*(Madj)*rkct9 !Lisst
T1 data #/m^3/sec
addAC2(i)=0!addAC(i)

if (SOC(i,0,0)>0) then
addHOC(i)=SHOCPOC(i,0,0)*addOC(i)/SOC(i,0,0)*addHOCadj*rkct9!set up
initial eroded HOC flux rate
else
addHOC(i)=0
end if

addHOCAC(i)=0

N(i,0,0)=0
C(i,0,0)=0
AC(i,0,0)=0
end do

AC(75,0,0)=SUM(OC)
addAC(75)=AC(75,0,0)*(6.77d-2*1d0/2.459d-2 -1d0)*(Madj)*rkct9 !Lisst
T1 data #/m^3/sec
addAC2(75)=addAC(75)
addHOCAC(75)=HOCSAC1(75,0,0)*addAC(75)/SAC1(75,0,0)*addHOCadj*rkct8*rkct9!set up initial eroded HOC flux rate
!
AC(75,0,0)=0
!
do i=1,mfc
OC(i,0,0)=0
end do
!AC(75,0,0)=4.4d-1
!HOCAC(75,0,0)=1d0
WRITE(*,'(4E14.6)') OC(75,0,0),AC(75,0,0),addHOC(75),addHOCAC(75)
temp22=Ws(mfc,0)

!-----Start main program loop
outter: do lp =1,hr,1
time: do t =1,mxt
t1=t-1
rk4: do rk=1,4
rkt=rk-1

!clean
do i=1,mfc
gain(i)=0
tempgain(i)=0

```

```

tempgc(i)=0
OCgain(i)=0
lost(i)=0
tempgHOC(i)=0
HOCPOCgain(i)=0
HOCPOClost(i)=0
tempgAC(i)=0
ACgain(i)=0
tempgHAC(i)=0
HACgain(i)=0
diffspoc1(i)=0
diffspoc2(i)=0
SOCmixedS1(i)=0
SOCmixedS2(i)=0
HSOCmixedS1(i)=0
HSOCmixedS2(i)=0
D(rk,i)=0
ocD(rk,i)=0
ACD(rk,i)=0
SAC1D(rk,i)=0
SAC2D(rk,i)=0
HOCACD(rk,i)=0
HOCSAC1D(rk,i)=0
HOCSAC2D(rk,i)=0
HOCPOCD(rk,i)=0
HOCwD(rk,i)=0
HSPOC2D(rk,i)=0
SOC2D(rk,i)=0
END do

!calculate beta
do i=1,mfc
  do j =1,mfc
    beta(i,j)=0
    temp10=0
    temp10=Ws(i,t1)-Ws(j,t1)
    if (temp10 >= 0) then
      temp10=temp10
    else
      temp10= -temp10
    end if
  !
  beta(i,j)=temp5(pi*((Df(i)+Df(j))**2d0)*(temp10))+temp5(4d0/3d0*((Df(i)
+Df(j))**3d0)*(shear))
  beta(i,j)=temp5(pi*((Df(i)+Df(j))**2d0)*temp21(temp10,temp22))*(2.5d-
1)*(1d-2)+temp5(4d0/3d0*((Df(i)+Df(j))**3d0)*(shear))
  end do
end do

!calculate gain
temp1=mfc-1
do i=1,temp1
  do j=1,temp1

    k=0
    k=nfvc(i,j)

```

```

    tempgain(k)=beta(i,j)*N(i,t1,rkt)*N(j,t1,rkt)*&
lalfa(alfa(i,t1,rkt),alfa(j,t1,rkt),N(i,t1,rkt),N(j,t1,rkt))*ndw(i,j)

    gain(k)=gain(k)+tempgain(k)

    temp10=0
    temp10= lalfa(alfa(i,t1,rkt),alfa(j,t1,rkt),N(i,t1,rkt),N(j,t1,rkt))

tempgc(k)=beta(i,j)*temp10*(N(i,t1,rkt)*OC(j,t1,rkt)+N(j,t1,rkt)*OC(i,t
1,rkt))*ndw(i,j)
    OCgain(k)=OCgain(k)+tempgc(k)

tempgHOC(k)=beta(i,j)*temp10*(N(i,t1,rkt)*HOCPOC(j,t1,rkt)+N(j,t1,rkt)*
HOCPOC(i,t1,rkt))*ndw(i,j)
    HOCPOCgain(k)=HOCPOCgain(k)+tempgHOC(k)

tempgAC(k)=beta(i,j)*temp10*(N(i,t1,rkt)*AC(j,t1,rkt)+N(j,t1,rkt)*AC(i,
t1,rkt))*ndw(i,j)
    ACgain(k)=ACgain(k)+tempgAC(k)

tempgHAC(k)=beta(i,j)*temp10*(N(i,t1,rkt)*HOCAC(j,t1,rkt)+N(j,t1,rkt)*H
OCAC(i,t1,rkt))*ndw(i,j)
    HACgain(k)=HACgain(k)+tempgHAC(k)

    end do
    end do

!calculatae lost
do i=1,temp1
    do j=1,mfc
        lost(i)=lost(i)+N(i,t1,rkt)*N(j,t1,rkt)*beta(i,j)*&
        lalfa(alfa(i,t1,rkt),alfa(j,t1,rkt),N(i,t1,rkt),N(j,t1,rkt))*floc
    lost
    end do

tempOClost(i)=0
HOCPOClost(i)=0
ACsink(i)=0
ACsink2(i)=0
HOCACsink(i)=0
HOCACsink2(i)=0

if (N(i,t1,rkt)>0) then
tempOClost(i)=(lost(i)/N(i,t1,rkt)+Ws(i,t1)/z*rkct1)*OC(i,t1,rkt)*rkct9
!OC lost
SOCflux(i)= (Ws(i,t1)/z*rkct1)*OC(i,t1,rkt)*z/sz*rkct9

HOCPOClost(i)=(lost(i)/N(i,t1,rkt)+Ws(i,t1)/z*rkct1)*HOCPOC(i,t1,rkt)*r
kct9
!HOCPOC lost
SHOCPOCsink(i)= (Ws(i,t1)/z*rkct1)*HOCPOC(i,t1,rkt)*z/sz*rkct9 !
convert it to the sediment concentration

```

```

ACsink(i)= (lost(i)/N(i,t1,rkt)+Ws(i,t1)/z*rkct1)*AC(i,t1,rkt)*rkct9
ACsink2(i)= (Ws(i,t1)/z*rkct1)*AC(i,t1,rkt)*z/sz*rkct9
HOCACsink(i)=(lost(i)/N(i,t1,rkt)+Ws(i,t1)/z*rkct1)*HOCAC(i,t1,rkt)*rkct8*rkct9
HOCACsink2(i)=(Ws(i,t1)/z*rkct1)*HOCAC(i,t1,rkt)*z/sz*rkct8*rkct9
else
tempOClost(i)=(Ws(i,t1)/z*rkct1)*OC(i,t1,rkt)*rkct9
SOCflux(i)= (Ws(i,t1)/z*rkct1)*OC(i,t1,rkt)*z/sz*rkct9
HOCPOClost(i)=(Ws(i,t1)/z*rkct1)*HOCPOC(i,t1,rkt)*rkct9
SHOCPOCsink(i)=(Ws(i,t1)/z*rkct1)*HOCPOC(i,t1,rkt)*z/sz*rkct9
ACsink(i)= (Ws(i,t1)/z*rkct1)*AC(i,t1,rkt)*rkct9
ACsink2(i)= (Ws(i,t1)/z*rkct1)*AC(i,t1,rkt)*z/sz*rkct9
HOCACsink(i)=(Ws(i,t1)/z*rkct1)*HOCAC(i,t1,rkt)*rkct8*rkct9
HOCACsink2(i)=(Ws(i,t1)/z*rkct1)*HOCAC(i,t1,rkt)*z/sz*rkct8*rkct9
END if
end do

!calculate new N
if (rk<4) then !start if loop for rk4
do i=1,mfc
temp4=0
!find a way to cal D for rk.....

D(rk,i)=dt*temp6(addN(i),gain(i),lost(i),Ws(i,t1),z,rk,N(i,t1,rkt),rkct1)
temp4=N(i,t1,0)+D(rk,i)
N(i,t1,rk)=temp5(temp4)
end do

!calculate new variable at time t1 for rk loop
diffgain=0
diffgainS1=0
diffgainS2=0
diffgainac=0
diffgainsa1=0
diffgainsac2=0
diffPWPW=0

do i=1,mfc
!calculate new OC
temp12=0
temp12=addOC2(i)*dt*z/sz
if (temp12 < SOC(i,t1,rkt)) then
addOC(i)=addOC2(i)*rkct9
else
addOC(i)=0
end if

ocD(rk,i)=dt*temp6b(addOC(i),OCgain(i),tempOClost(i),rk)
OC(i,t1,rk)=OC(i,t1,0)+ocD(rk,i)
OC(i,t1,rk)=temp5(OC(i,t1,rk))

!calculate diffusion exchange
HOCTemp=0
MTct(i)=Dm/(Df(i)/L0)*(pro2(i)*pro2(i)/((1d0-pro2(i))*Koc*foc(i,t1,rkt)*lof(i,0)+pro2(i)))
if (OC(i,t1,rkt)>0) then

```

```

HOCTemp=HOCPOC(i,t1,rkt)/OC(i,t1,rkt)/KOC
diffHOCPOC(i)=At(i)*MTct(i)*N(i,t1,rkt)*(HOCdissolved(t1,rkt)-
HOCTemp)*rkct2!OC
else
HOCTemp=0
diffHOCPOC(i)=0
end if
diffgain=diffgain+diffHOCPOC(i)

!calculate new HOCPOC
temp12=0
temp12=addHOC(i)*dt*z/sz
if (temp12 < SHOCPOC(i,t1,rkt) .and. SOC(i,t1,rkt)>0) then
addHOC(i)=SHOCPOC(i,t1,rkt)*addOC(i)/SOC(i,t1,rkt)*addHOCadj*rkct9
else
addHOC(i)=0
end if
HOCPOCD(rk,i)=dt*temp6c(addHOC(i),HOCPOCgain(i),HOCPOClost(i),diffHOCPOC(i),rk)
HOCPOC(i,t1,rk)=HOCPOC(i,t1,0)+HOCPOCD(rk,i)
HOCPOC(i,t1,rk)=temp5(HOCPOC(i,t1,rk))

!Calculate new SOC
temp11=0
temp13=addOC(i)*z/sz

SOCmixedS1(i)=W12/sz*(SOC2(i,t1,rkt)-SOC(i,t1,rkt))
SOCD(rk,i)=dt*temp6e(SOCflux(i),temp13,SOCmixedS1(i),rk)!(gain from
sink,lost from erosion,gain from particle excahnge,rk)
SOC(i,t1,rk)=SOC(i,t1,0)+SOCD(rk,i)
SOC(i,t1,rk)=temp5(SOC(i,t1,rk))

!Calculate SOC layer 2
temp11=0
SOCmixedS2(i)= -W12/sz2*(SOC2(i,t1,rkt)-SOC(i,t1,rkt))
SOC2D(rk,i)=dt*temp6e(temp11,temp11,SOCmixedS2(i),rk)!(?,?,lost from
particle exchange, rk)
SOC2(i,t1,rk)=SOC2(i,t1,0)+SOC2D(rk,i)
SOC2(i,t1,rk)=temp5(SOC2(i,t1,rk))

!calculate sediment-porewater diffusion exchange Layer 1
HOCTemp=0
temp14=lof(i,0)*(ld0 - pro(i))*V(i)
if (SOC(i,t1,rkt)>0) then
HOCTemp=SHOCPOC(i,t1,rkt)/SOC(i,t1,rkt)/SKOC
diffspoc1(i)=At(i)*skk*SOC(i,t1,rkt)/temp14*(pw1(t1,rkt)-
HOCTemp)*rkct3!OC
else
diffspoc1(i)=0
end if
diffgainS1=diffgainS1+diffspoc1(i)!/spro

!calculate sediment-porewater diffusion exchange Layer 2
HOCTemp=0
if (SOC2(i,t1,rkt)>0) then
HOCTemp=HSPOC2(i,t1,rkt)/SOC2(i,t1,rkt)/SKOC

```

```

diffspoc2(i)=At(i)*skk*SOC2(i,t1,rkt)/temp14*(pw2(t1,rkt)-
HOCTemp)*rkct5!OC
else
diffspoc2(i)=0
end if
diffgainS2=diffgainS2+diffspoc2(i)!/spro

!Calculate new SHOCPOC layer 1
temp13=addHOC(i)*z/sz
HSOCmixedS1(i)=W12/sz*(HSPOC2(i,t1,rkt)-SHOCPOC(i,t1,rkt))
SHOCPOCD(rk,i)=dt*temp6f(SHOCPOCSink(i),temp13,diffspoc1(i),HSOCmixedS1
(i),rk)
!(gain from sink,lost from erosion and mixed,gain from diffusion,rk)
SHOCPOC(i,t1,rk)=SHOCPOC(i,t1,0)+SHOCPOCD(rk,i)

!Calculate new HSPOC layer 2
temp11=0
HSOCmixedS2(i)= -W12/sz2*(HSPOC2(i,t1,rkt)-SHOCPOC(i,t1,rkt))
HSPOC2D(rk,i)=dt*temp6e(HSOCmixedS2(i),temp11,diffspoc2(i),rk)!(?,?,
gain from diffusion,rk)
HSPOC2(i,t1,rk)=HSPOC2(i,t1,0)+HSPOC2D(rk,i)

!calculate new AC
temp12=0
temp12=addAC2(i)*dt*z/sz
if (temp12 < SAC1(i,t1,rkt)) then
addAC(i)=addAC2(i)*rkct9
else
addAC(i)=0
end if

ACD(rk,i)=dt*temp6b(addAC(i),ACgain(i),ACsink(i),rk)
AC(i,t1,rk)=AC(i,t1,0)+ACD(rk,i)
AC(i,t1,rk)=temp5(AC(i,t1,rk))

!Calculate new SAC
temp11=0
temp13=addAC(i)*z/sz

SACmixedS1(i)=W12/sz*(SAC2(i,t1,rkt)-SAC1(i,t1,rkt))*rkct7
SAC1D(rk,i)=dt*temp6e(ACsink2(i),temp13,SACmixedS1(i),rk)!(gain from
sink,lost from erosion,gain from particle excahnge,rk)
SAC1(i,t1,rk)=SAC1(i,t1,0)+SAC1D(rk,i)
SAC1(i,t1,rk)=temp5(SAC1(i,t1,rk))

!Calculate SAC layer 2
temp11=0
SACmixedS2(i)= -W12/sz2*(SAC2(i,t1,rkt)-SAC1(i,t1,rkt))*rkct7
SAC2D(rk,i)=dt*temp6e(temp11,temp11,SACmixedS2(i),rk)!(?,?,lost from
particle exchange, rk)
SAC2(i,t1,rk)=SAC2(i,t1,0)+SAC2D(rk,i)
SAC2(i,t1,rk)=temp5(SAC2(i,t1,rk))

!calculate AC diffusion exchange
HOCTemp=0
ACMTct(i)=Dm/(Df(i)/L0)*(proAC*proAC/((temp16)*ACKoc*ACfoc*loAC+proAC))
if (AC(i,t1,rkt)>0) then

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HOCTemp=HOCAC(i,t1,rkt)/AC(i,t1,rkt)/ACKOC
diffHOCAC(i)=ACarea*ACMTct(i)*AC(i,t1,rkt)*(HOCdissolved(t1,rkt)-
HOCTemp)*rkct8
!diffHOCAC(i)=At(i)*MTct(i)*N(i,t1,rkt)*(HOCdissolved(t1,rkt)-
HOCTemp)*rkct8
else
HOCTemp=0
diffHOCAC(i)=0
end if
diffgainac=diffgainac+diffHOCAC(i)*rkct8

!calculate new HOCAC
temp12=0
temp12=addHOCAC(i)*dt*z/sz
if (temp12 < HOCSAC1(i,t1,rkt) .and. SAC1(i,t1,rkt)>0) then
addHOCAC(i)=HOCSAC1(i,t1,rkt)*addAC(i)/SAC1(i,t1,rkt)*addHOCadj*rkct8*r
kct9
else
addHOCAC(i)=0
end if

HOCACD(rk,i)=dt*temp6c(addHOCAC(i),HACgain(i),HOCACsink(i),diffHOCAC(i)
,rk)
HOCAC(i,t1,rk)=HOCAC(i,t1,0)+HOCACD(rk,i)
!HOCAC(i,t1,rk)=temp5(HOCAC(i,t1,rk))

!calculate sediment AC-porewater diffusion exchange Layer 1
HOCTemp=0
temp14=loAC*(ld0 - proAC)*V(i)
if (SAC1(i,t1,rkt)>0) then
HOCTemp=HOCSAC1(i,t1,rkt)/SAC1(i,t1,rkt)/ACKOC
diffsacl(i)=ACarea2*skk*SAC1(i,t1,rkt)/temp14*(pw1(t1,rkt)-
HOCTemp)*rkct8!*rkct3!OC
else
diffsacl(i)=0
end if
diffgainSacl=diffgainSacl+diffsacl(i)*rkct8

!calculate sediment AC-porewater diffusion exchange Layer 2
HOCTemp=0
if (SAC2(i,t1,rkt)>0) then
HOCTemp=HOCSAC2(i,t1,rkt)/SAC2(i,t1,rkt)/ACKOC
diffsacl2(i)=ACarea2*skk*SAC2(i,t1,rkt)/temp14*(pw2(t1,rkt)-
HOCTemp)*rkct8!*rkct5!OC
else
diffsacl2(i)=0
end if
diffgainSac2=diffgainSac2+diffsacl2(i)*rkct8

!Calculate new HOCSAC layer 1
temp13=addHOCAC(i)*z/sz*rkct8
HSACmixedS1(i)=W12b/sz*(HOCSAC2(i,t1,rkt)-
HOCSAC1(i,t1,rkt))*rkct7*rkct8
HOCSAC1D(rk,i)=dt*temp6f(HOCACSink2(i),temp13,diffsacl(i),HSACmixedS1(i)
),rk)
!(gain from sink,lost from erosion and mixed,gain from diffusion,rk)
HOCSAC1(i,t1,rk)=HOCSAC1(i,t1,0)+HOCSAC1D(rk,i)

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```

!Calculate new HOCSAC layer 2
temp11=0
HSACmixedS2(i)= -W12b/sz2*(HOCSAC2(i,t1,rkt)-
HOCSAC1(i,t1,rkt))*rkct7*rkct8
HOCSAC2D(rk,i)=dt*temp6e(HSACmixedS2(i),temp11,diffsac2(i),rk)!(?,?,
gain from diffusion,rk)
HOCSAC2(i,t1,rk)=HOCSAC2(i,t1,0)+HOCSAC2D(rk,i)

!calculate Dissolved HOC concentration
diffsw=0
diffsw=kksw/(sz*z)*(HOCdissolved(t1,rkt)-pw1(t1,rkt))*rkct6!pw-w
diffusion exchange
temp15=0
temp15=diffgainAC+diffgain
HOCwD(rk,i)=dt*temp6d(temp15,diffsw,rk)
HOCdissolved(t1,rk)=HOCdissolved(t1,0)+HOCwD(rk,i)
!HOCdissolved(t1,rk)=temp5(HOCdissolved(t1,rk))

!calculate porewater HOC concentration L1
diffPWPW=0
diffPWPW= kksw/(sz*sz)*(pw1(t1,rkt)-pw2(t1,rkt))*rkct4!L1 and L2's
porewater exchange[sediment volume adjusted here]
diffsw= -kksw/(sz*sz)*(HOCdissolved(t1,rkt)-
pw1(t1,rkt))*rkct6!+diffPWPW !!! sediment volume adjusted here
temp15=0
temp15=diffgainS1+diffgainSa1!*rkct8
pw1D(rk,i)=dt*temp6e(-temp15,diffsw,-diffPWPW,rk) !watch out the + -
term in the function
pw1(t1,rk)=pw1(t1,0)+pw1D(rk,i)
!pw1(t1,rk)=temp5(pw1(t1,rk))
!calculate porewater HOC concentration L2
diffPWPW= -kksw/(sz*sz2)*(pw1(t1,rkt)-pw2(t1,rkt))*rkct4!L1 and L2's
porewater exchange[sediment volume adjusted here]
temp15=0
temp15=diffgainS2+diffgainSac2*rkct8
pw2D(rk,i)=dt*temp6d(temp15,diffPWPW,rk)
pw2(t1,rk)=pw2(t1,0)+pw2D(rk,i)

!calculate new foc
if (C(i,t1,rkt)>0) then
temp7=OC(i,t1,rk)/C(i,t1,rkt)
foc(i,t1,rk)=temp8(temp7)
else
foc(i,t1,rk)=0
END if
alfa(i,t1,rk)=2d0*foc(i,t1,rk)*alfad+(1d0 -2d0*foc(i,t1,rk))*alfac
end do
!foc(75,t1,rk)=0.44
!claculate other variables
do i=1,mfc
lof(i,t)=2d0*foc(i,t1,rk)*lod+(1d0 -2d0*foc(i,t1,rk))*los
lofb(i,t)=lof(i,t)*(1d0 -pro(i))+low*pro(i)
C(i,t1,rk)=N(i,t1,rk)*V(i)*lof(i,t)*(1d0-pro(i))
C(i,t1,rk)=temp5(C(i,t1,rk))
vcf(i)=N(i,t1,rk)*V(i)

```

```

! Ws(i,t)=(lof(i,t)-low)*(Df(i)*Df(i))*9.8d0/1.8d1
!*wseff(vcmx,vcp,vcf(i))
  Ws(i,t)=(lof(i,t)-low)*((Df(i)**(nf-1d0))*(Df(1)**(3d0 -
nf)))*g/1.8d1*(hfc)/mu
! Ws(mfc,t)=(lof(i,t)-low)*(Df(i)*Df(i))*9.8d0/1.8d1
!*wseff(vcmx,vcp,vcf(i))
end do

!
else          !!!!!!!!!! rk=4 below
!
  do i=1,mfc
    temp4=0
    D(4,i)=
dt*temp6(addN(i),gain(i),lost(i),Ws(i,t1),z,rk,N(i,t1,rkt),rkct1)
    temp4=N(i,t1,0)+(D(1,i)+2d0*D(2,i)+2d0*D(3,i)+D(4,i))/(6d0)
    N(i,t,0)=temp5(temp4)
  end do

!calculate new variable at time t
diffgain=0
diffgainS1=0
diffgainS2=0
diffgainac=0
diffgainSac1=0
diffgainSac2=0
diffPWPW=0

  do i=1,mfc
!calculate new POC
temp12=0
temp12=addOC2(i)*dt*z/sz
if (temp12 < SOC(i,t1,rkt)) then
addOC(i)=addOC2(i)*rkct9
else
addOC(i)=0
end if
ocD(rk,i)=dt*temp6b(addOC(i),OCgain(i),tempOClost(i),rk)
OC(i,t,0)=OC(i,t1,0)+(ocD(1,i)+2d0*ocD(2,i)+2d0*ocD(3,i)+ocD(4,i))/(6d0
)
OC(i,t,0)=temp5(OC(i,t,0))

!calculate diffusion exchange
HOCTemp=0
MTct(i)=Dm/(Df(i)/L0)*(pro2(i)*pro2(i)/((1d0-
pro2(i))*Koc*foc(i,t1,rkt)*lof(i,0)+pro2(i)))
if (OC(i,t1,rkt)>0) then
HOCTemp=HOCPOC(i,t1,rkt)/OC(i,t1,rkt)/KOC
diffHOCPOC(i)=At(i)*MTct(i)*N(i,t1,rkt)*(HOCdissolved(t1,rkt)-
HOCTemp)*rkct2
else
HOCTemp=0
diffHOCPOC(i)=0
end if
diffgain=diffgain+diffHOCPOC(i)

!calculate new HOCPOC

```

```

temp12=0
temp12=addHOC(i)*dt*z/sz
if (temp12 < SHOCPOC(i,t1,rkt) .and. SOC(i,t1,rkt)>0) then
addHOC(i)=SHOCPOC(i,t1,rkt)*addOC(i)/SOC(i,t1,rkt)*addHOCadj*rkct9
else
addHOC(i)=0
end if
HOCPOCD(rk,i)=dt*temp6c(addHOC(i),HOCPOCgain(i),HOCPOClost(i),diffHOCPOC(i),rk)
HOCPOC(i,t,0)=HOCPOC(i,t1,0)+(HOCPOCD(1,i)+2d0*HOCPOCD(2,i)+2d0*HOCPOCD(3,i)+HOCPOCD(4,i))/(6d0)
HOCPOC(i,t,0)=temp5(HOCPOC(i,t,0))

!Calculate new SOC
temp13=addOC(i)*z/sz

SOCmixedS1(i)=W12/sz*(SOC2(i,t1,rkt)-SOC(i,t1,rkt))
SOC2D(rk,i)=dt*temp6e(SOCflux(i),temp13,SOCmixedS1(i),rk)!(gain from sink, lost from erosion,gain from particle exchange, rk)
SOC(i,t,0)=SOC(i,t1,0)+(SOC2D(1,i)+2d0*SOC2D(2,i)+2d0*SOC2D(3,i)+SOC2D(4,i))/(6d0)
SOC(i,t,0)=temp5(SOC(i,t,0))

!Calculate SOC layer 2
temp11=0
SOCmixedS2(i)= -W12/sz2*(SOC2(i,t1,rkt)-SOC(i,t1,rkt))
SOC2D(rk,i)=dt*temp6e(temp11,temp11,SOCmixedS2(i),rk)!(?,?,lost from particle exchange, rk)
SOC2(i,t,0)=SOC2(i,t1,0)+(SOC2D(1,i)+2d0*SOC2D(2,i)+2d0*SOC2D(3,i)+SOC2D(4,i))/(6d0)
SOC2(i,t,0)=temp5(SOC2(i,t,0))

!calculate sediment-porewater diffusion exchange layer 1
HOCTemp=0
temp14=lof(i,0)*(1d0 - pro(i))*V(i)
if (SOC(i,t1,rkt)>0) then
HOCTemp=SHOCPOC(i,t1,rkt)/SOC(i,t1,rkt)/SKOC
diffspoc1(i)=At(i)*skk*SOC(i,t1,rkt)/temp14*(pw1(t1,rkt)-HOCTemp)*rkct3!OC
else
diffspoc1(i)=0
end if
diffgainS1=diffgainS1+diffspoc1(i)!/spro

!calculate sediment-porewater diffusion exchange layer 2
HOCTemp=0
if (SOC2(i,t1,rkt)>0) then
HOCTemp=HSPOC2(i,t1,rkt)/SOC2(i,t1,rkt)/SKOC
diffspoc2(i)=At(i)*skk*SOC2(i,t1,rkt)/temp14*(pw2(t1,rkt)-HOCTemp)*rkct5!OC
else
diffspoc2(i)=0
end if
diffgainS2=diffgainS2+diffspoc2(i)!

!Calculate new HOCSOC layer 1
temp13=addHOC(i)*z/sz

```

```

HSOCmixedS1(i)=W12/sz*(HSPOC2(i,t1,rkt)-SHOCPOC(i,t1,rkt))
SHOCPOCD(rk,i)=dt*temp6f(SHOCPOCSink(i),temp13,diffspoc1(i),HSOCmixedS1
(i),rk)
!(gain from sink,lost from erosion,gain from diffusion,gain from
mixed,rk)
SHOCPOC(i,t,0)=SHOCPOC(i,t1,0)+(SHOCPOCD(1,i)+2d0*SHOCPOCD(2,i)+2d0*SHO
CPOCD(3,i)+SHOCPOCD(4,i))/(6d0)

!Calculate new HOCSOC layer 2
HSOCmixedS2(i)= -W12/sz2*(HSPOC2(i,t1,rkt)-SHOCPOC(i,t1,rkt))
HSPOC2D(rk,i)=dt*temp6e(HSOCmixedS2(i),temp11,diffspoc2(i),rk)!(?,?,gai
n from diffusion,rk)
HSPOC2(i,t,0)=HSPOC2(i,t1,0)+(HSPOC2D(1,i)+2d0*HSPOC2D(2,i)+2d0*HSPOC2D
(3,i)+HSPOC2D(4,i))/(6d0)

!calculate new AC
temp12=0
temp12=addAC2(i)*dt*z/sz
if (temp12 < SAC1(i,t1,rkt)) then
addAC(i)=addAC2(i)*rkct9
else
addAC(i)=0
end if

ACD(rk,i)=dt*temp6b(addAC(i),ACgain(i),ACsink(i),rk)
AC(i,t,0)=AC(i,t1,0)+(ACD(1,i)+2d0*ACD(2,i)+2d0*ACD(3,i)+ACD(4,i))/(6d0
)
AC(i,t,0)=temp5(AC(i,t,0))

!calculate AC diffusion exchange
HOCTemp=0
ACMTct(i)=Dm/(Df(i)/L0)*(proAC*proAC/((temp16)*ACKoc*ACfoc*loAC+proAC))
if (AC(i,t1,rkt)>0) then
HOCTemp=HOCAC(i,t1,rkt)/AC(i,t1,rkt)/ACKOC
diffHOCAC(i)=ACarea*ACMTct(i)*AC(i,t1,rkt)*(HOCdissolved(t1,rkt)-
HOCTemp)*rkct8
!diffHOCAC(i)=At(i)*MTct(i)*N(i,t1,rkt)*(HOCdissolved(t1,rkt)-
HOCTemp)*rkct8
else
HOCTemp=0
diffHOCAC(i)=0
end if
diffgainac=diffgainac+diffHOCAC(i)*rkct8

!calculate new HOCAC
temp12=0
temp12=addHOCAC(i)*dt*z/sz
if (temp12 < HOCSAC1(i,t1,rkt) .and. SAC1(i,t1,rkt)>0) then
addHOCAC(i)=HOCSAC1(i,t1,rkt)*addAC(i)/SAC1(i,t1,rkt)*addHOCadj*rkct8*r
kct9
else
addHOCAC(i)=0
end if

HOCACD(rk,i)=dt*temp6c(addHOCAC(i),HACgain(i),HOCACsink(i),diffHOCAC(i)
,rk)

```

```

HOCAC(i,t,0)=HOCAC(i,t1,0)+(HOCACD(1,i)+2d0*HOCACD(2,i)+2d0*HOCACD(3,i)
+HOCACD(4,i))/(6d0)
!HOCAC(i,t,0)=temp5(HOCAC(i,t,0))

!Calculate new SAC
temp13=addAC(i)*z/sz
SACmixedS1(i)=W12/sz*(SAC2(i,t1,rkt)-SAC1(i,t1,rkt))*rkct7
SAC1D(rk,i)=dt*temp6e(ACSink2(i),temp13,SACmixedS1(i),rk)!(gain from
sink, lost from erosion,gain from particle exchange, rk)
SAC1(i,t,0)=SAC1(i,t1,0)+(SAC1D(1,i)+2d0*SAC1D(2,i)+2d0*SAC1D(3,i)+SAC1
D(4,i))/(6d0)
SAC1(i,t,0)=temp5(SAC1(i,t,0))

!Calculate SAC layer 2
temp11=0
SACmixedS2(i)= -W12/sz2*(SAC2(i,t1,rkt)-SAC1(i,t1,rkt))*rkct7
SAC2D(rk,i)=dt*temp6e(temp11,temp11,SACmixedS2(i),rk)!(?,?,lost from
particle exchange, rk)
SAC2(i,t,0)=SAC2(i,t1,0)+(SAC2D(1,i)+2d0*SAC2D(2,i)+2d0*SAC2D(3,i)+SAC2
D(4,i))/(6d0)
SAC2(i,t,0)=temp5(SAC2(i,t,0))

!calculate sediment AC-porewater diffusion exchange Layer 1
HOCTemp=0
temp14=loAC*(1d0 - proAC)*V(i)
if (SAC1(i,t1,rkt)>0) then
HOCTemp=HOCSAC1(i,t1,rkt)/SAC1(i,t1,rkt)/ACKOC
diffsacl(i)=ACarea2*skk*SAC1(i,t1,rkt)/temp14*(pw1(t1,rkt)-
HOCTemp)!*rkct8!*rkct3
else
diffsacl(i)=0
end if
diffgainSacl=diffgainSacl+diffsacl(i)!*rkct8

!calculate sediment AC-porewater diffusion exchange Layer 2
HOCTemp=0
if (SAC2(i,t1,rkt)>0) then
HOCTemp=HOCSAC2(i,t1,rkt)/SAC2(i,t1,rkt)/ACKOC
diffsac2(i)=ACarea2*skk*SAC2(i,t1,rkt)/temp14*(pw2(t1,rkt)-
HOCTemp)*rkct8!*rkct5
else
diffsac2(i)=0
end if
diffgainSac2=diffgainSac2+diffsac2(i)*rkct8

!Calculate new HOCSAC layer 1
temp13=addHOCAC(i)*z/sz*rkct8
HSACmixedS1(i)=W12b/sz*(HOCSAC2(i,t1,rkt)-
HOCSAC1(i,t1,rkt))*rkct7*rkct8
HOCSAC1D(rk,i)=dt*temp6f(HOCACSink2(i),temp13,diffsacl(i),HSACmixedS1(i)
),rk)
!(gain from sink,lost from erosion,gain from diffusion,gain from
mixed,rk)
HOCSAC1(i,t,0)=HOCSAC1(i,t1,0)+(HOCSAC1D(1,i)+2d0*HOCSAC1D(2,i)+2d0*HOC
SAC1D(3,i)+HOCSAC1D(4,i))/(6d0)

!Calculate new HOCSAC layer 2

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```

temp11=0
HSACmixedS2(i)= -W12b/sz2*(HOCSAC2(i,t1,rkt)-
HOCSAC1(i,t1,rkt))*rkct7*rkct8
HOCSAC2D(rk,i)=dt*temp6e(HSOCmixedS2(i),temp11,diffsac2(i),rk)!(?,?,gai
n from diffusion,rk)
HOCSAC2(i,t,0)=HOCSAC2(i,t1,0)+(HOCSAC2D(1,i)+2d0*HOCSAC2D(2,i)+2d0*HOC
SAC2D(3,i)+HOCSAC2D(4,i))/(6d0)

!calculate Dissolved HOC concentration
diffsw=0
diffsw=kksw/(sz*z)*(HOCdissolved(t1,rkt)-pw1(t1,rkt))*rkct6!porewater-
water diffusion exchange
temp15=0
temp15=diffgainAC+diffgain

HOCwD(rk,i)=dt*temp6d(temp15,diffsw,rk)
HOCdissolved(t,0)=HOCdissolved(t1,0)+(HOCwD(1,i)+2d0*HOCwD(2,i)+2d0*HOC
wD(3,i)+HOCwD(4,i))/(6d0)
!HOCdissolved(t,0)=temp5(HOCdissolved(t,0))

!calculate porewater HOC concentration L1
diffPWPW=0
diffPWPW= kksw/(sz*sz)*(pw1(t1,rkt)-pw2(t1,rkt))*rkct4! L1 and L2's
porewater exchange[sediment volume adjusted here]
diffsw= -kksw/(sz*sz)*(HOCdissolved(t1,rkt)-
pw1(t1,rkt))*rkct6!+diffPWPW!pw-w diffusion exchange[sediment and pw
volume adjusted included]
temp15=0
temp15=diffgainS1+diffgainSacl!*rkct8
pw1D(rk,i)=dt*temp6e(-temp15,diffsw,-diffPWPW,rk)
pw1(t,0)=pw1(t1,0)+(pw1D(1,i)+2d0*pw1D(2,i)+2d0*pw1D(3,i)+pw1D(4,i))/(6
d0)
!pw1(t,0)=temp5(pw1(t,0))
!calculate porewater HOC concentration L2
temp11=0
diffPWPW= -kksw/(sz*sz2)*(pw1(t1,rkt)-pw2(t1,rkt))*rkct4!L1 and L2's
porewater exchange[sediment volume adjusted here]
temp15=0
temp15=diffgainS2+diffgainSAC2*rkct8
pw2D(rk,i)=dt*temp6d(diffgainS2,diffPWPW,rk)
pw2(t,0)=pw2(t1,0)+(pw2D(1,i)+2d0*pw2D(2,i)+2d0*pw2D(3,i)+pw2D(4,i))/(6
d0)

!cal foc
if (C(i,t1,rkt)>0) then
temp7=OC(i,t,0)/C(i,t1,rkt)
foc(i,t,0)=temp8(temp7)
else
foc(i,t,0)=0
END if
!cal other parameters
alfa(i,t,0)=2d0*foc(i,t,0)*alfad+(1d0 -2d0*foc(i,t,0))*alfac
end do!
!foc(75,t,0)=0.44
do i=1,mfc
lof(i,t)=2d0*foc(i,t,0)*lod+(1d0 -2d0*foc(i,t,0))*los
lofb(i,t)=lof(i,t)*(1d0 -pro(i))+low*pro(i)

```

```

C(i,t,0)=N(i,t,0)*V(i)*lof(i,t)*(ld0-pro(i))
C(i,t,0)=temp5(C(i,t,0))
vcf(i)=N(i,t,0)*V(i)
Ws(i,t)=(lof(i,t)-low)*((Df(i)**(nf-ld0))*(Df(1)**(3d0 -
nf)))*9.8d0/1.8d1*(hfc)/mu
!Ws(mfc,t)=(lof(i,t)-low)*(Df(i)*Df(i))*9.8d0/1.8d1
!*wseff(vcmx,vcp,vcf(i))
end do

end if !end if for rk loop

end do rk4

!checking volume conservation
Ncheck(t1)=0
Masscheck(t1)=0
Occheck(t1)=0
SOCcheck(t1)=0
SHOCPOCcheck(t1)=0
HOCPOCcheck(t1)=0
totalHOC(t1)=0
ACcheck(t1)=0
HOCACcheck(t1)=0
HOCSACcheck(t1)=0

do i=1,mfc
Ncheck(t1)=Ncheck(t1)+vcf(i)
Masscheck(t1)=Masscheck(t1)+C(i,t1,0)
Occheck(t1)=Occheck(t1)+OC(i,t1,0)*z
SOCcheck(t1)=SOCcheck(t1)+SOC(i,t1,0)*sz+SOC2(i,t1,0)*sz2+OC(i,t1,0)
ACcheck(t1)=ACcheck(t1)+AC(i,t1,0)!+SAC1(i,t1,0)*sz+SAC2(i,t1,0)*sz2
HOCPOCcheck(t1)=HOCPOCcheck(t1)+HOCPOC(i,t1,0)*z
SHOCPOCcheck(t1)=SHOCPOCcheck(t1)+SHOCPOC(i,t1,0)*sz+HSPOC2(i,t1,0)*sz2
HOCACcheck(t1)=HOCACcheck(t1)+HOCAC(i,t1,0)*z
HOCSACcheck(t1)=HOCSACcheck(t1)+HOCSAC1(i,t1,0)*sz+HOCSAC2(i,t1,0)*sz2
end do
totalHOC(t1)=HOCdissolved(t1,0)*z+HOCACcheck(t1)+HOCPOCcheck(t1)+SHOCPOC
Ccheck(t1)+(pw1(t1,0)*sz+pw2(t1,0)*sz2)+HOCSACcheck(t1)
!calculate D50
AAVCF=0
D50(t1)=0
do i=1,mfc
AVCF(i)=vcf(i)/Ncheck(t1)
AAVCF=AAVCF+AVCF(i)
  if (AAVCF > 5d-1 ) then
    D50(t1)=Df(i)*(ld6)
    exit
  end if
end do
! end calculate D50

END do time ! End small time loop (mxt)

!output time interval 6 min
WRITE(*,*) lp
outputdata: do i=0,mxt-1,10
!outputdata: do i=1,mxt-1,12

```

```

write(31,'(13E14.6)')
Ncheck(i),OCcheck(i),Masscheck(i),ACcheck(i),HOCPOCcheck(i),HOCACcheck(
i),HOCdissolved(i,0),&
SHOCPOCcheck(i),pw1(i,0),pw2(i,0),totalHOC(i),SOCcheck(i),D50(i)

write(*,'(8E14.6)')
OCcheck(i),Masscheck(i),ACcheck(i),totalHOC(i),HOCdissolved(i,0),HOCACc
heck(i),HOCPOCcheck(i),pw1(i,0)
!write(31,'(4E14.6)')
totalHOC(i),SHOCPOC(75,mxt,0),HOCSAC1(75,mxt,0),pw1(i,0)
!write(*,'(4E14.6)')
totalHOC(i),SHOCPOC(75,mxt,0),HOCSAC1(75,mxt,0),pw1(i,0)
end do outputdata

!reset initial values as the latest time data
k=mxt-1
do i=1,mfc
N(i,0,0)=N(i,k,0)
C(i,0,0)=C(i,k,0)
OC(i,0,0)=OC(i,k,0)
SOC(i,0,0)=SOC(i,k,0)
SOC2(i,0,0)=SOC2(i,k,0)
AC(i,0,0)=AC(i,k,0)
SAC1(i,0,0)=SAC1(i,k,0)
SAC2(i,0,0)=SAC2(i,k,0)
HOCAC(i,0,0)=HOCAC(i,k,0)
HOCSAC1(i,0,0)=HOCSAC1(i,k,0)
HOCSAC2(i,0,0)=HOCSAC2(i,k,0)
HOCPOC(i,0,0)=HOCPOC(i,k,0)
SHOCPOC(i,0,0)=SHOCPOC(i,k,0)
HSPOC2(i,0,0)=HSPOC2(i,k,0)
Ws(i,0)=Ws(i,k)
foc(i,0,0)=foc(i,k,0)
alfa(i,0,0)=alfa(i,k,0)
lof(i,0)=lof(i,k)
lofb(i,0)=lofb(i,k)
end do
HOCdissolved(0,0)=HOCdissolved(k,0)
pw1(0,0)=pw1(k,0)
pw2(0,0)=pw2(k,0)
END do outter

!close (30, status = 'keep')
close (31, status = 'keep')
!close (32, status = 'keep')
!close (33, status = 'keep')
!close (34, status = 'keep')
!close (36, status = 'keep')
!close (37, status = 'keep')
stop
end program simpleflocculation

function lalfa(a,b,c,d)
implicit none
REAL(KIND=8), INTENT(IN) :: a,b,c,d
REAL(KIND=8) :: e,lalfa
e=c+d

```

```

if (e >0) then
lalfa=(a*c+b*d)/e
else
lalfa=b
end if
return
END function lalfa

function temp5(a)
implicit none
REAL(KIND=8), INTENT(IN) :: a
REAL(KIND=8) :: temp5
if (a >0) then
temp5=a
else
temp5=0
end if
return
END function temp5

function temp6(h,a,b,c,d,e,f,g)
implicit none
REAL(KIND=8) :: temp6
REAL(KIND=8), INTENT(IN) :: h,a,b,c,d,f,g
INTEGER, INTENT(IN) :: e
if (e == 2 .or. e == 3) then
temp6= (h+(5d-1)*a-b-c/d*f*g)*(5d-1)
ELSE
temp6= (h+(5d-1)*a-b-c/d*f*g)
end if
return
END function temp6

function temp6b(h,a,b,e)
implicit none
REAL(KIND=8), INTENT(IN) :: h,a,b
INTEGER, INTENT(IN) :: e
Real(kind=8) :: c,temp6b
c=5d-1
if (e == 2 .or. e == 3) then
temp6b= (h+c*a-b)*(5d-1)
ELSE
temp6b= (h+c*a-b)
end if
return
END function temp6b

function temp6c(d,a,b,c,e)
implicit none
REAL(KIND=8), INTENT(IN) :: a,b,c,d
INTEGER, INTENT(IN) :: e
REAL(KIND=8) :: temp6c

if (e == 2 .or. e == 3) then
temp6c= (d+(5d-1)*a-b+c)*(5d-1)
ELSE
temp6c= (d+(5d-1)*a-b+c)

```

```

end if
return
END function temp6c

function temp6d(a,b,e)
implicit none
REAL(KIND=8) :: temp6d
REAL(KIND=8), INTENT(IN) :: a,b
INTEGER, INTENT(IN) :: e

if (e == 2 .or. e == 3) then
  temp6d= -(a+b)*(5d-1)
ELSE
  temp6d= -(a+b)
end if
return
END function temp6d

function temp6e(a,b,c,e)
implicit none
REAL(KIND=8), INTENT(IN) :: a,b,c
INTEGER, INTENT(IN) :: e
REAL(KIND=8) :: temp6e

if (e == 2 .or. e == 3) then
  temp6e= (a-b+c)*(5d-1)
ELSE
  temp6e= (a-b+c)
end if
return
END function temp6e

function temp6f(a,b,c,d,e)
implicit none
REAL(KIND=8), INTENT(IN) :: a,b,c,d
INTEGER, INTENT(IN) :: e
REAL(KIND=8) :: temp6f

if (e == 2 .or. e == 3) then
  temp6f= (a-b+c+d)*(5d-1)
ELSE
  temp6f= (a-b+c+d)
end if
return
END function temp6f

function temp8(a)
implicit none
REAL(KIND=8), INTENT(IN) :: a
REAL(KIND=8) :: temp8
if (a > 5d-1) then ! replace 1.0 by 0.5
  temp8=5d-1
ELSE IF (a<0) then
  temp8=0
else
  temp8=a
end if
end if

```

```

return
END function temp8

function temp20(a)
implicit none
REAL(KIND=8), INTENT(IN) :: a
REAL(KIND=8) :: temp20
if (a > 1d0) then
temp20=1d0
ELSE IF (a<0) then
temp20=1d-5
else
temp20=a
end if
return
END function temp20

function temp21(a,b)
implicit none
REAL(KIND=8), INTENT(IN) :: a,b
REAL(KIND=8) :: temp21
if (a > b) then
temp21=b
ELSE IF (a<0) then
temp21=0
else
temp21=a
end if
return
END function temp21

```

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